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Hypersonic Flows for Reentry Problems

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Hypersonic Flows for Reentry Problems

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Contents

Lectures on experimental results are marked with *

Problems for Analysis

Flow over a Slender Cone	6
Problem 2: Turbulent Base Flow	8
Problem 3: Flow over a 2D Ramp	10
Problem 4: Flow over a 3D Obstacle	13
Problem 5: Corner Flow	15
Problem 6: Double (Simple) Ellipsoid	17
Problem 7: Flow over a Delta Wing	25
Problem 8: Non-equilibrium Flow in an Arc Jet or a Shock Tube	31
Appendix: Non-equilibrium Model for Problem 6	37

Problems 1-2

Hypersonic Boundary Layer and Base Flow* By P.A. Denman, J.K. Harvey, R. Hillier (With 14 Figures) . . .	43
Computation of Hypersonic Turbulent Flow over a Rearward Facing Step By M.P. Netterfield (With 11 Figures)	63
Hypersonic Cone Flow Predictions Using an Implicit Upwind Space-Marching Code By S.L. Lawrence (With 16 Figures)	75

Synthesis

A Synthesis of Results for Test Cases 1 and 2:
Hypersonic Boundary Layer and Base Flow
By O. Pironneau 92

Problems 3-4

Experiments on Shock-Wave/Boundary-Layer Interactions
Produced by Two-Dimensional Ramps
and Three-Dimensional Obstacles*
By J. Delery, M.-C. Coet (With 25 Figures) 97

An Experimental Contribution to the Flat Plate 2D Compression
Ramp, Shock/Boundary Layer Interaction Problem at Mach 14:
Test Case 3.7*
By G. Simeonides, J.F. Wendt (With 16 Figures) 129

Viscous, 2-D, Laminar Hypersonic Flows
over Compression Ramps
By Z. Alsalihi, H. Deconinck (With 16 Figures) 152

Computational Results for 2-D and 3-D Ramp Flows
with an Upwind Navier-Stokes Solver
By E. Venkatapathy (With 20 Figures) 167

Application of the Galerkin/Least-Squares Formulation
to the Analysis of Hypersonic Flows:
I. Flow over a Two-Dimensional Ramp
By F. Chalot, T.J.R. Hughes, Z. Johan, F. Shakib
(With 20 Figures) 181

The Application of an Adaptive Upwind Unstructured Grid
Solution Algorithm to the Simulation of Compressible Laminar
Viscous Flows over Compression Corners
By M. Vahdati, K. Morgan, J. Peraire (With 13 Figures) 201

Computation of Flows over 2D Ramps
By R. Radespiel, U. Herrmann, J.M.A. Longo (With 18 Figures) 212

Hypersonic Viscous Flow over Two-Dimensional Ramps
By Dachun Jiang, B.E. Richards (With 20 Figures) 228

Grid-Refinement Study of Hypersonic Laminar Flow over a 2-D Ramp By J.L. Thomas, D.H. Rudy, A. Kumar, B. van Leer (With 4 Figures)	244
Contribution to Problem 3 Using a Galerkin Least Square Finite Element Method By M. Mallet, B. Mantel, J. Périaux, B. Stoufflet (With 12 Figures)	255
Computational Results for Flows over Compression Ramps By W. Haase (With 20 Figures)	268
Implicit Upwind Finite-Difference Simulation of Laminar Hypersonic Flow over a 2D Ramp By B. Müller (With 20 Figures)	285
Synthesis	
A Synthesis of Results on the Calculation of Flow over a 2D Ramp and a 3D Obstacle: Antibes Test Cases 3 and 4 By J.F. Wendt, M. Mallet, B. Oskam (With 8 Figures)	301
Problem 5	
Experimental Study of the Longitudinal Hypersonic Corner Flow Field* By A. Henckels, F. Maurer (With 14 Figures)	315
Problem 6	
6.1 Non-reacting Flows	
Experimental Study of the Flow Around a Double Ellipsoid Configuration* By D. Aymer, T. Alziary, L. de Luca, G. Carlomagno (With 22 Figures)	335
Solution of the Euler Equations Around a Double Ellipsoidal Shape Using Unstructured Meshes and Including Real Gas Effects By F. Dubois, O. Michaux (With 18 Figures)	358

Computation of 2D Inviscid Hypersonic Flows Using Unstructured Polygonal Meshes By P. Vankeirsbilck, H. Deconinck (With 17 Figures)	374
An Upwind Relaxation Method for Hypersonic Viscous Flows over a Double-Ellipsoidal Body By R. Schwane, D. Hänel (With 20 Figures)	396
Navier-Stokes Calculations over a Double Ellipse and a Double Ellipsoid by an Implicit Non-centered Method By C. Marmignon, H. Hollanders, F. Coquel (With 18 Figures)	414
Application of the Galerkin/Least-Squares Formulation to the Analysis of Hypersonic Flows: II. Flow Past a Double Ellipse By F. Chalot, T.J.R. Hughes, Z. Johan, F. Shakib (With 25 Figures)	427
The Application of an Adaptive Unstructured Grid Method to the Solution of Hypersonic Flows Past Double Ellipse and Double Ellipsoid Configurations By O. Hassan, J. Peiro, J. Peraire, K. Morgan (With 20 Figures)	451
Computation of the Hypersonic Flow over a Double Ellipsoid By R. Radespiel, U. Herrmann, J.M.A. Longo (With 18 Figures)	472
High Resolution Schemes for Steady Hypersonic Flow By Z.J. Wang, B.E. Richards (With 20 Figures)	494
Numerical Simulation of Laminar Hypersonic Flow Past a Double-Ellipsoid By S. Riedelbauch (With 20 Figures)	517
Inviscid and Viscous Flow Calculation over a Double Ellipse with an Implicit Centered Method By K. Khalfallah, G. Lacombe, A. Lerat (With 16 Figures)	535
2D Hypersonic Viscous Flow Past a Double Ellipse Geometry By P. Leyland (With 12 Figures)	553
Hypersonic Flows over a Double or Simple Ellipse By F. Angrand, B. Tessieras, B. Dubroca, J.P. Morreeuw (With 16 Figures)	566

Solving Flow Equations for High Mach Numbers on Overlapping Grids By B. Gustafsson, E. Pärt-Enander, B. Sjögren (With 19 Figures)	585
An Explicit Finite-Difference Solution of Hypersonic Flows Using Rational Runge-Kutta Scheme By N. Satofuka, K. Morinishi (With 18 Figures)	600
Viscous and Inviscid Hypersonic Flow About a Double Ellipsoid By K. Dortmann (With 19 Figures)	616
Numerical Simulation of Hypersonic Flow over a Double Ellipse Using a Taylor-Galerkin Finite Element Formulation with Adaptive Grids By E. Oñate, F. Quintana, J. Miquel (With 12 Figures)	635
Hypersonic Viscous Flow Past Double Ellipse and Past Double Ellipsoid – Numerical Results By J. Argyris, I. St. Doltsinis, H. Friz (With 18 Figures)	655
Adaptive Mesh Embedding for Reentry Flow Problems By F. Grasso, M. Passalacqua (With 16 Figures)	673
Synthesis	
Attempt to Evaluate the Computation for Test Case 6.1 – Cold Hypersonic Flow Past Ellipsoidal Shapes By W. Kordulla, J. Périaux, T. Alziary de Roquefort (With 11 Figures)	689
6.2/6.3 Reacting Flows	
A Staggered Mesh Finite Difference Scheme for the Computation of Hypersonic Euler Flows By R. Sanders (With 4 Figures)	713
Reactive Flow Computations by Upwind Finite Elements By N. Botta, M.-C. Ciccoli, J.-A. Désidéri, L. Fezoui, N. Glinisky, E. Hettena, C. Olivier (With 22 Figures)	731
Numerical Analysis of Chemically Reacting Inviscid Flow in 2-D By M. Fey, R. Jeltsch, P. Karmann (With 18 Figures)	749

A Contribution to the Prediction of Hypersonic Non-equilibrium Flows By S. Borrelli, M. Pandolfi (With 16 Figures)	765
Reactive and Inert Inviscid Flow Solutions by Quasi-linear Formulations and Shock Fitting By F. Sabetta, B. Favini, G. Moretti, M. Onofri, M. Valorani (With 20 Figures)	782
Inviscid Calculations by an Upwind Finite Element Method of Hypersonic Flows over a Double (Single) Ellipse By V. Selmin, L. Formaggia (With 21 Figures)	798
Contribution to Problem 6 Using an Upwind Euler Solver with Unstructured Meshes By A. Descamps, M. Mallet, J. Périaux, B. Stoufflet (With 19 Figures)	815
Inviscid Hypersonic Flow Simulations Using an Explicit Scheme By J.B. Vos, C.M. Bergman (With 16 Figures)	832
Computation of Thermochemical Non-equilibrium Flows Around a Simple and a Double Ellipse By T. Gökçen (With 18 Figures)	848
Synthesis	
Some Comments on the Numerical Computations of Reacting Flows over the Double-Ellipse (Double Ellipsoid) By J.-A. Désidéri (With 12 Figures)	871
6.4 Rarefied Flows	
DSMC Calculations for the Double Ellipse By J.N. Moss, J.M. Price, M.C. Celenligil (With 13 Figures) . . .	882
Rarefied Flow Around a Double Ellipse By S. Igarashi, K. Nanbu (With 18 Figures)	896
The Hypersonic Double Ellipse in Rarefied Flow By W.J. Feiereisen (With 7 Figures)	912

Problem 7

7.1 Vortex Flows

Leeside Flow over Delta Wing at $M = 7.15$ Experimental Results for Test Case 7.1.2* By M. Linde	927
Hypersonic Delta Wing Flow Calculations Using a Multidomain MUSCL Euler Solver By Ph. Guillen, M. Borrel (With 18 Figures)	933
Finite Volume 3DNS and PNS Solutions of Hypersonic Viscous Flow Around a Delta Wing Using Osher's Flux Difference Splitting By N. Qin, B.E. Richards (With 9 Figures)	947
Inviscid Hypersonic Flow over a Delta Wing By S.M. Hitzel (With 20 Figures)	960
Hypersonic Leeside Delta-Wing-Flow Computations Using Centered Schemes By P. Eliasson, A. Rizzi, S. Srinivasan, B. Winzell (With 21 Figures)	981

Synthesis

Evaluation of Contributions for Test Cases 7.1.1 and 7.1.2 By A. Dervieux, M. Linde, A. Rizzi (With 10 Figures)	1006
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7.2 Rarefied Flows

Experimental Density Flowfields over a Delta Wing Located in Rarefied Hypersonic Flows* By J. Allegre, X. Heriard Dubreuilh, M. Raffin (With 16 Figures)	1017
Experiments on the Heat Transfer and on the Aerodynamic Coefficients of a Delta Wing in Rarefied Hypersonic Flows* By Ch.-H. Chun (With 17 Figures)	1031

DSMC Calculations for the Delta Wing By M.C. Celenligil, J.N. Moss (With 12 Figures)	1051
---	------

Rarefied Gas Flow Around a 3D-Deltawing By F. Gropengiesser, H. Neunzert, J. Struckmeier, B. Wiesen (With 14 Figures)	1067
---	------

Synthesis (Problems 6.4/7.2)

Appraisal of the Rarefied Flow Computations By J.K. Harvey (With 1 Figure)	1070
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Problem 8

Study of a Nitrogen D.C. Plasma Torch at Different Pressures Using Optical and Thermal Diagnostics* By J.M. Léger, J.F. Coudert, A. Grimaud, P. Fauchais (With 11 Figures)	1089
---	------

Non-equilibrium Flow in an Arc-Jet By P. Rostand, R.W. MacCormack (With 13 Figures)	1102
--	------

Non-equilibrium Flow in an Arc Heated Wind Tunnel By P.C. Sleziona, M. Auweter-Kurtz, B. Glocker, T. Gogel, T. Gölz, E. Messerschmid, H.O. Schrade (With 15 Figures) . . .	1116
--	------

Calculation of Non-equilibrium Flows in a High-Enthalpy Wind Tunnel By R. Abgrall (With 19 Figures)	1131
---	------

Application of Program LAURA to Thermochemical Non-equilibrium Flow Through a Nozzle By P.A. Gnoffo (With 11 Figures)	1145
---	------

Non-equilibrium Flow in a Hypersonic Wind Tunnel Nozzle By S.P. Spekreijse (With 23 Figures)	1159
---	------

Non-equilibrium Flow in a Hypersonic Wind Tunnel Nozzle By D. Zeitoun, P. Colas, M. Imbert, R. Brun (With 18 Figures) . . .	1174
--	------

Quasi Monodimensional Inviscid Non-equilibrium Nozzle Flow Computation By L. Marraffa (With 12 Figures)	1193
---	------

Synthesis**Non-equilibrium Flow in an Arcjet or a Shock Tube****By J. Muylaert, B. Stoufflet (With 11 Figures) 1206****Conclusion****General Synthesis****By P. Perrier 1221****Index of Contributors 1229**

Problems for Analysis

In the perspective of European and American space projects, the knowledge of flows at high Mach numbers and large angles of attack, occurring during the atmospheric reentry of space vehicles, has regained recently considerable interest. The extreme physical conditions realized during this critical phase are at the limit of possible experimental measurements, because of the rarefied atmosphere and/or the very high temperatures. The experimentation is therefore very difficult (if available) and costly. In addition, safety requirements cannot be achieved by over-dimensioning the structures without compromising the operational performance. Hence, the numerical simulation is led to play a major role in the process of the project validation.

During reentry, three main types of flow are encountered:

- rarefied gas flow,
- reacting flow,
- non-reacting flow.

The governing equations to be solved are then respectively - the Boltzmann equations, - the Euler or Navier-Stokes equations combined with a chemistry model, - the Navier-Stokes equations possibly combined with a turbulence model.

The following eight critical problems have been proposed to the participants:

Problem 1: Flow over a Slender Cone

Problem 2: Base Flow

Problem 3: Flow over a 2D Ramp

Problem 4: Flow over a 3-D Obstacle

Problem 5: Corner Flow

Problem 6: Flow over a Double Ellipse/ Double Ellipsoid

Problem 7: Flow over a Delta Wing

Problem 8: Non-equilibrium Flow in an Arc Jet or a Shock Tube

Due to the interest or the difficulty of each test case, a level of gradation, ranging from no "star"(0) to four "star" (****) has been attributed. The mandatory condition for participation is to collect a sum of at least four stars, by any combination of test cases solved (for example four cases with one star (*), or two cases with two stars (**), etc...).

Fluid Characteristics and Notations

• Fluid characteristics

1. Non-reactive flows

- The fluid is a perfect gas with specific heats c_p , c_v with ratio

$$\gamma = \frac{c_p}{c_v} = 1.4$$

- If p is the static pressure, the pitot pressure p_t is defined by :

$$\frac{p_t}{p} = \left(\frac{(\gamma + 1)M^2}{2} \right)^{\gamma/(\gamma-1)} \left(\frac{\gamma + 1}{2\gamma M^2 - (\gamma - 1)} \right)^{1/(\gamma-1)}$$

- The viscosity coefficients λ , μ verify the Stokes relation:

$$3\lambda + 2\mu = 0$$

- The Prandtl number is constant: $Pr = \mu c_p / \kappa = .72$ where κ is the thermal conductivity coefficient.
- The viscosity coefficient μ is a function of temperature T in Kelvin according to Sutherland's law:

- If $T_\infty \geq 120K$:

$$\begin{aligned} \mu(T) &= \mu_\infty \left(\frac{T}{T_\infty} \right)^{1.5} \left(\frac{T_\infty + 110.}{T + 110.} \right) & \text{if } T \geq 120K \\ \mu(T) &= \mu(120) \frac{T}{120} & \text{if } T \leq 120K \end{aligned}$$

- If $T_\infty \leq 120K$:

$$\begin{aligned} \mu(T) &= \mu_\infty \frac{T}{T_\infty} & \text{if } T \leq 120K \\ \mu(T) &= \mu(120) \left(\frac{T}{120} \right)^{1.5} \left(\frac{120. + 110.}{T + 110.} \right) & \text{if } T \geq 120K \end{aligned}$$

μ_∞ is obtained from the Reynolds number : $Re_\infty / m = u_\infty \rho_\infty / \mu_\infty$, a value of temperature in Kelvin is required to adjust the constant in Sutherland's law (T_∞ in Kelvin).

- When a turbulent simulation is required, the choice of the model is completely free.

2. Non-equilibrium reactive flows

- A computation with Park's model is required. This model is described in Appendix 1 where the unit for the rate constants is $cm^3/mole/s$. In addition, participants are encouraged to also use more accurate physical models to assess the influence of the model.
- The body surface is supposed to be fully non-catalytic in all test cases.

- Contributors should be aware that Sutherland's law is not physically relevant and alternate models should be used.

3. Equilibrium reactive flows

- Similar models as for non-equilibrium reactive flows but with the assumption of local equilibrium.

4. Rarefied gases

- The internal energy model to be considered is the Larsen - Borgnakke model. Calculations should be performed for a monoatomic gas for comparison or if the participants have not implemented the Larsen - Borgnakke model.
- The fluid is characterized by the mean free path at infinity λ_∞ .

• Notations

The following non dimensional coefficients are defined:

- Pressure coefficient: $C_p = \frac{p - p_\infty}{0.5\rho_\infty u_\infty^2}$ where ρ_∞ and u_∞ are the density and the velocity at infinity.
- Skin friction coefficient: $C_f = \frac{\tau_w}{0.5\rho_\infty u_\infty^2}$ (τ_w : wall shear stress)
- Heat flux coefficient: $C_h = \frac{q}{0.5\rho_\infty u_\infty^3}$ where q is the heat flux; in the continuum regime: $q = -\kappa \nabla T \cdot n$ where n is the outward normal to the wall.
- Stanton number : $St = \frac{q}{\rho_\infty u_\infty c_{p\infty}(T_{0_\infty} - T_w)}$ where $c_{p\infty}$ is the free stream specific heat at constant pressure, T_{0_∞} is the reservoir isentropic stagnation temperature and T_w is the wall temperature.
- All dimensions are given in the meter, kilogram, second (MKS) system.