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tables of dimensions,
indices, and
branching rules
for representations
of simple lie algebras

W. G. McKay
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Tables of Dimensions, Indices, and Branching Rules for Representations of Simple Lie Algebras

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Preface

The nature of applications of semisimple Lie algebras and groups frequently requires the knowledge of how representations of an algebra reduce to representations of its subalgebras. The problem, although fairly standard from the point of view of mathematics, has no known easy general solution. There arises therefore a need to compile the results, within certain practical limits, in the form of tables. This volume contains the reduction (branching rules) of representations of complex simple Lie algebras of ranks not exceeding 8 to representations of their maximal semisimple subalgebras; the dimensions of representations given here are limited by 5000 for the classical Lie algebras, and 10000 for the five exceptional algebras. Also, there are tables of auxiliary quantities (dimensions, second and fourth indices of representations) useful in applications.

In 1971 a book of similar nature but much more limited content (dimensions < 1000) was prepared [1]. Since then the computation procedure [2] has been considerably speeded up. This volume contains results of an independent computation of branching rules; some conventions were modified in comparison with Ref. 1 and a few corrections were made. The text is completely rewritten.

We make no attempt to collect a complete bibliography on branching rules, which have been the subject of a number of publications during the last fifteen years. Efficient methods are known for particular algebra-subalgebra pairs [3] and extensive tables have been computed for some

others often used in spectroscopy [4].

The present book is intended for quick reference by people knowing basic facts about semisimple Lie algebras and their finite dimensional representations [5] who use representations of semisimple Lie algebras as tools in their own work.

Since the publication of Ref. 1 we have received many comments and helpful suggestions. We are particularly grateful to Dr. F.W. Lemire, Dr. R.T. Sharp, and Dr. R. Slansky.

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1/Notations, Conventions, and Some Properties of Representations

Semisimple Lie Algebras

There are nine types of complex simple Lie algebras: four classical series A_n, B_n, C_n, D_n of ranks $n = 1, 2, \dots$, and five types of isolated exceptional algebras E_6, E_7, E_8, F_4 and G_2 . A semisimple Lie algebra L decomposes into simple algebras (ideals); we write

$$L = L_1 + L_2 + \dots + L_k \quad (1.1)$$

There exist only the following isomorphisms among semisimple Lie algebras:

$$A_1 \sim B_1 \sim C_1 \sim D_1, \quad B_2 \sim C_2, \quad D_2 \sim A_1 + A_2, \quad A_3 \sim D_3 \quad (1.2)$$

The symbols B_1, C_1, D_1, B_2, D_2 , and D_3 do not therefore appear further in this book.

With every semisimple Lie algebra of rank n one associates a system of n real vectors α_i , $i = 1, 2, \dots, n$, called simple roots. The system is conveniently presented in the form of a Dynkin diagram (Fig. 1). A simple root is denoted by a circle. Two circles which are connected by a single, double or triple line correspond to two roots spanning an angle of 120° , 135° , and 150° respectively; two circles which are not connected directly correspond to mutually orthogonal simple roots. The Dynkin diagram of a

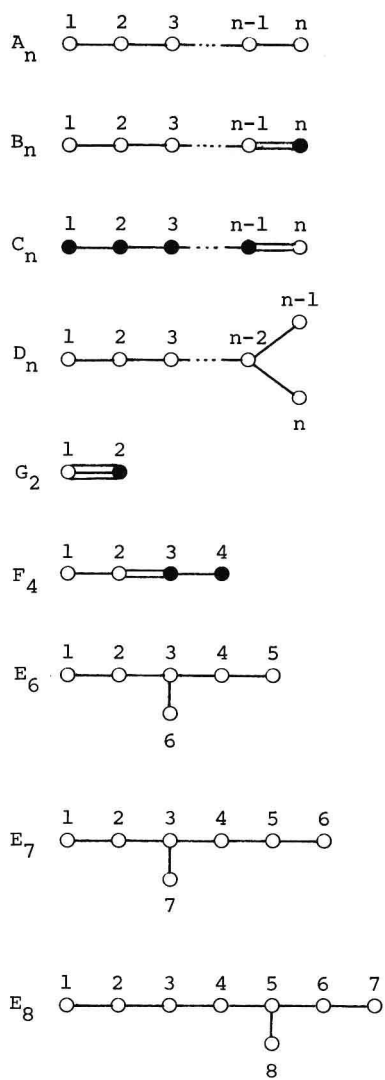


FIG. 1. Numbering of simple roots of simple Lie algebras

simple algebra is connected, the diagram of a nonsimple algebra (1.1) consists of k disconnected subdiagrams. Each simple algebra has simple roots of no more than two different lengths, the black circles denoting the shorter roots. The length of the longer roots is fixed by the convention that the scalar product (α, α) , has value 2,

$$(\alpha, \alpha) = 2 \quad (1.3)$$

Furthermore, one has $(\alpha, \alpha) = 1$ for the shorter roots in the case of algebras B_n , C_n , and F_4 , and $(\alpha, \alpha) = 2/3$ for G_2 .

Throughout the book the simple roots are numbered as on Fig. 1.

Contragredient, Orthogonal, and Symplectic Representations

An irreducible representation $\phi(L)$ of a simple Lie algebra L is often denoted by its highest weight Λ . The highest weight Λ of a particular irreducible representation is specified by the coordinates

$$a_i = \frac{2(\Lambda, \alpha_i)}{(\alpha_i, \alpha_i)}, \quad i = 1, 2, \dots, n \quad (1.4)$$

Thus one writes $\Lambda = (a_1 a_2 \dots a_n)$. In particular, the convention (1.3) implies that an irreducible representation of A_1 (the algebra of "angular momentum, spin, isospin", etc. of quantum mechanics) is given by (a_1) , where a_1 is an integer equal to *twice* the corresponding angular momentum.

An irreducible representation of a nonsimple algebra, say $L_1 + L_2$, is specified by

$$\phi(L_1 + L_2) = (a_1 a_2 \dots a_n) (b_1 b_2 \dots b_m) \quad (1.5)$$

where $(a_1 a_2 \dots a_n)$ and $(b_1 b_2 \dots b_m)$ are representations of L_1 and L_2 , respectively. In order to improve the readability of the Table 2 below, we use there $(a_1 a_2 \dots a_n - b_1 b_2 \dots b_m)$ instead of (1.5). Sometimes it is advantageous to give a representation by the Dynkin diagram of the corresponding Lie algebra with the coordinates (1.4) shown at each simple root.

Some representations are of special interest. We list them here. The

nontrivial representations of the lowest dimension are as follows:

$$\begin{array}{llll}
 (10\dots 0) \text{ and } (0\dots 01) & \text{for } A_n, & n \geq 1 \\
 (10\dots 0) & \text{for } B_n, & n \geq 3 \\
 (10\dots 0) & \text{for } C_n, & n \geq 2 \\
 (1000), (0010), \text{ and } (0001) & \text{for } D_4 \\
 (10\dots 0) & \text{for } D_n, & n \geq 5 \\
 (100000) \text{ and } (000010) & \text{for } E_6 \\
 (0000010) & \text{for } E_7 \\
 (1000000) & \text{for } E_8 \\
 (0001) & \text{for } F_4 \\
 (01) & \text{for } G_2
 \end{array} \tag{1.6}$$

The adjoint representations are

$$\begin{array}{llll}
 (2) & \text{for } A_1 \\
 (10\dots 01) & \text{for } A_n, & n \geq 2 \\
 (010\dots 0) & \text{for } B_n, & n \geq 3 \\
 (20\dots 0) & \text{for } C_n, & n \geq 2 \\
 (010\dots 0) & \text{for } D_n, & n \geq 4 \\
 (000001) & \text{for } E_6 \\
 (1000000) & \text{for } E_7 \\
 (10000000) & \text{for } E_8 \\
 (1000) & \text{for } F_4 \\
 (10) & \text{for } G_2
 \end{array} \tag{1.7}$$

A representation $\phi(L)$ acting in a space S_ϕ is selfcontragredient if there exists an invariant bilinear form (x,y) , where $x,y \in S_\phi$, such that

$$(Xx,y) + (x,Xy) = 0, \quad X \in \phi(L) \tag{1.8}$$

A selfcontragredient representation is orthogonal or symplectic if the form is symmetric, $(x,y) = (y,x)$, or skew symmetric, $(x,y) = -(y,x)$. The simple Lie algebras of types

$$A_1, B_n, C_n, D_{2k}, E_7, E_8, F_4, \text{ and } G_2 \tag{1.9}$$

have only selfcontragredient representations. Among the irreducible representations of the remaining simple Lie algebras only the following are selfcontragredient

$$\begin{aligned}
 (a_1 a_2 \dots a_2 a_1) &\equiv \begin{array}{c} a_1 \quad a_2 \quad \dots \quad a_2 \quad a_1 \\ \circ \text{---} \circ \text{---} \dots \text{---} \circ \text{---} \circ \end{array} \\
 (a_1 a_2 \dots a_{2k-1} a_{2k} a_{2k}) &\equiv \begin{array}{c} a_1 \quad a_2 \quad \dots \quad a_{2k-1} \quad \circ \\ \circ \text{---} \circ \text{---} \dots \text{---} \circ \swarrow \searrow \\ \quad \quad \quad \quad \quad \quad \quad \quad \circ \quad a_{2k} \\ \quad \quad \quad \quad \quad \quad \quad \quad \circ \quad a_{2k} \end{array} \\
 (a_1 a_2 a_3 a_2 a_1 a_6) &\equiv \begin{array}{c} a_1 \quad a_2 \quad a_3 \quad a_2 \quad a_1 \\ \circ \text{---} \circ \text{---} \circ \text{---} \circ \text{---} \circ \\ \quad \quad \quad \quad \quad \quad \quad \quad \circ \\ \quad \quad \quad \quad \quad \quad \quad \quad a_6 \end{array}
 \end{aligned} \tag{1.10}$$

Here on the right hand side the representations are denoted using the Dynkin diagram which specifies the type of the algebra.

The irreducible selfcontragredient representations of

$$A_{2k}, A_{4k+3}, B_{4k}, B_{4k+3}, D_{2k+1}, D_{4k}, E_6, E_8, F_4, \text{ and } G_2 \tag{1.11}$$

are all orthogonal. An irreducible selfcontragredient representation of any of the remaining simple Lie algebras is orthogonal (symplectic) if a certain linear combination of the coordinates (1.4) of the highest weight is even (odd). These linear combinations are

$$\begin{aligned}
 a_{2k+1} &\quad \text{for } A_{4k+1}, \quad k \geq 1 \\
 a_{4k+1} &\quad \text{for } B_{4k+1}, \quad k \geq 1 \\
 a_{4k+2} &\quad \text{for } B_{4k+2}, \quad k \geq 1 \\
 a_1 + a_3 + a_5 + \dots &\quad \text{for } C_n, \quad n \geq 2 \\
 a_{4k+1} + a_{4k+2} &\quad \text{for } D_{4k+2}, \quad k \geq 1 \\
 a_4 + a_6 + a_7 &\quad \text{for } E_7
 \end{aligned} \tag{1.12}$$

Two irreducible representations $\phi(L)$ and $\bar{\phi}(L)$ of a simple Lie algebra

L are mutually contragredient if they are one of the pairs

$$\phi(A_n) = (a_1 a_2 \dots a_n) \quad \text{and} \quad \bar{\phi}(A_n) = (a_n a_{n-1} \dots a_2 a_1)$$

$$\begin{aligned} \phi(D_{2k+1}) &= (a_1 a_2 \dots a_{2k-1} a_{2k} a_{2k+1}) \quad \text{and} \\ \bar{\phi}(D_{2k+1}) &= (a_1 a_2 \dots a_{2k-1} a_{2k+1} a_{2k}) \end{aligned} \quad (1.13)$$

$$\phi(E_6) = (a_1 a_2 a_3 a_4 a_5 a_6) \quad \text{and} \quad \bar{\phi}(E_6) = (a_5 a_4 a_3 a_2 a_1 a_6)$$

A direct sum of selfcontragredient, orthogonal, symplectic representations is respectively selfcontragredient, orthogonal, or symplectic.

The direct sum $\phi(L) \oplus \bar{\phi}(L)$ of two mutually contragredient representations is both orthogonal and symplectic, i.e., there exists a symmetric as well as skew symmetric invariant bilinear form in the corresponding representation space $S_{\phi \oplus \bar{\phi}}$.

An irreducible representation

$$(a_1^1 a_2^1 \dots a_{n_1}^1) (a_1^2 a_2^2 \dots a_{n_2}^2) \dots (a_1^k a_2^k \dots a_{n_k}^k) \quad (1.14)$$

of a nonsimple algebra (1.1) is selfcontragredient if all the representations

$$(a_1^i a_2^i \dots a_{n_i}^i), \quad i = 1, 2, \dots, k \quad (1.15)$$

are selfcontragredient. A selfcontragredient representation (1.14) is orthogonal (symplectic) if it contains an even (odd) number of symplectic components (1.15).

Congruence Classes of Representations

Another convenient characterization of an irreducible representation of a semisimple algebra L is by its congruence class [6]. A particular case

of that property well known in particle physics is the triality of SU(3) representations.

We say that two irreducible representations Λ and Λ' of a semisimple Lie algebra L belong to the same congruence class or are mutually congruent if and only if the difference $\Lambda - \Lambda'$ of their highest weights is an integer linear combination of simple roots of L . Each congruence class is characterized by a value of the congruence number c which, for an irreducible representation $(a_1 a_2 \dots a_n)$ of a simple Lie algebra, is given by

$$c = \sum_{k=1}^n a_k \pmod{n+1} \quad \text{for } A_n \quad (1.16)$$

$$c = a_n \pmod{2} \quad \text{for } B_n \quad (1.17)$$

$$c = a_1 + a_3 + a_5 + \dots \pmod{2} \quad \text{for } C_n \quad (1.18)$$

$$c = a_1 - a_2 + a_4 - a_5 \pmod{3} \quad \text{for } E_6 \quad (1.19)$$

$$c = a_4 + a_6 + a_7 \pmod{2} \quad \text{for } E_7 \quad (1.20)$$

$$c = 0 \quad \text{for all representations of } E_8, F_4, \text{ and } G_2 \quad (1.21)$$

Finally, when L is of the type D_n the congruence number c is a two component vector

$$c = (a_{n-1} + a_n, 2a_1 + 2a_3 + \dots + 2a_{n-2} + (n-2)a_{n-1} + na_n) \quad \text{for } n \text{ odd} \quad (1.22)$$

$$c = (a_{n-1} + a_n, 2a_1 + 2a_3 + \dots + 2a_{n-3} + (n-2)a_{n-1} + na_n) \quad \text{for } n \text{ even} \quad (1.23)$$

In (1.22) and (1.23) the first component is calculated mod 2 and the second one mod 4. One easily verifies that there are four congruence