# Langmuir-Blodgett filmes An introduction

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### An introduction

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## **Preface**

Langmuir-Blodgett (LB) films have been the subject of scientific curiosity for most of the twentieth century. However, interest has grown significantly since the 1970s – a direct result of the work of Hans Kuhn and colleagues on energy transfer in multilayer systems. This introduced the idea of molecular engineering, i.e., using the LB technique to position certain molecular groups at precise distances to others. In this way new thin film materials could be built up at the molecular level and the relationship between these artificial structures and the natural world explored.

There are already several books that cover LB and related thin films. So why another? My own background is in electronics. While I have been involved in LB film research I have spent many hours pondering on chemical formulae, struggling with biological nomenclature and trying to understand the finer points of thermodynamics. The scope of the subject is continuing to grow and anyone now starting work in the area must assimilate an enormous amount of information. My intention therefore has been to provide a gentle introduction to newcomers with an emphasis on the multidisciplinary and interdisciplinary nature of the field.

Each chapter addresses a different issue. Chapter 1 describes the various bulk phases of matter and outlines physical principles that can be used to model these. Monolayer phases are introduced in chapter 2. Much is still to be learnt about the nature of even the simplest monolayers and I have concentrated on experimental aspects rather than theory. The LB process itself is discussed in some detail in chapter 3. Again, my aim has been to provide useful practical information. I have also briefly reviewed related thin film techniques. Materials that can be manipulated at the air/water interface are described in chapter 4. In chapter 5 the main methods of investigation are introduced and some important results outlined. Chapter 6 covers the

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electrical behaviour. This is an area that has attracted considerable interest, not least because of the possibilities of exploiting LB layers in electronic device structures; I have included a section outlining such ideas. Finally, chapter 7 provides an overview of the optical properties of mono- and multilayer films.

At the end of each chapter, I have given the key references. These should be good starting points for the reader wishing to delve deeper. Important subjects such as crystallography, chemical bonding and the interaction of electromagnetic radiation with thin films are dealt with in separate appendices.

I hope that the result is a book readily accessible to both new research students and more experienced workers.

# Acknowledgements

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Original photographs were provided by Jas Pal Badyal and Simon Evenson, David Batchelder, Harald Fuchs and Mathias Lösche. Dirk Hönig spent many hours working on the cover photograph. Julie Morgan helped with the rest of the photography. Kay Seaton put in a tremendous effort drawing all the diagrams and deserves my very special thanks.

## Notes for reference

#### Constants

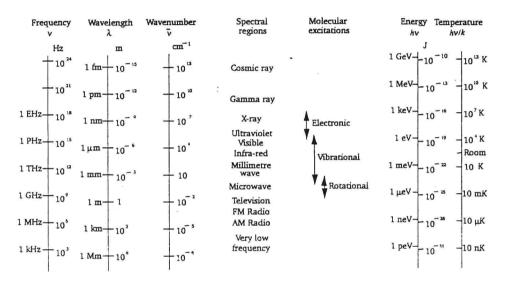
velocity of light in free space c permittivity of free space,  $\epsilon_0$  electronic charge, q Planck's constant, h Boltzmann's constant, k Avogadro's number,  $N_A$  universal gas constant, R

 $2.998 \times 10^{8} \text{ m s}^{-1}$   $8.854 \times 10^{-12} \text{ F m}^{-1}$   $1.602 \times 10^{-19} \text{ C}$   $6.626 \times 10^{-34} \text{ J s}$   $1.381 \times 10^{-23} \text{ J K}^{-1}$   $6.022 \times 10^{26} \text{ kilomole}^{-1}$  $8.314 \times 10^{3} \text{ J kilomole}^{-1} \text{ K}^{-1}$ 

#### Useful relationships

```
1 electronvolt (eV) = 1.602 \times 10^{-19} J for vacuum, energy in eV = 1.240/(wavelength in \mu m) 1 eV = 8066 cm<sup>-1</sup> 1 eV per particle = 23\,060 kcal kilomole<sup>-1</sup> (=23.06 kcal mole<sup>-1</sup>) 1 calorie = 4.186 J at 300 K, kT \approx 1/40 eV 1 atmosphere = 1.013 \times 10^5 N m<sup>-2</sup>
```

#### The electromagnetic spectrum



 $[1Hz \equiv 3.00 \times 10^8 \text{ m} \equiv 4.14 \times 10^{-15} \text{ eV} \equiv 4.80 \times 10^{-11} \text{ K}]$ 

#### Simple functional groups —

Compound type	Functional group
Alkene (double bond)	)c=c(
Alkene (triple bond)	-c≡c-
Alkene (aromatic ring)	
Halide	C - X (X - F, Cl, Br, I)
Alcohol	
Ether	-\c-o-c <del>\</del>
Carbonyl	)c=o
Aldehyde	H_c=0
Ketone	co=o
Ester	_c_o_c_

Compound type	Functional group
Carboxylic acid	О    -С-О-Н
Carboxylic acid chloride	_c_a
Amide	_C_n′
Amine	\c-n'\
Nitrîle	-с≡и
Nitro	-)c-n'0
Sulfide	-\_c-s-c\_
Sulfoxide	
Sulfone	-\c-\s^2+c\/
Organometallic	

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# The bulk phases of matter

#### 1.1 Gases, liquids and solids

The three most common states, or phases, of matter, gases, liquids and solids are very familiar (Walton, 1976). Phases that are not so well known are plasmas and liquid crystals (although these are both found in electrical and electronic devices in everyday use). All these states are generally distinguished by the degree of translational and orientational order of the constituent molecules. On this basis some phases may be further subdivided. For example, solids, consisting of a rigid arrangement of molecules, can be crystalline or amorphous. In an amorphous solid (a good example is a glass), the molecules are fixed in place, but with no pattern in their arrangement. As shown in figure 1.1, the crystalline solid state is characterized by long-range translational order of the constituent molecules (the molecules are constrained to occupy specific positions in space) and long-range orientational order (the molecules orient themselves with respect to each other). The molecules are, of course, in a constant state of thermal agitation, with a mean translational kinetic energy of 3kT/2 (k is Boltzmann's constant, T is temperature; kT/2 for each component of their velocity). However, this energy is considerably less than that associated with the chemical bonds in the material and the motion does not disrupt the highly ordered molecular arrangement.

In the gaseous state, the intermolecular forces are not strong enough to hold the molecules together. They are therefore free to diffuse about randomly, spreading evenly throughout any container they occupy, no matter how large it is. The average interatomic distance is determined by the number of molecules and the size of the container. A gas is easily compressed, as it takes comparatively little force to move the molecules closer together.

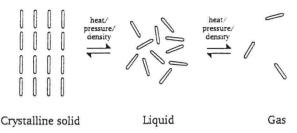


Figure 1.1 The three most common bulk phases of matter.

On the microscopic level, the liquid state is generally thought of as a phase that is somewhere between that of a solid and that of a gas. The molecules in a liquid neither occupy a specific average position nor remain oriented in a particular way. They are free to move around and, as in the gaseous state, this motion is random. The physical properties of both liquids and gases are *isotropic*; i.e., they do not depend upon direction (direction-dependent properties are called *anisotropic*). Therefore, there is a similarity between liquids and gases and under certain conditions it is impossible to distinguish between these two states. When placed in a container, the liquid will fill it to the level of a free surface. Liquids flow and change their shape in response to weak outside forces. The forces holding a liquid together are much less than those in a solid. Liquids are highly incompressible, a characteristic that is exploited in hydraulic systems.

Almost all elements and chemical compounds possess a solid, a liquid and a gaseous phase. The gaseous phase is favoured by high temperatures and low pressures. A transition from one phase to another can be provoked by a change in temperature, pressure, density or volume.

#### 1.2 Liquid crystals

The melting of a solid is normally a very sharp transition. As the material is heated, the order of the system abruptly changes from that of the three-dimensional order of the solid to the zero order associated with the liquid. However, it is not uncommon for an *organic* solid to pass through intermediate phases as it is heated from a solid to a liquid. Perhaps one in every few hundred organic compounds exhibits such behaviour. Such phases are known as *mesophases*. When such a state is formed, the translational order of the solid phase may be lost, but the orientational order remains. The resulting phase is termed a *thermotropic liquid crystal* (Collings, 1990). Sometimes a mesophase may display translational order but no orientational order; this is termed a *plastic crystal*. A very wide range of liquid crystal phases is now known. These are identified by the degree of short range translational order and by the shape of the molecules.

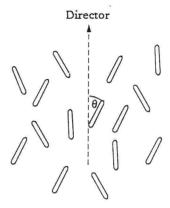


Figure 1.2 Arrangement of rod-shaped molecules in a liquid crystalline phase. Each molecule makes an angle of  $\theta$  with the director.

The *nematic* (from the Greek word for thread) phase is the least ordered liquid crystal phase and is exploited extensively in electrooptic applications. This phase has no long-range translational order and only orientational order. There is only one nematic phase and, on heating, this will eventually become an isotropic liquid. A schematic diagram, showing the arrangement of rod-shaped molecules in a nematic phase, is given in figure 1.2. The molecules in the liquid crystal are free to move about in much the same fashion as in a liquid; as they do so they tend to remain orientated in a certain direction. The direction of preferred orientation is called the *director* of the liquid crystal. Each molecule is oriented at some angle  $\theta$  about the director. The degree of orientational order is given by the *order parameter S* which is defined

$$S = \frac{1}{2} \langle 3\cos^2 \theta - 1 \rangle \tag{1.1}$$

A value of S = 1 indicates perfect orientational order whereas no orientational order results in S = 0.

Smectic (from the Greek word for soap) phases are usually formed by thermotropic liquid crystals at lower temperature than the nematic phase. Besides the orientational order, smectic phases possess one-dimensional translational ordering into layers. The smectic phases can be further subdivided and, at present, twelve different types have been identified. These are designated  $S_A$ ,  $S_B$ , . . .  $S_K$ ; there are two smectic  $S_B$  phases – the crystal B and hexatic (for a discussion of this phase, see Tredgold 1994)  $S_B$  phases. Some of these mesophases (crystal B,  $S_E$ ,  $S_G$ ,  $S_H$  and  $S_I$ ) have very longrange correlation of position over many layers and are more similar to crystalline solids.

The S<sub>A</sub> phase is the least ordered of the thermotropic smectic phases. The molecules are arranged in disordered layers, each layer having a liquid-like

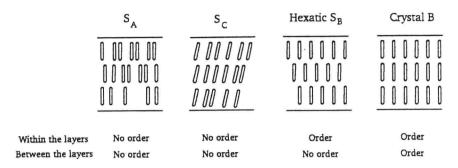


Figure 1.3 Different smectic liquid crystalline arrangements of rod-shaped molecules.  $S_C$  is a tilted version of  $S_A$ .

freedom of motion of its constituent molecules in two dimensions, with the director perpendicular to the layer planes. By contrast, in the  $S_C$  phase the molecules are tilted from this direction by about 35°. This tilt is correlated between molecules within each layer and from one layer to another. In the hexatic  $S_B$  phase, there is again ordering of orientationally aligned molecules into layers. The molecules are arranged in an hexagonal array but the translational order is short-range only. The orientation of the hexagonal net is, however, maintained over a long range and unlike the ordering of molecular positions, is correlated between layers. Two tilted variations of the  $S_B$  phase exist in which the tilt direction is constrained to point either towards one face of the hexagonal lattice ( $S_F$ ) or towards one apex ( $S_I$ ). The  $S_A$ ,  $S_C$ , hexatic  $S_B$  and crystal  $S_B$  arrangements of rod-shaped molecules are contrasted in figure 1.3 (Lacey, 1994). Further details about these smectic liquid crystal phases are to be found in the book by Gray and Goodby (1984).

The final distinct type of liquid crystalline mesophase is the *cholesteric* or *chiral nematic*. The molecules in such a phase are arranged in a unique helical structure in which the director gradually rotates from one plane of molecules to the next. Such phases can exhibit vivid colour effects that change with temperature (as the pitch of the helix changes) and are exploited as temperature sensors. These molecules may also form smectic phases.

#### 1.3 Phase changes

The science of thermodynamics allows the properties of the various phases of matter to be correlated. The *state* of any thermodynamic system (e.g., a fixed amount of matter) is specified by values of certain experimentally variable quantities called *state variables* or *properties*. Familiar examples are the temperature of a system, its pressure and the volume occupied by it. Those properties that are proportional to the mass of the system are called *extensive*; examples are the total volume and total energy of the