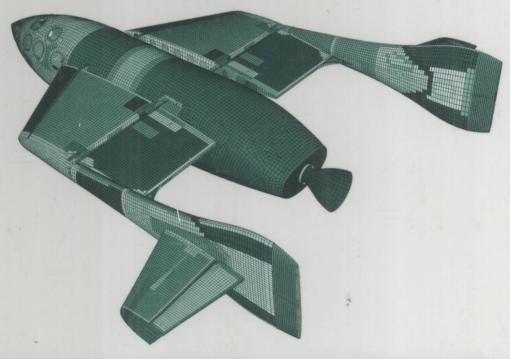
Applied Calculus of Variations for Engineers



Louis Komzsik



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To my daughter, Stella

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Louis Komzsik

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Preface

The topic of this book has a long history. Its fundamentals were laid down by icons of mathematics like Euler and Lagrange. It was once heralded as the panacea for all engineering optimization problems by suggesting that all one needs to do was to apply the Euler-Lagrange equation form and solve the resulting differential equation.

This, as most all encompassing solutions, turned out to be not always true and the resulting differential equations are not necessarily easy to solve. On the other hand, many of the differential equations commonly used by engineers today are derived from a variational problem. Hence, it is important and useful for engineers to delve into this topic.

The book is organized into two parts: theoretical foundation and engineering applications. The first part starts with the statement of the fundamental variational problem and its solution via the Euler-Lagrange equation. This is followed by the gradual extension to variational problems subject to constraints, containing functions of multiple variables and functionals with higher order derivatives. It continues with the inverse problem of variational calculus, when the origin is in the differential equation form and the corresponding variational problem is sought. The first part concludes with the direct solution techniques of variational problems, such as the Ritz, Galerkin and Kantorovich methods.

With the emphasis on applications, the second part starts with a detailed discussion of the geodesic concept of differential geometry and its extensions to higher order spaces. The computational geometry chapter covers the variational origin of natural splines and the variational formulation of B-splines under various constraints.

The final two chapters focus on analytic and computational mechanics. Topics of the first include the variational form and subsequent solution of several classical mechanical problems using Hamilton's principle. The last chapter discusses generalized coordinates and Lagrange's equations of motion. Some fundamental applications of elasticity, heat conduction and fluid mechanics as well as their computational technology conclude the book.

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Part I

Mathematical foundation



The foundations of calculus of variations

The problem of the calculus of variations evolves from the analysis of functions. In the analysis of functions the focus is on the relation between two sets of numbers, the independent (x) and the dependent (y) set. The function f creates a one-to-one correspondence between these two sets, denoted as

$$y = f(x)$$
.

The generalization of this concept is based on allowing the two sets not being restricted to be real numbers and to be functions themselves. The relationship between these sets is now called a functional. The topic of the calculus of variations is to find extrema of functionals, most commonly formulated in the form of an integral.

1.1 The fundamental problem and lemma of calculus of variations

The fundamental problem of the calculus of variations is to find the extremum (maximum or minimum) of the functional

$$I(y) = \int_{x_0}^{x_1} f(x, y, y') dx,$$

where the solution satisfies the boundary conditions

$$y(x_0) = y_0$$

and

$$y(x_1) = y_1.$$

This problem may be generalized to the cases when higher derivatives or multiple functions are given and will be discussed in Chapters 3 and 4, respectively. These problems may also be extended with constraints, the topic of Chapter 2.

A solution process may be arrived at with the following logic. Let us assume that there exists such a solution y(x) for the above problem that satisfies

the boundary conditions and produces the extremum of the functional. Furthermore, we assume that it is twice differentiable. In order to prove that this function results in an extremum, we need to prove that any alternative function does not attain the extremum.

We introduce an alternative solution function of the form

$$Y(x) = y(x) + \epsilon \eta(x),$$

where $\eta(x)$ is an arbitrary auxiliary function of x, that is also twice differentiable and vanishes at the boundary:

$$\eta(x_0) = \eta(x_1) = 0.$$

In consequence the following is also true:

$$Y(x_0) = y(x_0) = y_0$$

and

$$Y(x_1) = y(x_1) = y_1.$$

A typical relationship between these functions is shown in Figure 1.1 where the function is represented by the solid line and the alternative function by the dotted line. The dashed line represents the arbitrary auxiliary function.

Since the alternative function Y(x) also satisfies the boundary conditions of the functional, we may substitute into the variational problem.

$$I(\epsilon) = \int_{x_0}^{x_1} f(x, Y, Y') dx.$$

where

$$Y'(x) = y'(x) + \epsilon \eta'(x).$$

The new functional in terms of ϵ is identical with the original in the case when $\epsilon = 0$ and has its extremum when

$$\frac{\partial I(\epsilon)}{\partial \epsilon}|_{\epsilon=0} = 0.$$

Executing the derivation and taking the derivative into the integral, since the limits are fixed, with the chain rule we obtain

$$\frac{\partial I(\epsilon)}{\partial \epsilon} = \int_{x_0}^{x_1} \left(\frac{\partial f}{\partial Y} \frac{dY}{d\epsilon} + \frac{\partial f}{\partial Y'} \frac{dY'}{d\epsilon} \right) dx.$$

Clearly

$$\frac{dY}{d\epsilon} = \eta(x),$$

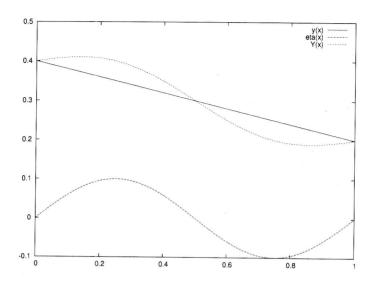


FIGURE 1.1 Alternative solutions example

and

$$\frac{dY'}{d\epsilon} = \eta'(x),$$

resulting in

$$\frac{\partial I(\epsilon)}{\partial \epsilon} = \int_{x_0}^{x_1} (\frac{\partial f}{\partial Y} \eta(x) + \frac{\partial f}{\partial Y'} \eta'(x)) dx.$$

Integrating the second term by parts yields

$$\int_{x_0}^{x_1} (\frac{\partial f}{\partial Y'} \eta'(x)) dx = \frac{\partial f}{\partial Y'} \eta(x)|_{x_0}^{x_1} - \int_{x_0}^{x_1} (\frac{d}{dx} \frac{\partial f}{\partial Y'}) \eta(x) dx.$$

Due to the boundary conditions, the first term vanishes. With substitution and factoring the auxiliary function, the problem becomes

$$\frac{\partial I(\epsilon)}{\partial \epsilon} = \int_{x_0}^{x_1} (\frac{\partial f}{\partial Y} - \frac{d}{dx} \frac{\partial f}{\partial Y'}) \eta(x) dx.$$

The extremum is achieved when $\epsilon = 0$ as stated above, hence

$$\frac{\partial I(\epsilon)}{\partial \epsilon}|_{\epsilon=0} = \int_{x_0}^{x_1} (\frac{\partial f}{\partial y} - \frac{d}{dx} \frac{\partial f}{\partial y'}) \eta(x) dx.$$

Let us now consider the following integral:

$$\int_{x_0}^{x_1} \eta(x) F(x) dx,$$

where $x_0 \le x \le x_1$ and F(x) is continuous, while $\eta(x)$ is continuously differentiable, satisfying

$$\eta(x_0) = \eta(x_1) = 0.$$

The fundamental lemma of calculus of variations states that if for all such $\eta(x)$

$$\int_{x_0}^{x_1} \eta(x) F(x) dx = 0,$$

then

$$F(x) = 0$$

in the whole interval.

The following proof by contradiction is from [13]. Let us assume that there exists at least one such location $x_0 \leq \zeta \leq x_1$ where F(x) is not zero, for example

$$F(\zeta) > 0.$$

By the condition of continuity of F(x) there must be a neighborhood of

$$\zeta - h \le \zeta \le \zeta + h$$

where F(x) > 0. In this case, however, the integral becomes

$$\int_{x_0}^{x_1} \eta(x) F(x) dx > 0,$$

for the right choice of $\eta(x)$, which contradicts the original assumption. Hence the statement of the lemma must be true.

Applying the lemma to this case results in the Euler-Lagrange differential equation specifying the extremum

$$\frac{\partial f}{\partial y} - \frac{d}{dx} \frac{\partial f}{\partial y'} = 0.$$

1.2 The Legendre test

The Euler-Lagrange differential equation just introduced represents a necessary, but not sufficient, condition for the solution of the fundamental variational problem.

The alternative functional of

$$I(\epsilon) = \int_{x_0}^{x_1} f(x, Y, Y') dx,$$

may be expanded as

$$I(\epsilon) = \int_{x_0}^{x_1} f(x, y + \epsilon \eta(x), y' + \epsilon \eta'(x)) dx.$$

Assuming that the f function has continuous partial derivatives, the mean-value theorem is applicable:

$$\begin{split} f(x,y+\epsilon\eta(x),y'+\epsilon\eta'(x)) &= f(x,y,y') + \\ \epsilon(\eta(x)\frac{\partial f(x,y,y')}{\partial y} + \eta'(x)\frac{\partial f(x,y,y')}{\partial y'}) + O(\epsilon^2). \end{split}$$

By substituting we obtain

$$\begin{split} I(\epsilon) &= \int_{x_0}^{x_1} f(x,y,y') dx + \\ \epsilon \int_{x_0}^{x_1} (\eta(x) \frac{\partial f(x,y,y')}{\partial y} + \eta'(x) \frac{\partial f(x,y,y')}{\partial y'}) dx + O(\epsilon^2). \end{split}$$

With the introduction of

$$\delta I_1 = \epsilon \int_{x_0}^{x_1} (\eta(x) \frac{\partial f(x, y, y')}{\partial y} + \eta'(x) \frac{\partial f(x, y, y')}{\partial y'}) dx,$$

we can write

$$I(\epsilon) = I(0) + \delta I_1 + O(\epsilon^2),$$

where δI_1 is called the first variation. The vanishing of the first variation is a necessary, but not sufficient, condition to have an extremum. To establish a sufficient condition, assuming that the function is thrice continuously differentiable, we further expand as

$$I(\epsilon) = I(0) + \delta I_1 + \delta I_2 + O(\epsilon^3).$$

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