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# Classification and Identification of Lie Algebras

Libor Šnobl Pavel Winternitz



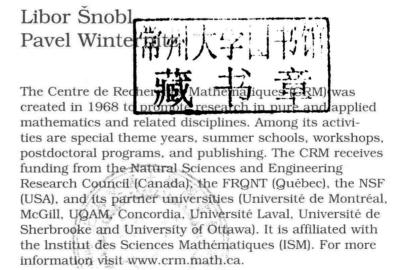
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### C R M

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Classification and Identification of Lie Algebras

#### Preface

The purpose of this book is to serve as a tool for practitioners of Lie algebra and Lie group theory, i.e., for those who apply Lie algebras and Lie groups to solve problems arising in science and engineering. It is not intended to be a textbook on Lie theory, nor is it oriented towards one specific application, for instance the analysis of symmetries of differential equations. We restrict our attention to finite-dimensional Lie algebras over the fields of complex and real numbers.

In any application Lie algebras typically arise as sets of linear operators that commute with a given operator, say the Hamiltonian of a physical system. Alternatively, Lie groups arise as groups of (local) transformations leaving some object invariant; the corresponding Lie algebra then consists of vector fields generating 1-parameter subgroups. The object may be for instance the set of all solutions of a system of equations. The equations can be differential, difference, algebraic or integral ones, or some combination of such equations. They may be linear or nonlinear. In any case, the Lie algebra is realized by some operators in a basis that is usually not the standard one and that depends crucially on the manner in which it was obtained. The structure constants of Lie algebras can be calculated in any basis, but they in turn are basis dependent and reveal very little about the actual structure of the given Lie algebra.

After the Lie algebra  $\mathfrak g$  associated with a studied problem is found, the next task that faces the researcher is to identify the Lie algebra as an abstract Lie algebra. In some cases  $\mathfrak g$  may be isomorphic to a known algebra given in some accessible list. This is certainly the case for semisimple Lie algebras in view of Cartan's classification of all simple Lie algebras over the complex numbers, and subsequent classification of their real forms.

The fundamental Levi theorem, stating that every finite dimensional Lie algebra is isomorphic to a semidirect sum of a semisimple Lie algebra and the maximal solvable ideal (the radical) greatly simplifies the task of identifying a given Lie algebra. The weak link is that no complete classification of solvable Lie algebras exists, nor can one be expected to be produced in the future.

The problem addressed in this book is that of transforming a randomly obtained basis of a Lie algebra into a "canonical basis" in which all basis independent features of the Lie algebra are directly visible. For low dimensional Lie algebras (of dimension less or equal six) this makes it possible to identify the Lie algebra completely. In this book we give a representative list of all such Lie algebras. As stated above, in any dimension a complete identification can be performed for semisimple Lie algebras. We also describe some classes of nilpotent and solvable Lie algebras of arbitrary finite dimensions for which a complete classification exists and hence an exact identification is possible.

x PREFACE

The book has four parts. The first presents some general results and concepts that are used in the subsequent chapters. In particular such invariant notions as the dimension of ideals in the characteristic series, and the invariants of the coadjoint representation are introduced.

In Part 2 we present algorithms that accomplish the following tasks:

- (1) An algorithm for determining whether the algebra  $\mathfrak g$  can be decomposed into a direct sum. If  $\mathfrak g$  is decomposable the algorithm provides a basis in which  $\mathfrak g$  is explicitly decomposed into a direct sum of indecomposable Lie subalgebras.
- (2) A further algorithm is presented to find the radical  $R(\mathfrak{g})$  and the Levi factor, i.e., the semisimple component of  $\mathfrak{g}$ .
- (3) If the Lie algebra is solvable, for instance if it is the radical of a larger algebra, then it is necessary to identify its nilradical, i.e., the maximal nilpotent ideal. A rational (i.e., avoiding calculation of eigenvalues) algorithm for performing this is presented.

The text includes many examples illustrating various situations that may arise in such computations. All these algorithms have been implemented on computers.

Part 3 is devoted to solvable and nilpotent Lie algebras. While a complete classification of such algebras seems not to be feasible, it is possible to take a class of nilpotent Lie algebras and construct all extensions of these algebras to solvable ones. Finite-dimensional solvable Lie algebras with Abelian, Heisenberg, Borel, filiform and quasifiliform nilradicals are presented in Part 3.

Part 4 of the book consists of tables of all indecomposable Lie algebras of dimension n where  $1 \le n \le 6$ . They are ordered in such a way as to make the identification of any given low-dimensional Lie algebra written in an arbitrary basis as simple as possible. Any Lie algebra up to dimension 6 is isomorphic to precisely one entry in the tables. Essential characteristics of each algebra including its Casimir invariants are also provided.

The book is based on material that was previously dispersed in journal articles, many of them written by one or both of the authors of this book together with collaborators. The tables in Part 4 are based on older results and have been independently verified, in some cases corrected, unified and ordered by structural properties of the algebras (rather than by the way they were originally obtained).

Libor Šnobl and Pavel Winternitz

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L. Šnobl dedicates this book to his parents Libuše and Zdeněk. P. Winternitz dedicates this book to his wife Milada and his sons Peter and Michael. We both thank them for their support and encouragement.

#### Contents

Preface		ix
Acknowle	edgements	xi
Part 1.	General Theory	1
Chapter	1. Introduction and Motivation	3
Chapter 2.1. 2.2. 2.3. 2.4.	2. Basic Concepts Definitions Levi theorem Classification of complex simple Lie algebras Chevalley cohomology of Lie algebras	11 11 17 17 20
Chapter 3.1. 3.2.	3. Invariants of the Coadjoint Representation of a Lie Algebra Casimir operators and generalized Casimir invariants Calculation of generalized Casimir invariants using the infinitesimal	23 23
	method	24
3.3.	Calculation of generalized Casimir invariants by the method of moving frames	32
Part 2. Constan	Recognition of a Lie Algebra Given by Its Structure ats	37
Chapter 4.1. 4.2.	4. Identification of Lie Algebras through the Use of Invariants Elementary invariants More sophisticated invariants	39 39 42
Chapter 5.1. 5.2. 5.3.	5. Decomposition into a Direct Sum General theory and criteria Algorithm Examples	47 47 56 57
Chapter 6.1. 6.2. 6.3.	<ol> <li>Levi Decomposition. Identification of the Radical and Levi Factor</li> <li>Original algorithm</li> <li>Modified algorithm</li> <li>Examples</li> </ol>	63 63 65 66
Chapter 7.1. 7.2.	7. The Nilradical of a Lie Algebra General theory Algorithm	71 71 75

7.3. 7.4.	Examples Identification of the nilradical using the Killing form	79 84
Part 3.	Nilpotent, Solvable and Levi Decomposable Lie Algebras	87
Chapter	8. Nilpotent Lie Algebras	89
8.1. 8.2.	Maximal Abelian ideals and their extensions Classification of low-dimensional nilpotent Lie algebras	89 93
Chapter	9. Solvable Lie Algebras and Their Nilradicals	99
9.1. 9.2.	General structure of a solvable Lie algebra General procedure for classifying all solvable Lie algebras with a	99
9.3.	given nilradical Upper bound on the dimension of a solvable extension of a given	99
9.4.	nilradical Particular classes of nilradicals and their solvable extensions	103
9.5.	Vector fields realizing bases of the coadjoint representation of a	105
	solvable Lie algebra	106
Chapter		107
10.1.	Basic structural theorems	107
10.2.	Decomposability properties of the solvable Lie algebras	114
10.3. 10.4.	Solvable Lie algebras with centers of maximal dimension	116
	Solvable Lie algebras with one nonnilpotent element and an <i>n</i> -dimensional Abelian nilradical	121
10.5.	Solvable Lie algebras with two nonnilpotent elements and n-dimensional Abelian nilradical	123
10.6.	Generalized Casimir invariants of solvable Lie algebras with Abelian	1
	nilradicals	125
Chapter	11. Solvable Lie Algebras with Heisenberg Nilradical	131
11.1.	The Heisenberg relations and the Heisenberg algebra	131
11.2.	Classification of solvable Lie algebras with nilradical $\mathfrak{h}(m)$	132
11.3.	The lowest dimensional case $m=1$	134
11.4.	The case $m=2$	135
11.5.	Generalized Casimir invariants	136
Chapter	12. Solvable Lie Algebras with Borel Nilradicals	141
12.1.	Outer derivations of nilradicals of Borel subalgebras	141
12.2.	Solvable extensions of the Borel nilradicals $NR(\mathfrak{b}(\mathfrak{g}))$	146
12.3.	Solvable Lie algebras with triangular nilradicals	153
12.4.	Casimir invariants of nilpotent and solvable triangular Lie algebras	162
Chapter		s175
13.1.	Classification of solvable Lie algebras with the model filiform nilradical $\mathfrak{n}_{n,1}$	176
13.2.	Classification of solvable Lie algebras with the nilradical $\mathfrak{n}_{n,2}$	182
13.3.	Solvable Lie algebras with other filiform nilradicals	189
13.4.	Example of an almost filiform nilradical	190
13.5.	Generalized Casimir invariants of $\mathfrak{n}_{n,3}$ and of its solvable extensions	

CONTENTS	vii
----------	-----

Chapter 14. Levi Decomposable Algebras 14.1. Levi decomposable algebras with a nilpotent radical 14.2. Levi decomposable algebras with nonnilpotent radicals 14.3. Levi decomposable algebras of low dimensions	203 204 207 208
Part 4. Low-Dimensional Lie Algebras	215
Chapter 15. Structure of the Lists of Low-Dimensional Lie Algebras 15.1. Ordering of the lists 15.2. Computer-assisted identification of a given Lie algebra	217 217 218
Chapter 16. Lie Algebras up to Dimension 3 16.1. One-dimensional Lie algebra 16.2. Solvable two-dimensional Lie algebra with the nilradical $\mathfrak{n}_{1,1}$ 16.3. Nilpotent three-dimensional Lie algebra 16.4. Solvable three-dimensional Lie algebras with the nilradical $2\mathfrak{n}_{1,1}$ 16.5. Simple three-dimensional Lie algebras	225 225 225 225 226 226
Chapter 17. Four-Dimensional Lie Algebras 17.1. Nilpotent four-dimensional Lie algebra 17.2. Solvable four-dimensional algebras with the nilradical $\mathfrak{3n}_{1,1}$ 17.3. Solvable four-dimensional Lie algebras with the nilradical $\mathfrak{n}_{3,1}$ 17.4. Solvable four-dimensional Lie algebras with the nilradical $\mathfrak{2n}_{1,1}$	227 227 227 228 229
Chapter 18. Five-Dimensional Lie Algebras  18.1. Nilpotent five-dimensional Lie algebras  18.2. Solvable five-dimensional Lie algebras with the nilradical $\mathfrak{n}_{3,1} \oplus \mathfrak{n}_{3,1} \oplus \mathfrak{n}_$	231 231 232 11,1 235 239 240 241 241
Chapter 19. Six-Dimensional Lie Algebras  19.1. Nilpotent six-dimensional Lie algebras with the nilradical $5\mathfrak{n}_{1,1}$ 19.2. Solvable six-dimensional Lie algebras with the nilradical $\mathfrak{n}_{3,1} \oplus 2\mathfrak{n}_{1,1}$ 19.3. Solvable six-dimensional Lie algebras with the nilradical $\mathfrak{n}_{4,1} \oplus \mathfrak{n}_{1,1}$ 19.4. Solvable six-dimensional Lie algebras with the nilradical $\mathfrak{n}_{4,1} \oplus \mathfrak{n}_{1,1}$ 19.5. Solvable six-dimensional Lie algebras with the nilradical $\mathfrak{n}_{5,1}$ 19.6. Solvable six-dimensional Lie algebras with the nilradical $\mathfrak{n}_{5,2}$ 19.7. Solvable six-dimensional Lie algebras with the nilradical $\mathfrak{n}_{5,3}$ 19.8. Solvable six-dimensional Lie algebras with the nilradical $\mathfrak{n}_{5,4}$ 19.9. Solvable six-dimensional Lie algebras with the nilradical $\mathfrak{n}_{5,5}$ 19.10. Solvable six-dimensional Lie algebras with the nilradical $\mathfrak{n}_{5,6}$ 19.11. Solvable six-dimensional Lie algebras with the nilradical $\mathfrak{n}_{3,1} \oplus \mathfrak{n}_{3,1}$ 19.12. Solvable six-dimensional Lie algebras with the nilradical $\mathfrak{n}_{3,1} \oplus \mathfrak{n}_{3,1}$ 19.13. Solvable six-dimensional Lie algebra with the nilradical $\mathfrak{n}_{4,1}$	1,1 266 271 277 279 283 285 286 286
19.14. Simple six-dimensional Lie algebra  19.15. Six-dimensional Levi decomposable Lie algebras	296 296

viii	CONTENTS
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Bibliography	299
Index	305

# Part 1 General Theory

#### CHAPTER 1

#### Introduction and Motivation

Lie groups and Lie algebras appear in science in many different guises. They may be a priori parts of the theory, like Lorentz or Galilei invariance of most physical theories, or the (semi)simple Lie groups of the Standard model in particle theory.

Alternatively, specific Lie groups may appear as consequences of specific dynamics. Consider any physical system with dynamics described by a system of ordinary or partial differential equations. This system of equations will be invariant under some local Lie group of local point transformations, taking solutions into solutions. This symmetry group G and its Lie algebra  $\mathfrak g$  can be determined in an algorithmic manner [86]. The Lie algebra  $\mathfrak g$  is obtained as an algebra of vector fields, usually in some nonstandard basis, depending on the way in which the algorithm is applied.

An immediate task is to identify the algebra found as being isomorphic to some known abstract Lie algebra. To do this we must transform it to a canonical basis in which all basis independent properties are manifest. Thus, if  $\mathfrak g$  is decomposable into a direct sum, it should be explicitly decomposed into components that are further indecomposable

$$\mathfrak{g}=\mathfrak{g}_1\oplus\mathfrak{g}_2\oplus\cdots\oplus\mathfrak{g}_k.$$

Each indecomposable component must be further identified. Let  $\mathfrak g$  now denote such an indecomposable Lie algebra. A fundamental theorem due to E. E. Levi [59,71] tells us that any finite-dimensional Lie algebra can be represented as the semidirect sum

$$\mathfrak{g}=\mathfrak{p} \ni \mathfrak{r}, \quad [\mathfrak{p},\mathfrak{p}]=\mathfrak{p}, \quad [\mathfrak{r},\mathfrak{r}] \subset \mathfrak{r}, \quad [\mathfrak{p},\mathfrak{r}] \subseteq \mathfrak{r},$$

where the subalgebra  $\mathfrak p$  is semisimple and  $\mathfrak r$  is the radical of  $\mathfrak g$ , i.e., its maximal solvable ideal. If  $\mathfrak g$  is simple, we have  $\mathfrak r=0$  (an indecomposable semisimple algebra is actually simple). If  $\mathfrak g$  is solvable, we have  $\mathfrak p=0$ . Algorithms realizing decompositions (1.1), (1.2) exist [102] and are presented below.

In view of the Levi theorem (1.2) the classification of all finite-dimensional Lie algebras can be reduced to three steps:

- Classification of all simple Lie algebras. This also provides a classification of all semisimple ones.
- (2) Classification of all solvable Lie algebras.
- (3) Classification of all possible linear actions of the semisimple algebra  $\mathfrak{p}$  on the radical  $\mathfrak{r}$ .

Semisimple Lie algebras over the field of complex numbers  $\mathbb{C}$  have been completely classified by W. Killing and É. Cartan [22,60], over the field of real numbers  $\mathbb{R}$  by É. Cartan in [20,21] (the analysis in the real case was later simplified by F. Gantmacher in [50]).

The third step is basically a task of the representation theory of semisimple Lie algebras. The main unsolved problem is the classification of all finite dimensional solvable Lie algebras, a task that does not have a realistic solution. Even all nilpotent Lie algebras are impossible to classify. (For instance in [127] the authors claim that they have found 24 168 nonisomorphic 9-dimensional nilpotent algebras Lie algebras with a maximal Abelian ideal of dimension 7 alone.) A realistic partial classification problem is to classify all solvable Lie algebras with a given nilradical of an arbitrary finite dimension n. So far this has been done for certain series of nilpotent Lie algebras, namely Abelian, Heisenberg and Borel nilradicals, as well as certain filiform and quasifiliform algebras. The results of the original articles [83, 84, 106, 115-118, 123, 124] are presented in a unified manner in Part 3 of this book.

An interesting physical application of the classification of low-dimensional Lie algebras is in general relativity. Indeed, the classification of Einstein spaces according to their isometry groups [95] is based on the work of Bianchi and his successors [6,65]. The Petrov classification concerns Einstein spaces of dimension 4 and hence involves isometry groups of relatively low dimensions [95, 120]. String theory, brane cosmology and some other elementary particle theories going beyond the standard model require the use of higher-dimensional spaces. Any attempt at a Lie group classification of such spaces will require knowledge of higher-dimensional Lie groups, including solvable ones.

One of the very useful applications of Lie group analysis in science is the identification of seemingly different problems that are mathematically equivalent. Indeed, let us consider two systems, say (A) and (B), of differential equations

(1.3) 
$$E_{a}^{A}(\vec{X}, \vec{U}; \partial_{X_{i}}U_{\alpha}, \partial_{X_{i}X_{j}}^{2}U_{\alpha}, \dots) = 0,$$

$$E_{a}^{B}(\vec{x}, \vec{u}; \partial_{x_{i}}u_{\alpha}, \partial_{x_{i}x_{j}}^{2}u_{\alpha}, \dots) = 0,$$

$$\vec{X} = (X_{1}, \dots, X_{p}), \qquad \vec{U} = (U_{1}, \dots, U_{q}),$$

$$\vec{x} = (x_{1}, \dots, x_{p}), \qquad \vec{u} = (u_{1}, \dots, u_{q}),$$

$$1 \le a \le N, \quad 1 \le \alpha \le q, \quad 1 \le i, j, \dots \le p$$

of the same order describing different physical processes. From the mathematical viewpoint, we may consider these two systems equivalent if there exists a local invertible transformation of the independent and dependent variables

(1.4) 
$$X_i = \Lambda_i(\vec{x}, \vec{u}), \qquad U_\alpha = \Omega_\alpha(\vec{x}, \vec{u})$$

transforming the systems (A) and (B) into each other. The transformation (1.4) together with the mapping between the spaces of functions  $\vec{u}(\vec{x})$  and  $\vec{U}(\vec{X})$  induced by it is called a *point transformation*.

A necessary condition for the equivalence of the systems (A) and (B) is that they have isomorphic Lie algebras of infinitesimal point symmetries  $\mathfrak{g}_A$  and  $\mathfrak{g}_B$ . The algebras  $\mathfrak{g}_A$  and  $\mathfrak{g}_B$  are by construction realized by vector fields in variables  $(\vec{X}, \vec{U})$  and  $(\vec{x}, \vec{u})$ , respectively.

If  $\mathfrak{g}_A$  and  $\mathfrak{g}_B$  are isomorphic then a local point transformation (1.4) may exist such that it transforms the vector fields of  $\mathfrak{g}_A$  into those of  $\mathfrak{g}_B$  and vice versa. The transformation is unique up to point transformations leaving  $\mathfrak{g}_A$  or  $\mathfrak{g}_B$  invariant. In any case, the transformation taking  $\mathfrak{g}_A$  into  $\mathfrak{g}_B$  will also take the system (A)