Insulating Films on Semiconductors

Editors: M. Schulz and G. Pensl



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Editors: M. Schulz and G. Pensl

With 240 Figures



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Preface

The INFOS 81 Conference on Insulating Films on Semiconductors was held at the University of Erlangen-Nürnberg in Erlangen from 27 to 29 April 1981. This conference was a sequel to the first conference INFOS 79 held in Durham. INFOS 81 attracted 170 participants from universities, research institutes and industry. Attendants were registered from 15 nations. The biannual topical conference series will be continued by INFOS 83 to be held

in Eindhoven, The Netherlands, in April 1983.

The conference proceedings include all the invited (9) and contributed (42) papers presented at the meeting. The topics range from the basic physical understanding of the properties of insulating films and their interface to semiconductors to the discussion of stability and dielectric strength as well as growing and deposition techniques which are relevant for technical applications. Strong emphasis was given to the semiconductor silicon and its native oxide; however, sessions on compound semiconductors and other insulating films also raised strong interest. The proceedings survey the present state of our understanding of the system of insulating films on semiconductors. As a new aspect of the topic, the properties of semiconductors deposited and laser processed on insulating films was included for the first time.

The organisational load of the conference was distributed on many collaborators. The service of the session chairmen who together with many participants reviewed all the publication manuscripts during the conference, was very much appreciated. Special thanks go to all the colleagues and students of the Institute of Applied Physics who assisted in the local organization. The conference administration and the retyping of many manuscripts in the proceedings was performed with great care and engagement by Mrs. Gabriele

Loy. All the help is gratefully acknowledged.

Erlangen, April 1981

M. Schulz · G. Pensl

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Part I:

Si-SiO₂ Interface

Electronic Structure of the Si-SiO₂ Interface

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After sketching various theoretical models that have been proposed for describing the Si/SiO2 interface on an atomic scale, we will discuss an idealized model which is designed to simulate the average contact between a silicon crystal and its oxide overlayer. According to this model, the Si/SiO2 interface is represented by the boundary between two crystalline domains, the first being the silicon substrate, and the second an idealized crystalline form of SiO₂, diamond-like beta cristobalite. Nearly perfect registry between Si and SiO_2 is obtained by placing the (100) face of the former next to the 45° rotated (100) face of the latter. Half the Si atoms at the interface are fourfold coordinated, connecting the Si and SiO2 regions to one another. The remaining half are connected only to the silicon substrate and are thus twofold coordinated. Localized electronic states in the thermal gap at the interface are associated with the dangling bonds at the twofold coordinated Si atoms. If H, O, or OH groups are attached to these Si atoms, eliminating the dangling bonds, the thermal gap interface states are also eliminated. Even though this model ignores structural disorder, it does provide considerable insight into the nature of the Si/SiO2 interface. We will discuss some of the implications of this model, as well as the results of model-based theoretical calculations. In particular, we will consider the electronic structure of the interface including the effects produced by the introduction of Si and O vacancies and interstitials and other types of localized defects.

1. Introduction

Because of the important role played by the Si/SiO₂ interface in modern semi-conductor technology, considerable experimental and theoretical effort has been devoted to the study of this interface [2-8]. Oxidation of a silicon surface produces an amorphous oxide layer whose atomic arrangement is not well

understood. Many authors have suggested that there is a non-stoichiometric transition layer, often described as $\mathrm{Si0}_{\mathrm{X}}$, separating the silicon substrate from the stoichiometric $\mathrm{Si0}_2$ layer. (Here x lies between 1 and 2.) The detailed atomic structure of the transition layer, including its uniformity and width, its structural disorder, its affinity for impurity atoms, etc., continues to be a subject of debate. Of course, the nature of the transition layer is expected to depend on the manner in which the oxide layer is formed. By the same token, a better understanding of the transition layer could lead to further improvements in the preparation of $\mathrm{Si/Si0}_2$ interfaces.

The technological importance of the $\mathrm{Si/Si0}_2$ interface arises from the fortunate circumstance that suitably oxidized silicon substrates are protected from environmental contaminants and have rather low densities of localized interface states in the thermal gap. There continues to be considerable interest in the nature and identity of the residual defects that occur near well-prepared interfaces, giving rise to residual trapping sites and fixed charge. Many imaginative models of defects at the $\mathrm{Si/Si0}_2$ interface and in $\mathrm{Si0}_2$ have been proposed [2-11], but relatively few of these defects have been unambiguously identified, even by such precise measurements as ESR [11]. In view of the difficulties encountered in practice in attempting to interpret most experimental measurements in terms of specific defects, we thought it would be useful to develop a simple model for the $\mathrm{Si/Si0}_2$ interface which could provide a conceptual framework for thinking about interface defects as well as a theoretical framework for carrying out numerical calculations.

In the first paper on this subject [12], we described the construction of idealized $\mathrm{Si/Si0}_2$ interfaces and superlattices and demonstrated by first-principles band structure calculations that there are no localized interface states in the thermal gap if all the Si atoms at the interface are saturated. In a second paper [13] we showed how localized defects in $\mathrm{Si0}_2$ could be investigated theoretically by introducing periodic arrays of noninteracting defects in a $\mathrm{Si0}_2$ supercell normally containing 24 atoms. In a third paper [14] we introduced periodic arrays of noninteracting defects in idealized $\mathrm{Si/Si0}_2$ superlattices, thereby simulating the presence of localized defects at or near the $\mathrm{Si/Si0}_2$ interface. In the present paper we will first mention some recent attempts by others to develop a theoretical model of $\mathrm{Si0}_2$ and the $\mathrm{Si/Si0}_2$ interface, and then we will discuss our own model briefly and summarize some of our principal conclusions.

2. Theoretical Models of the Si/SiO₂ Interface

The extensive literature on crystalline and non-crystalline ${\rm SiO}_2$ and on ${\rm Si/SiO}_2$ interfaces [2-11] contains a variety of models of the oxide layer and the ${\rm Si/SiO}_2$ interface. In this paper, we will focus on atomic scale rather than on phenomenological models. Nearly all the models that have been proposed for describing the atomic arrangements at the interface are speculative. Some of the proposals are quite specific. They assume particular types of linkages between Si and ${\rm SiO}_2$, well-defined structural imperfections arising from modifications in the normal linkages, as well as chemical imperfections due to the presence of specific foreign species. The remaining proposals tend to be descriptive or impressionistic, invoking dislocations, disjunctions, varying degrees of stoichiometry and crystallinity, etc. The first set of proposals can be distinguished from the second by the ease with which ball-and-stick models can be constructed to illustrate the proposed arrangements.

There has also been considerable debate regarding the initial formation of the oxide layer, and how this layer might change as it grows [2-8]. Some authors believe that a better understanding of the initial stages of oxidation could provide valuable clues concerning the nature of the ultimate oxide layer; other authors disagree. In any event, it is very important to have heuristic atomic models of the formation and ultimate structure of the Si/SiO₂ interface, because such ideas could lead to improved processing techniques, new device concepts, and better understanding of still more complex semiconductor/insulator interfaces.

Up to now, most $\mathrm{Si/Si0}_2$ interface modeling studies have been qualitative, being based on simple physical and chemical reasoning. Many of the arguments put forward in favor of specific atomic models seem quite plausible, but it is usually difficult to verify these interface models experimentally. Therefore, a constructive step would be to study some of these models theoretically and obtain quantitative estimates of crucial electronic properties. If one wishes to carry out detailed electronic structure calculations for model systems, it is obviously necessary to have specific atomic arrangements in mind. The essential difficulty is that we do not know the actual structure of the $\mathrm{Si/Si0}_2$ interface on an atomic scale.

The overall problem can be dealt with in three parts: In the first and easiest part, one makes certain assumptions about the silicon substrate. For example, one chooses a specific crystalline orientation, and then assumes that the silicon is structurally perfect up to the interface. Relaxation and reconstruction effects can be deferred to a later stage or expressly ignored.

The second part concerns the description of the oxide film property, allowing if possible for its noncrystalline nature. In dealing with non-crystalline SiO₂, one can build on our improved knowledge of amorphous and glassy semiconductors [9,15-19]. The third part, the most difficult of all, involves connecting the silicon substrate to the non-crystalline oxide layer in a physically and chemically plausible manner.

One popular model for noncrystalline ${\rm SiO}_2$ and the ${\rm Si/SiO}_2$ interface is the cluster Bethe lattice model [20,21], which takes topological aspects of structural disorder into account in an elegant manner. Because this model is rather simple in form and can be easily parameterized, it can be used to study electronic and lattice vibrational spectral properties associated with ideal and defective interfaces. But because these models are based on adjustable parameters that are assigned ad hoc values, they are not particularly trustworthy from a quantitative point of view [22]. Nevertheless, they do have considerable pedagogical value.

Another popular model of non-crystalline ${\rm SiO}_2$ is the continuous random network model [15-19], according to which ${\rm SiO}_2$ is composed of a randomly oriented collection of ${\rm SiO}_4$ tetrahedra which are linked to one another to form a continuous network of ${\rm Si-O-Si}$ bonds. It is assumed that the ${\rm SiO}_4$ tetrahedra tend to maintain their structural integrity (bond lengths and bond angles) in different environments, and that the necessary structural versatility is provided by the ${\rm Si-O-Si}$ bond, whose angle may assume a wide range of values at little cost in energy [4]. Models of this type have recently been used to study bulk ${\rm SiO}_2$ [23] as well as the ${\rm Si/SiO}_2$ interface [24]. In the study of bulk ${\rm SiO}_2$, cyclic boundary conditions were imposed on a large disordered molecular ${\rm SiO}_2$ cluster, producing a computationally tractable model of non-crystalline ${\rm SiO}_2$. The object here was to investigate electron and hole localization induced by the structural disorder.

In the study of the $\mathrm{Si/Si0}_2$ interface, it was demonstrated that a rather abrupt interface could be constructed within the framework of a continuous random network model, a rather important result. Unfortunately, the resulting atomic model is much too complex for detailed electronic structure studies. In both [23] and [24], the construction of continuous random network models is based on ideas regarding the nature of non-crystallinity which are highly subjective and which may not necessarily describe the experimental situation. It is possible that non-crystalline $\mathrm{Si0}_2$ near a $\mathrm{Si/Si0}_2$ interface does not actually form a completely continuous random network. That is to say, the network may contain discontinuous elements such as non-bridging 0 atoms or three-

fold coordinated Si atoms connected to H or OH groups, or voids composed of aggregates of such discontinuous elements [19].

There are many theoretical efforts currently underway [18,19,25] whose aim is to formulate general guiding principles for constructing atomic models of non-crystalline materials. Combined with intuitive notions [26] and computer modeling [27,28], these theoretical studies may eventually lead us to objective models which we can use with confidence for electronic structure studies of non-crystalline ${\rm SiO}_2$. In the meanwhile, the atomic architecture of the ${\rm Si/SiO}_2$ interface remains an open question.

We already know from extensive experimental and theoretical studies [2,3,29] that the gross features of the electronic structure of non-crystalline and various crystalline forms of SiO_2 are essentially the same, reflecting the common SiO_4 building block. Only the finer details of electronic structure are affected by the special manner in which the SiO_4 tetrahedra are linked together in a particular structural modification. In view of the overwhelming importance of the short-range order (linked SiO_4 tetrahedra) in determining the essential features of the electronic structure of SiO_2 , and the likelihood that the long-range disorder affects only the finer details, we will explicitly ignore the structural disorder of SiO_2 and concentrate on models of the interface which represent the average contact between Si and SiO_2 .

3. Idealized Si/SiO₂ Superlattices and Interfaces

Our overall program is as follows. First, we will disregard structural disorder in SiO_2 , and represent SiO_2 by a suitably chosen periodic structure. We will then form ordered $\mathrm{Si/SiO}_2$ interfaces by attaching crystalline SiO_2 to crystalline Si . Next, we will construct a $\mathrm{Si/SiO}_2$ superlattice by stacking Si and SiO_2 slabs on top of one another. The slabs will be made sufficiently thick so that successive interfaces will not interact with one another. Because the $\mathrm{Si/SiO}_2$ superlattice is a periodic structure in three dimensions, we can use the highly developed methods of band theory [30] to calculate the electronic structure. The results of such calculations include a description of localized interface states, including their energy levels and charge distributions. Finally, we can introduce structural and chemical imperfections, and study their effects on the electronic structure.

To carry out the above program, we begin by representing the oxide layer by an idealized crystalline form of SiO_2 having as simple a unit cell as pos-

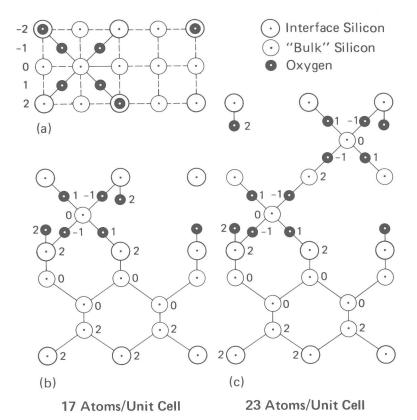


Fig. 1a-c. Si/SiO_2 superlattices and interfaces. Panels (a) and (b) depict the top and front views of a superlattice having 17 atoms per unit cell, and panel (c) depicts the front view of a superlattice having 23 atoms per unit cell. Both superlattices have five layers of silicon (2 atoms each) and two oxygen atoms saturating the twofold coordinated Si substrate atoms. In panel (b) there is one tier of SiO_4 tetrahedra (5 atoms), while in panel (c) there are three SiO_4 tiers (11 atoms). The numbers denote the planes in which the various atoms lie, the 2 representing the front-most plane, etc. Another view of these superlattices appears in [8]

sible. For this purpose we will use beta cristobalite [31], except that we will straighten out the Si-O-Si bonds in the actual structure, making them all linear rather than crooked (all bond angles 180° rather than about 140°). Assuming an Si-O bond length of 1.633 Å [32], the linear Si-O-Si bond length becomes twice this, and the unit cube edge of idealized diamond-like SiO_2 becomes 7.543 Å, which is 5 percent larger than the unit cube edge of actual beta cristobalite, 7.16 Å. This 5 percent difference reflects the contraction produced by the bent bonds in the actual structure. Except as otherwise noted,