

TURING

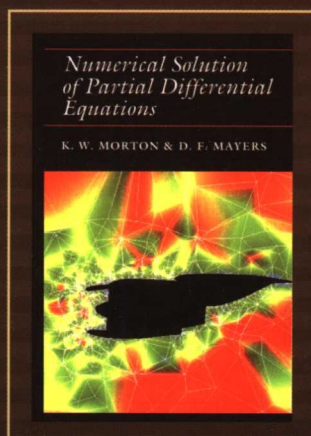
图灵原版数学·统计学系列

Numerical Solution of Partial Differential Equations

偏微分方程数值解

(英文版·第2版)

[英] K. W. Morton 著
D. F. Mayers



人民邮电出版社
POSTS & TELECOM PRESS

TURING

图灵原版数学·统计学系列

偏微分方程数值解

(英文版·第2版)

Numerical Solution of Partial Differential Equations

[英] K. W. Morton 著
D. F. Mayers



人民邮电出版社

POSTS & TELECOM PRESS

图书在版编目(CIP)数据

偏微分方程数值解: 第2版 / (英)莫顿著, 迈耶斯著. —北京: 人民邮电出版社, 2006.1
(图灵原版数学·统计学系列)

ISBN 7-115-14075-8

I. 偏... II. ①莫...②迈... III. 偏微分方程—数值计算—英文 IV. 0241.82

中国版本图书馆 CIP 数据核字 (2005) 第 116947 号

内 容 提 要

偏微分方程是构建科学、工程学和其他领域的数学模型的主要手段。一般情况下, 这些模型都需要用数值方法去求解。本书提供了标准数值技术的简明介绍。借助抛物线型、双曲线型和椭圆型方程的一些简单例子介绍了常用的有限差分方法、有限元方法、有限体方法、修正方程分析、辛积分格式、对流扩散问题、多重网格、共轭梯度法。利用极大值原理、能量法和离散傅里叶分析清晰严格地处理了稳定性问题。本书全面讨论了这些方法的性质, 并附有典型的图像结果, 提供了不同难度的例子和练习。

本书可作为数学、工程学及计算机专业本科教材, 也可供工程技术人员和应用工作者参考。

图灵原版数学·统计学系列

偏微分方程数值解 (英文版·第2版)

-
- ◆ 著 [英] K.W.Morton D.F.Mayers
责任编辑 王丽萍
 - ◆ 人民邮电出版社出版发行 北京市崇文区夕照寺街14号
邮编 100061 电子函件 315@ptpress.com.cn
网址 <http://www.ptpress.com.cn>
北京顺义振华印刷厂印刷
新华书店总店北京发行所经销
 - ◆ 开本: 800×1000 1/16
印张: 18.25
字数: 389千字 2006年1月第1版
印数: 1~3000册 2006年1月北京第1次印刷

著作权合同登记号 图字: 01-2005-5233 号

ISBN 7-115-14075-8/TP · 5016

定价: 39.00 元

读者服务热线: (010)88593802 印装质量热线: (010)67129223

版权声明

K. W. Morton and D. F. Mayers: *Numerical Solution of Partial Differential Equations* (ISBN 0-521-60793-0).

Originally published by Cambridge University Press in 2005.

This reprint edition is published with the permission of the Syndicate of the Press of the University of Cambridge, Cambridge, England.

Copyright © 2005 by Cambridge University Press.

This edition is licensed for distribution and sale in the People's Republic of China only, excluding Hong Kong, Taiwan and Macao, and may not be distributed and sold elsewhere.

本书原版由剑桥大学出版社出版。

本书英文影印版由剑桥大学出版社授权人民邮电出版社出版。

此版本仅限在中华人民共和国境内（中国香港、台湾、澳门地区除外）销售发行。

版权所有，侵权必究。

Preface to the first edition

The origin of this book was a sixteen-lecture course that each of us has given over the last several years to final-year Oxford undergraduate mathematicians; and its development owes much to the suggestions of some of our colleagues that the subject matter could readily be taught somewhat earlier as a companion course to one introducing the theory of partial differential equations. On the other hand, we have used much of the same material in teaching a one-year Master's course on mathematical modelling and numerical analysis. These two influences have guided our choice of topics and the level and manner of presentation. -

Thus we concentrate on finite difference methods and their application to standard model problems. This allows the methods to be couched in simple terms while at the same time treating such concepts as stability and convergence with a reasonable degree of mathematical rigour. In a more advanced text, or one with greater emphasis on the finite element method, it would have been natural and convenient to use standard Sobolev space norms. We have avoided this temptation and used only discrete norms, specifically the maximum and the l_2 norms. There are several reasons for this decision. Firstly, of course, it is consistent with an aim of demanding the minimum in prerequisites – of analysis, of PDE theory, or of computing – so allowing the book to be used as a text in an early undergraduate course and for teaching scientists and engineers as well as mathematicians.

Equally importantly though, the decision fits in with our widespread use of discrete maximum principles in analysing methods for elliptic and parabolic problems, our treatment of discrete energy methods and conservation principles, and the study of discrete Fourier modes on finite domains. We believe that treating all these ideas at a purely discrete level helps to strengthen the student's understanding of these important

mathematical tools. At the same time this is a very practical approach, and it encourages the interpretation of difference schemes as direct models of physical principles and phenomena: all calculations are, after all, carried out on a finite grid, and practical computations are checked for stability, etc. at the discrete level. Moreover, interpreting a difference scheme's effect on the Fourier modes that can be represented on the mesh, in terms of the damping and dispersion in one time step is often of greater value than considering the truncation error, which exemplifies the second justification of our approach.

However, the limiting process as a typical mesh parameter h tends to zero is vital to a proper understanding of numerical methods for partial differential equations. For example, if U^n is a discrete approximation at time level n and evolution through a time step Δt is represented as $U^{n+1} = C_h U^n$, many students find great difficulty in distinguishing the limiting process when $n \rightarrow \infty$ on a fixed mesh and with fixed Δt , from that in which $n \rightarrow \infty$ with $n\Delta t$ fixed and $h, \Delta t \rightarrow 0$. Both processes are of great practical importance: the former is related to the many iterative procedures that have been developed for solving the discrete equations approximating steady state problems by using the analogy of time stepping the unsteady problem; and understanding the latter is crucial to avoiding instability when choosing methods for approximating the unsteady problems themselves. The notions of uniform bounds and uniform convergence lie at the heart of the matter; and, of course, it is easier to deal with these by using norms which do not themselves depend on h . However, as shown for example by Palencia and Sanz-Serna,¹ a rigorous theory can be based on the use of discrete norms and this lies behind the approach we have adopted. It means that concepts such as well-posedness have to be rather carefully defined; but we believe the slight awkwardness entailed here is more than compensated for by the practical and pedagogical advantages pointed out above.

The ordering of the topics is deliberate and reflects the above concerns. We start with parabolic problems, which are both the simplest to approximate and analyse and also of widest utility. Through the addition of convection to the diffusion operator, this leads naturally to the study of hyperbolic problems. It is only after both these cases have been explored in some detail that, in Chapter 5, we present a careful treatment of the concepts of consistency, convergence and stability for evolutionary problems. The final two chapters are devoted respectively

¹ Palencia, C. and Sanz-Serna, J. M. (1984), An extension of the Lax-Richtmyer theory, *Numer. Math.* **44** (2), 279–283.

to the discretisation of elliptic problems, with a brief introduction to finite element methods, and to the iterative solution of the resulting algebraic equations; with the strong relationship between the latter and the solution of parabolic problems, the loop of linked topics is complete. In all cases, we present only a small number of methods, each of which is thoroughly analysed and whose practical utility we can attest to. Indeed, we have taken as a theme for the book that all the model problems and the methods used to approximate them are simple but generic.

Exercises of varying difficulty are given at the end of each chapter; they complete, extend or merely illustrate the text. They are all analytical in character, so the whole book could be used for a course which is purely theoretical. However, numerical analysis has very practical objectives, so there are many numerical illustrations of the methods given in the text; and further numerical experiments can readily be generated for students by following these examples. Computing facilities and practices develop so rapidly that we believe this open-ended approach is preferable to giving explicit practical exercises.

We have referred to the relevant literature in two ways. Where key ideas are introduced in the text and they are associated with specific original publications, full references are given in footnotes – as earlier in this Preface. In addition, at the end of each chapter we have included a brief section entitled ‘Bibliographic notes and recommended reading’ and the accumulated references are given at the end of the book. Neither of these sets of references is intended to be comprehensive, but they should enable interested students to pursue further their studies of the subject. We have, of course, been greatly guided and influenced by the treatment of evolutionary problems in Richtmyer and Morton (1967); in a sense the present book can be regarded as both introductory to and complementary to that text.

We are grateful to several of our colleagues for reading and commenting on early versions of the book (with Endre Süli’s remarks being particularly helpful) and to many of our students for checking the exercises. The care and patience of our secretaries Jane Ellory and Joan Himpson over the long period of the book’s development have above all made its completion possible.

Preface to the second edition

In the ten years since the first edition of this book was published, the numerical solution of PDEs has moved forward in many ways. But when we sought views on the main changes that should be made for this second edition, the general response was that we should not change the main thrust of the book or make very substantial changes. We therefore aimed to limit ourselves to adding no more than 10%–20% of new material and removing rather little of the original text: in the event, the book has increased by some 23%.

Finite difference methods remain the starting point for introducing most people to the solution of PDEs, both theoretically and as a tool for solving practical problems. So they still form the core of the book. But of course finite element methods dominate the elliptic equation scene, and finite volume methods are the preferred approach to the approximation of many hyperbolic problems. Moreover, the latter formulation also forms a valuable bridge between the two main methodologies. Thus we have introduced a new section on this topic in Chapter 4; and this has also enabled us to reinterpret standard difference schemes such as the Lax–Wendroff method and the box scheme in this way, and hence for example show how they are simply extended to nonuniform meshes. In addition, the finite element section in Chapter 6 has been followed by a new section on convection–diffusion problems: this covers both finite difference and finite element schemes and leads to the introduction of Petrov–Galerkin methods.

The theoretical framework for finite difference methods has been well established now for some time and has needed little revision. However, over the last few years there has been greater interaction between methods to approximate ODEs and those for PDEs, and we have responded to this stimulus in several ways. Firstly, the growing interest in applying

symplectic methods to Hamiltonian ODE systems, and extending the approach to PDEs, has led to our including a section on this topic in Chapter 4 and applying the ideas to the analysis of the staggered leap-frog scheme used to approximate the system wave equation. More generally, the revived interest in the method of lines approach has prompted a complete redraft of the section on the energy method of stability analysis in Chapter 5, with important improvements in overall coherence as well as in the analysis of particular cases. In that chapter, too, is a new section on modified equation analysis: this technique was introduced thirty years ago, but improved interpretations of the approach for such as the box scheme have encouraged a reassessment of its position; moreover, it is again the case that its use for ODE approximations has both led to a strengthening of its analysis and a wider appreciation of its importance.

Much greater changes to our field have occurred in the practical application of the methods we have described. And, as we continue to have as our aim that the methods presented should properly represent and introduce what is used in practice, we have tried to reflect these changes in this new edition. In particular, there has been a huge improvement in methods for the iterative solution of large systems of algebraic equations. This has led to a much greater use of implicit methods for time-dependent problems, the widespread replacement of direct methods by iterative methods in finite element modelling of elliptic problems, and a closer interaction between the methods used for the two problem types. The emphasis of Chapter 7 has therefore been changed and two major sections added. These introduce the key topics of multigrid methods and conjugate gradient methods, which have together been largely responsible for these changes in practical computations.

We gave serious consideration to the possibility of including a number of MATLAB programs implementing and illustrating some of the key methods. However, when we considered how very much more powerful both personal computers and their software have become over the last ten years, we realised that such material would soon be considered outmoded and have therefore left this aspect of the book unchanged. We have also dealt with references to the literature and bibliographic notes in the same way as in the earlier edition: however, we have collected both into the reference list at the end of the book.

Solutions to the exercises at the end of each chapter are available in the form of L^AT_EX files. Those involved in teaching courses in this area can obtain copies, by email only, by applying to solutions@cambridge.org.

We are grateful to all those readers who have informed us of errors in the first edition. We hope we have corrected all of these and not introduced too many new ones. Once again we are grateful to our colleagues for reading and commenting on the new material.

Contents

1	Introduction	1
2	Parabolic equations in one space variable	7
2.1	Introduction	7
2.2	A model problem	7
2.3	Series approximation	9
2.4	An explicit scheme for the model problem	10
2.5	Difference notation and truncation error	12
2.6	Convergence of the explicit scheme	16
2.7	Fourier analysis of the error	19
2.8	An implicit method	22
2.9	The Thomas algorithm	24
2.10	The weighted average or θ -method	26
2.11	A maximum principle and convergence for $\mu(1 - \theta) \leq \frac{1}{2}$	33
2.12	A three-time-level scheme	38
2.13	More general boundary conditions	39
2.14	Heat conservation properties	44
2.15	More general linear problems	46
2.16	Polar co-ordinates	52
2.17	Nonlinear problems	54
	<i>Bibliographic notes</i>	56
	<i>Exercises</i>	56

3	2-D and 3-D parabolic equations	62
3.1	The explicit method in a rectilinear box	62
3.2	An ADI method in two dimensions	64
3.3	ADI and LOD methods in three dimensions	70
3.4	Curved boundaries	71
3.5	Application to general parabolic problems	80
	<i>Bibliographic notes</i>	83
	<i>Exercises</i>	83
4	Hyperbolic equations in one space dimension	86
4.1	Characteristics	86
4.2	The CFL condition	89
4.3	Error analysis of the upwind scheme	94
4.4	Fourier analysis of the upwind scheme	97
4.5	The Lax–Wendroff scheme	100
4.6	The Lax–Wendroff method for conservation laws	103
4.7	Finite volume schemes	110
4.8	The box scheme	116
4.9	The leap-frog scheme	123
4.10	Hamiltonian systems and symplectic integration schemes	128
4.11	Comparison of phase and amplitude errors	135
4.12	Boundary conditions and conservation properties	139
4.13	Extensions to more space dimensions	143
	<i>Bibliographic notes</i>	146
	<i>Exercises</i>	146
5	Consistency, convergence and stability	151
5.1	Definition of the problems considered	151
5.2	The finite difference mesh and norms	152
5.3	Finite difference approximations	154
5.4	Consistency, order of accuracy and convergence	156
5.5	Stability and the Lax Equivalence Theorem	157
5.6	Calculating stability conditions	160
5.7	Practical (strict or strong) stability	166
5.8	Modified equation analysis	169
5.9	Conservation laws and the energy method of analysis	177
5.10	Summary of the theory	186
	<i>Bibliographic notes</i>	189
	<i>Exercises</i>	190

6	Linear second order elliptic equations in two dimensions	194
6.1	A model problem	194
6.2	Error analysis of the model problem	195
6.3	The general diffusion equation	197
6.4	Boundary conditions on a curved boundary	199
6.5	Error analysis using a maximum principle	203
6.6	Asymptotic error estimates	213
6.7	Variational formulation and the finite element method	218
6.8	Convection-diffusion problems	224
6.9	An example	228
	<i>Bibliographic notes</i>	231
	<i>Exercises</i>	232
7	Iterative solution of linear algebraic equations	235
7.1	Basic iterative schemes in explicit form	237
7.2	Matrix form of iteration methods and their convergence	239
7.3	Fourier analysis of convergence	244
7.4	Application to an example	248
7.5	Extensions and related iterative methods	250
7.6	The multigrid method	252
7.7	The conjugate gradient method	258
7.8	A numerical example: comparisons	261
	<i>Bibliographic notes</i>	263
	<i>Exercises</i>	263
	<i>References</i>	267
	<i>Index</i>	273

1

Introduction

Partial differential equations (PDEs) form the basis of very many mathematical models of physical, chemical and biological phenomena, and more recently their use has spread into economics, financial forecasting, image processing and other fields. To investigate the predictions of PDE models of such phenomena it is often necessary to approximate their solution numerically, commonly in combination with the analysis of simple special cases; while in some of the recent instances the numerical models play an almost independent role.

Let us consider the design of an aircraft wing as shown in Fig. 1.1, though several other examples would have served our purpose equally well – such as the prediction of the weather, the effectiveness of pollutant dispersal, the design of a jet engine or an internal combustion engine,

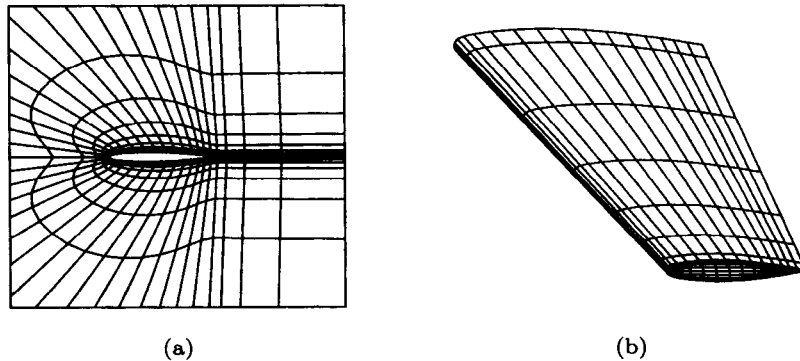


Fig. 1.1. (a) A typical (inviscid) computational mesh around an aerofoil cross-section; (b) a corresponding mesh on a wing surface.

the safety of a nuclear reactor, the exploration for and exploitation of oil, and so on.

In steady flight, two important design factors for a wing are the lift generated and the drag that is felt as a result of the flow of air past the wing. In calculating these quantities for a proposed design we know from boundary layer theory that, to a good approximation, there is a thin boundary layer near the wing surface where viscous forces are important and that outside this an inviscid flow can be assumed. Thus near the wing, which we will assume is locally flat, we can model the flow by

$$u \frac{\partial u}{\partial x} - \nu \frac{\partial^2 u}{\partial y^2} = (1/\rho) \frac{\partial p}{\partial x}, \quad (1.1)$$

where u is the flow velocity in the direction of the tangential co-ordinate x , y is the normal co-ordinate, ν is the viscosity, ρ is the density and p the pressure; we have here neglected the normal velocity. This is a typical *parabolic* equation for u with $(1/\rho)\partial p/\partial x$ treated as a forcing term.

Away from the wing, considered just as a two-dimensional cross-section, we can suppose the flow velocity to be inviscid and of the form $(u_\infty + u, v)$ where u and v are small compared with the flow speed at infinity, u_∞ in the x -direction. One can often assume that the flow is irrotational so that we have

$$\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} = 0; \quad (1.2a)$$

then combining the conservation laws for mass and the x -component of momentum, and retaining only first order quantities while assuming homentropic flow, we can deduce the simple model

$$(1 - M_\infty^2) \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 \quad (1.2b)$$

where M_∞ is the Mach number at infinity, $M_\infty = u_\infty/a_\infty$, and a_∞ is the sound speed.

Clearly when the flow is subsonic so that $M_\infty < 1$, the pair of equations (1.2a, b) are equivalent to the Cauchy-Riemann equations and the system is *elliptic*. On the other hand for supersonic flow where $M_\infty > 1$, the system is equivalent to the one-dimensional wave equation and the system is *hyperbolic*. Alternatively, if we operate on (1.2b) by $\partial/\partial x$ and eliminate v by operating on (1.2a) by $\partial/\partial y$, we either obtain an equivalent to Laplace's equation or the second order wave equation.

Thus from this one situation we have extracted the three basic types of partial differential equation: we could equally well have done so from the other problem examples mentioned at the beginning. We know from PDE theory that the analysis of these three types, what constitutes a well-posed problem, what boundary conditions should be imposed and the nature of possible solutions, all differ very markedly. This is also true of their numerical solution and analysis.

In this book we shall concentrate on model problems of these three types because their understanding is fundamental to that of many more complicated systems. We shall consider methods, mainly finite difference methods and closely related finite volume methods, which can be used for more practical, complicated problems, but can only be analysed as thoroughly as is necessary in simpler situations. In this way we will be able to develop a rigorous analytical theory of such phenomena as stability and convergence when finite difference meshes are refined. Similarly, we can study in detail the speed of convergence of iterative methods for solving the systems of algebraic equations generated by difference methods. And the results will be broadly applicable to practical situations where precise analysis is not possible.

Although our emphasis will be on these separate equation types, we must emphasise that in many practical situations they occur together, in a system of equations. An example, which arises in very many applications, is the Euler–Poisson system: in two space dimensions and time t , they involve the two components of velocity and the pressure already introduced; then, using the more compact notation ∂_t for $\partial/\partial t$ etc., they take the form

$$\begin{aligned}\partial_t u + u \partial_x u + v \partial_y u + \partial_x p &= 0 \\ \partial_t v + u \partial_x v + v \partial_y v + \partial_y p &= 0 \\ \partial_x^2 p + \partial_y^2 p &= 0.\end{aligned}\tag{1.3}$$

Solving this system requires the combination of two very different techniques: for the final elliptic equation for p one needs to use the techniques described in Chapters 6 and 7 to solve a large system of simultaneous algebraic equations; then its solution provides the driving force for the first two hyperbolic equations, which can generally be solved by marching forward in time using techniques described in Chapters 2 to 5. Such a model typically arises when flow speeds are much lower than in aerodynamics, such as flow in a porous medium, like groundwater flow. The two procedures need to be closely integrated to be effective and efficient.

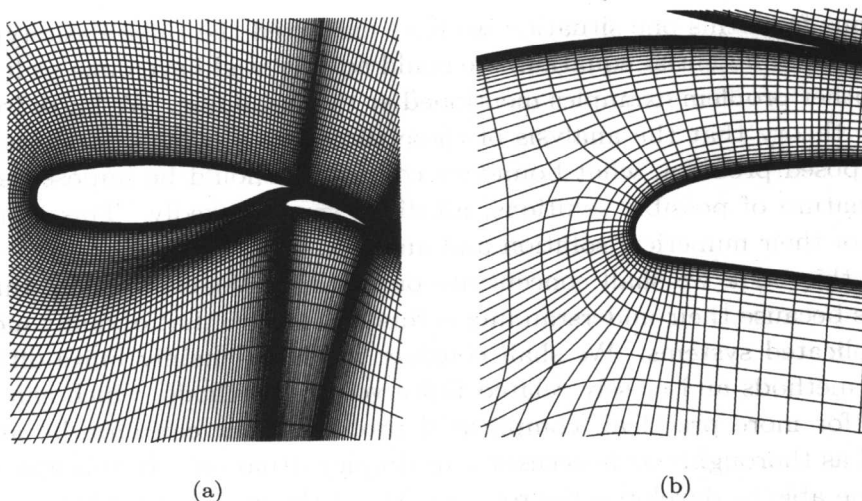


Fig. 1.2. A typical multi-aerofoil: (a) a general view; (b) a detail of the mesh that might be needed for a Navier-Stokes calculation. (Courtesy of DRA, Farnborough.)

Returning to our wing design example, however, it will be as well to mention some of the practical complications that may arise. For a civil aircraft most consideration can be given to its behaviour in steady flight at its design speed; but, especially for a military aircraft, manoeuvrability is important, which means that the flow will be unsteady and the equations time-dependent. Then, even for subsonic flow, the equations corresponding to (1.2a, b) will be hyperbolic (in one time and two space variables), similar to but more complicated than the Euler-Poisson system (1.3). Greater geometric complexity must also be taken into account: the three-dimensional form of the wing must be taken into consideration particularly for the flow near the tip and the junction with the aircraft body; and at landing and take-off, the flaps are extended to give greater lift at the slower speeds, so in cross-section it may appear as in Fig. 1.2.

In addition, rather than the smooth flow regimes which we have so far implicitly assumed, one needs in practice to study such phenomena as shocks, vortex sheets, turbulent wakes and their interactions. Developments of the methods we shall study are used to model all these situations but such topics are well beyond the scope of this book. Present capabilities within the industry include the solution of approximations to the Reynolds-averaged Navier-Stokes equations for unsteady viscous flow around a complete aircraft, such as that shown in Fig. 1.3.