Statistical Mechanics

Part A: Equilibrium Techniques

Edited by
Bruce J. Berne

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Part A: Equilibrium Techniques

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Part 8: Applications

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Bruce L Berne

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Preface

The last decade has been marked by a rapid growth in statistical mechanics, especially in connection with the physics and chemistry of the fluid state. Our understanding in these areas has been considerably advanced and enriched by the discovery of new techniques and the sharpening of old techniques, ranging all the way from computer simulation to mode—mode coupling theories.

Statistical mechanics brings together under one roof a broad spectrum of mathematical techniques. The aim of these volumes is to provide a didactic treatment of those techniques that are most useful for the study of problems of current interest to theoretical chemists. The emphasis throughout is on the techniques themselves and not on reviewing the enormous literature in statistical mechanics. Each author was charged with the following task. Given N pages, (a) pose the problem, (b) present those aspects of the particular technique that clearly illustrate its internal workings, (c) apply the technique to the solution of several illustrative examples, and (d) write the chapter so that it will enable the reader to approach key citations to the literature intelligently.

These volumes are designed for graduate students and research workers in statistical mechanics. Nevertheless, because of the range of techniques and their general utility, they should be useful in other areas as well.

The choice of topics was dictated not only by the taste and interests of the editor, but also by the proviso that there did not already exist a didactic treatment in the literature. The topics fall rather neatly into two categories: equilibrium and nonequilibrium properties of fluids. Thus, this volume is devoted to equilibrium techniques and the companion volume to the nonequilibrium techniques.

This volume begins with a chapter on modern cluster methods in equilibrium statistical mechanics and shows how topological reduction can be used to renormalize bonds. A general discussion of renormalization methods is given and the formalism is applied to the study of polar gases, ionic solutions, perturbation theory of fluids, hydrogen-bonded fluids, and integral equations.

This chapter is followed by two chapters on the treatment of fluids with long-range forces. In particular, special techniques are presented for treating

Coulomb forces and permanent multipolar forces. These chapters draw on some of the concepts developed in the first chapter.

Much of what is known about the fluid state springs from Monte Carlo studies of model fluids. Chapter 4 is concerned with the application of Monte Carlo techniques to the study of fluids with short-range forces. Chapter 5 extends this method to the study of fluids with long-range forces and introduces some innovative techniques. The volume closes with a chapter on nucleation theory.

Bruce J. Berne

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MODERN THEORETICAL CHEMISTRY

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Part A

Edited by William H. Miller

Volume 2 DYNAMICS OF MOLECULAR COLLISIONS
Part B
Edited by William H. Miller

Volume 3 METHODS OF ELECTRONIC STRUCTURE THEORY
Edited by Henry F. Schaefer III

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Volume 5 STATISTICAL MECHANICS
Part A: Equilibrium Techniques
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STRUCTURE CALCULATION
Part B: Applications
Edited by Gerald A. Segal

Statistical Mechanics

Part A: Equilibrium Techniques

Edited by

Bruce J. Berne

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Cluster Methods in Equilibrium Statistical Mechanics of Fluids

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1. Introduction was all property of poil of groups derived and the statements

Cluster expansion methods have proven to be exceedingly useful and general for discussing the equilibrium structure and thermodynamic properties of fluids. They have been used to rederive older results, such as the virial expansions for gases⁽¹⁾ and the Debye-Hückel theory for ionic solutions, ⁽²⁾ and then extend them in a systematic way. They have been used to derive approximate results that had not been obtained in other ways, such as the Γ and γ expansions⁽³⁾ and the EXP approximation⁽⁴⁾ for classical fluids. Moreover, they provide a unifying language for discussing and comparing the variety of theories of liquids that have been proposed; for example, many integral equations such as the Percus-Yevick equation⁽⁵⁾ and the mean spherical model equation⁽⁶⁻⁸⁾ for the pair correlation function of a fluid can be regarded as devices for summing certain diagrams in a cluster series.

statismual-mechanical quantities related to fluid structure and thermodynamics are defined and are expressed in terms of objector expressions. Section 4 objections

In the early work of Mayer, (1,2) the theoretical manipulation of cluster expansions involved solving a number of rather complex combinatorial problems and was formidable enough to deter all but the most stalwart of theoreticians. A curious aspect of this work is that the route from a rather simple starting point (the formal definition of a canonical ensemble partition function) to an elegant and simple result (an expression for the free energy or pressure in terms of irreducible cluster integrals) required such a complicated path. It was as if all the complications introduced in the combinatorial analysis magically

disappeared at the end to give a simple answer. The work of Morita and Hiroike⁽⁹⁾ (see also Stell⁽¹⁰⁾) cleared up many of the difficulties by pointing out some of the formal properties of cluster series that are responsible for obtaining simple results. In the more recent work on cluster expansions, combinatorial problems play a very minor role; they are needed for defining cluster series and for evaluating cluster integrals but not for the formal manipulation of cluster series.

In this article we will deal with graphical methods for calculating the free energy and pair correlation function for homogeneous classical fluids consisting of atoms and rigid molecules away from the critical point. The formalism for fluids of molecules with orientation-dependent forces and for mixtures of molecules is no more complicated than for atomic liquids, and so we will include these features in the discussion. We will discuss some general properties of cluster expansions and then discuss techniques for reexpressing cluster series in forms that lead to computationally tractable and accurate approximations for fluids of various types.

Section 2 defines the graph-theoretic terms we will use. In Section 3, the statistical-mechanical quantities related to fluid structure and thermodynamics are defined and are expressed in terms of cluster expansions. Section 4 contains a discussion of "topological reduction," which is the basic theoretic tool for manipulating cluster expansions. Section 5 discusses the various strategies used in applying cluster theory to particular problems. In Section 6, some examples are discussed to illustrate the various strategies. Section 7 contains some suggestions for further reading.

Two other chapters in this volume, namely, Chapter 2 by Stell and Chapter 3 by Friedman and Dale also contain discussions of particular applications of cluster expansion theory.

2. Graph-Theoretic Definitions

In this section we will define some of the important terms associated with the graphs in cluster expansion theory.

2.1. Definition of a Graph and the Ideas of Topological Equivalence and Connectivity

A graph is a collection of points and bonds that connect these points. In a pictorial representation of a graph, a point is drawn as a small circle and a bond as a line from one circle to another. (In some applications, it is convenient to define bonds that connect three or more points but we will not discuss this possibility here.) There are two different kinds of points, namely, root points

and field points. Root points are represented as open circles and field points as filled circles. (In some applications there is more than one type of root and field point, but we will not consider this here.) In some applications, there is only one type of bond, which is usually represented as a solid line. In other applications, there is more than one type of bond, and so we might use a solid line to denote one type of bond, a dashed line to represent the second type, etc. (For present applications, bonds have no intrinsic direction.) For theoretical manipulations in cluster theory, it is usually most convenient to deal with graphs in which the root points are labeled but the field points are unlabeled. That is, each root point is given a label, usually a number, and no two roots have the same label. Unless otherwise specified, all graphs discussed in this article are of this type. Some examples of graphs are given in the various figures.

We now wish to consider the idea of whether two graphs are topologically different or topologically equivalent. We will first discuss this idea informally and then give a more precise definition.

First of all, the question of whether two graphs are topologically equivalent or different arises only when they have the same number of root points and those roots have the same set of labels. Roughly speaking, two graphs are topologically equivalent if one can be "distorted" in such a way that it is superimposable on the other. To envisage the types of "distortions" that are allowed, imagine the points as buttons and the bonds as flexible rubber bands that connect the buttons. The buttons may be freely moved around on the graph, and when they move they drag along the ends of any rubber bands that are attached to them. The rubber bands can freely pass through each other, however. Some examples of topologically equivalent graphs are given in Fig. 1.

To define this concept more precisely, we need to define a labeled graph. A labeled graph has exactly the same definition as that given above for a graph, except that the field points (if any) as well as the root points have labels, usually numbers, and no two points have the same label. It is easy to define topological equivalence for labeled graphs. Two labeled graphs, which have the same number of root points, the same set of labels on the root points, the same

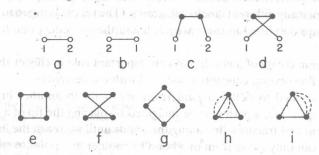


Fig. 1. Examples of topologically equivalent graphs. Graphs a and b are equivalent; as are c and d; e, f, and g; h and i. (The symmetry numbers of graphs a-i are 1, 1, 1, 1, 8, 8, 8, 2, and 2, respectively.)