Thermal Vibrations in Crystallography

B. T. M. Willis

A. W. Pryor

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B. T. M. Willis

Atomic Energy Research Establishment, Harwell

A. W. Pryor

Australian Atomic Energy Commission, Lucas Heights, and Visiting Fellow, School of Mathematics and Physics Macquarie University

Published by the Syndics of the Cambridge University Press Bentley House, 200 Euston Road, London NW1 2DB American Branch: 32 East 57th Street, New York, N.Y.10022

© Cambridge University Press 1975

Library of Congress Catalogue Card Number: 73-94357

ISBN: 0 521 20447 X

First published 1975

Printed in Great Britain at the University Printing House, Cambridge (Brooke Crutchley, University Printer)

Dedicated to the memory of Barrie Dawson

Preface

In the past fifteen years, the work of solid-state physicists has enormously increased our knowledge of the vibrational properties of atoms in solids. This book describes the impact of this new understanding of lattice dynamics on crystallography. Whether or not a crystallographer is concerned with the vibrational properties of a crystal, his experimental measurements always relate to atoms in thermal motion and so he must know how to account for the effect of this motion in extracting the maximum amount of information from his data. Much of our discussion is in terms of simple models of the type used in solid-state physics, but we show how ideas originating from these models may be carried over into the interpretation of the more complicated systems studied by the crystallographer. We have restricted ourselves to general points which must always be considered in carrying out diffraction work, and for this reason special topics, such as the connection between lattice vibrations and structural phase transitions, are omitted.

The book is written in three parts, of three chapters each. The first part deals with the basic theory of lattice dynamics. The traditional approach to the theory begins with the vibrations of a linear chain of point masses. We begin instead with the elements of the theory of the vibrations of an isolated molecule, and then proceed from a molecule to the Born-von Kármán treatment of a crystal by introducing the twin concepts of the Brillouin zone and of running waves in a finite but unbounded medium. The treatment of the vibrations of rigid molecules in a molecular crystal follows as a straightforward extension of the Born-von Kármán procedure. Our approach has the merit of bringing the reader immediately into contact with the real world of molecules and three-dimensional crystals.

The second part covers the influence of lattice vibrations on the interpretation of Bragg intensities. This is a subject dating back to the early years of X-ray diffraction, and yet there are a number of comparatively recent developments which warrant a new discussion. These include anharmonic effects, which are often surprisingly large and may even give rise to 'forbidden' Bragg reflections, and the use of higher cumulants and of the rigid-molecule hypothesis in crystal-structure refinements. These developments have gone hand-in-hand with the modern quest for making

Preface | x

highly accurate intensity measurements - as required, for instance, in determining the bonding-electron distributions in solids.

The last part gives a brief survey of the thermal diffuse scattering (TDS) of X-rays and neutrons. The TDS of neutrons, usually referred to as inelastic neutron scattering, provides the principal source of experimental data on phonon dispersion relations of solids. We sketch the information about interatomic forces in crystals given by these dispersion relations; more detailed discussions are given in several new books on lattice dynamics, cited in the bibliography (p. 261). In the last chapter, we describe the correction of experimentally-measured Bragg intensities for the contribution of TDS.

The book is addressed primarily to those engaged in diffraction studies but the first part would be suitable reading for undergraduate students of physics or chemistry. Matrix methods are used throughout, and an account of the relevant matrix algebra is presented in an appendix.

Professor B. N. Brockhouse, Professor W. Cochran, Dr J. C. Speakman and Dr M. W. Thomas made many suggestions for improving various parts of our manuscript. The typing was carried out by Mrs Shirley Hinder. We wish to thank them all, and also those authors and publishers who have granted us permission to reproduce diagrams.

B.T.M.W.

Harwell and Sydney May 1973

Glossary of important symbols

The equation numbers in parentheses refer to the place in the text where the symbol is introduced

 \mathbf{a}_i (i = 1, 2, 3) base vectors of direct unit cell (1.1a). a, a_0 cell edge of cubic crystal. |A(i)|amplitude of excitation of the mode (i) of an isolated molecule (2.26) $|A(j\mathbf{q})|$ amplitude of excitation of the lattice mode (iq) of a crystal (3.17). \mathbf{b}_i (i = 1, 2, 3) base vectors of reciprocal unit cell (1.3). \boldsymbol{b}_{κ} coherent neutron scattering amplitude of atom κ (1.13). beoh } isotropic B factor ($\equiv 8\pi^2 \langle u^2 \rangle$) (4.72). B(K) 3×3 matrix representing mean-square displacement of atom κ (4.54). Binol(K) 6×6 matrix representing mean-square translationallibrational displacement of rigid molecule κ (6.4). elastic constants of cubic crystal (3.38). C11, C12, C44 specific heat. CU D $3n \times 3n$ dynamical matrix (3.7). D. dynamical matrix without mass-adjustment (3.10a). Dinol $6n \times 6n$ dynamical matrix of molecular crystal (3.50c). Dinoi dynamical matrix of molecular crystal without massadjustment (3.50a). e(i)eigenvector describing pattern of atomic displacements in an isolated molecule due to the mode (i) (2.26). e(jq)eigenvector describing pattern of atomic displacements in a crystal due to the lattice mode (jq) (3.13). emol(ia) eigenvector describing pattern of molecular (translational and librational) displacements in molecular crystal due to the lattice mode $(i\mathbf{q})$ (3.51). $e(\kappa | j)$ polarisation vector of atom κ when vibrating in the mode (i) of an isolated molecule. The assemblage of n polarisation vectors ($\kappa = 1, 2, ..., n$) constitutes the 3*n*-vector $\mathbf{e}(j)$ (2.31).

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Glossary | xii
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e(\kappa | iq)
                   polarisation vector of atom \kappa when vibrating in the mode
                   (i\mathbf{q}) of a crystal. The assemblage of n polarisation vectors
                    (\kappa = 1, 2, ..., n) constitutes the 3n-vector e(i\mathbf{q}) (3.13).
e^{\text{mol}}(\kappa | jq)
                   polarisation vector of molecule \kappa when vibrating in the
                   lattice mode (ja) of a crystal. The assemblage of n polarisa-
                   tion vectors (\kappa = 1, 2, ..., n) constitutes the 6n-vector
                    e^{\text{mol}}(i\mathbf{q}).
                   average energy of mode (i) of isolated molecule
E_i
                    (2.40).
E_i(\mathbf{q})
                   average energy of lattice mode (iq) of crystal (3.17).
E_0
                   energy of incident X-rays (neutrons);
                    or integrated intensity of Bragg scattering.
\boldsymbol{E}
                   energy of scattered X-rays (neutrons).
F(\mathbf{O})
F(\mathbf{H})
                   structure factor for Bragg scattering (1.11).
F_{hkl}
f_{\kappa}(\mathbf{Q})
                   X-ray scattering factor of atom \kappa (1.12).
g(v)
                   frequency distribution function (1.20).
g(\omega)
                   frequency distribution function for translational modes
g_{T}(\omega)
                   (6.15a).
                   frequency distribution function for librational modes
g_{\rm L}(\omega)
                   (6.15b).
G(j\mathbf{q})
                   structure factor for first-order (one-phonon) scattering
                   associated with lattice mode (iq) (7.4).
                   3×3 matrix representing metric tensor with elements
g
                   g_{ii} = \mathbf{a}_i \cdot \mathbf{a}_i
g^{-1}
                   3 \times 3 matrix with elements g_{ij}^{-1} = \mathbf{b}_i \cdot \mathbf{b}_i.
                   vector of reciprocal lattice (\equiv h_1 b_1 + h_2 b_2 + h_3 b_3) (1.4).
H
h_1h_2h_3
                   Miller indices.
hkl
h
                   Miller index;
                   or Planck's constant.
ħ
                   Planck's constant \div 2\pi.
1
                   unit matrix (2.25).
(\kappa)
                   3 \times 3 moment-of-inertia matrix-of-molecule \kappa (3.49b).
                   intensity of Bragg scattering.
I_0
I_1, I_2, \dots
                   intensity of first-order (one-phonon), second-order (two-
                   phonon)...scattering.
I_1(j\mathbf{q})
                   intensity of first-order scattering associated with single
                   lattice mode (jq) (4.22).
                   label for modes of vibration of isolated molecule con-
                   taining n atoms (j = 1, 2, ..., 3n) (2.26);
```

	or label for branch of dispersion relations of crystal con-
	taining n atoms per unit cell $(j = 1, 2,, 3n)$ (3.13).
(<i>j</i> q)	lattice mode of vibration belonging to jth branch of
	dispersion relations and with wave vector q.
\mathbf{k}_{0}	wave vector of incident radiation $(1.8a)$.
k	wave vector of scattered radiation (1.8a).
k	Miller index;
2	or wave number of scattered radiation.
$k_{ m B}$	Boltzmann's constant (1.19).
1	Miller index;
1	or label for unit cell $(l = 1, 2,, N)$.
L(K)	3×3 libration matrix of rigid molecule κ (6.5b).
$m(\kappa)$	mass of atom κ (2.11);
	or mass of molecule κ (3.49).
$\mathbf{m}(\kappa)$	3×3 diagonal matrix of masses $m(\kappa)$ (2.17).
m	$3n \times 3n$ diagonal mass matrix (2.16).
$\mathbf{m}^{\mathrm{mol}}$	$6n \times 6n$ matrix of molecular masses and moments of inertia
	(3.49b).
M_{o}	$3n \times 3n$ dynamical matrix for the isolated molecule formed
0	by assembling $n \times n$ force-constant matrices (2.18).
M	mass-adjusted M ₀ -matrix (2.23).
$m_{\rm n}$	neutron mass.
n	number of atoms in isolated molecule (2.2);
	or number of atoms in (primitive) unit cell (3.13);
	or number of molecules in (primitive) unit cell.
$\langle n \rangle$	average quantum number of a mode of vibration of a
. ,	molecule or crystal (2.43).
N	number of unit cells in crystal.
N_{0}	number of atoms in crystal $(= nN)$.
$p_{\kappa}(\mathbf{u})$	probability density function of atom κ (4.45).
Q	scattering vector (1.8).
\overline{Q}	magnitude of scattering vector (= $4\pi \sin \theta_B/\lambda$ for Bragg
~	scattering).
q	wave vector of lattice mode of vibration.
q	wave number of lattice mode.
$\mathbf{r}(\kappa)$	equilibrium position of atom κ with respect to origin of
	unit cell (1.12).
$\mathbf{r}(\kappa l)$	equilibrium position of atom κ in cell l with respect to
	crystal origin (3.10).
$\mathbf{r}(\kappa\alpha)$	equilibrium position of atom α in molecule κ of molecular
	crystal (6.7).
$\mathbf{r}(l)$	vector defining position of origin of cell <i>l</i> in crystal (4.2).
$S(\kappa)$	3×3 translation-libration matrix of rigid molecule κ (6.5c).
-(-)	

Glossary	xiv

Glossary XIV	
T	absolute temperature (1.19).
T_{κ}	temperature factor of atom κ (4.1).
$T(\kappa)$	3×3 translation matrix of rigid molecule κ (6.5a).
u(K)	instantaneous displacement from its equilibrium position
	of atom κ in isolated molecule (2.1);
	or instantaneous translational displacement from its
	equilibrium posttion of a rigid molecule κ (6.6).
$\mathbf{u}(\kappa l)$	instantaneous displacement from its equilibrium position
	of atom κ in cell l of crystal (3.19).
$\langle u^2 \rangle$	mean-square displacement from equilibrium position.
$\mathbf{U_0}$	$3n \times 1$ column matrix representing amplitudes of displace-
	ment of n atoms in isolated molecule due to all modes of
	vibration (2.15);
	or $3n \times 1$ column matrix representing amplitudes of dis-
	placement of n atoms in unit cell due to all modes of
	vibration (3.8).
$\mathbf{U_0}(\mathbf{jq})$	$3n \times 1$ column matrix representing amplitudes of displace-
	ment of n atoms in unit cell due to single lattice mode (jq)
	(3.16).
U	mass-adjusted U_0 -matrix (2.22).
$U(j\mathbf{q})$	mass-adjusted $U_0(j\mathbf{q})$ -matrix.
$U(\kappa)$	amplitude of displacement of atom κ from all modes of
	vibration of molecule (2.12).
$\mathbf{U}(\kappa j)$	amplitude of displacement of atom κ due to mode of
	vibration (j) of molecule.
$\mathbf{U}(\kappa j\mathbf{q})$	amplitude of displacement of atom κ due to mode of
	vibration ($j\mathbf{q}$) of crystal (3.2);
	or amplitude of translational displacement of molecule κ
• mal	due to mode of vibration (jq) of crystal (6.5).
$\mathbf{U}_0^{\mathrm{mol}}$	$6n \times 1$ column matrix representing amplitudes of transla-
	tional and librational displacements of n molecules in the
$\mathbf{U}_0^{\mathrm{mol}}(j\mathbf{q})$	unit cell of a molecular crystal.
$O_0^{-1}(\mathbf{q})$	6n×1 column matrix representing amplitudes of transla-
	tional and librational displacements of n molecules in unit cell due to single lattice mode ($j\mathbf{q}$).
Umol	mass-adjusted $\mathbf{U}_0^{\text{mol}}$ matrix.
-	mass-adjusted $\mathbf{U}_0^{\text{matrix}}$.
	mass-adjusted 6×1 column matrix representing amplitude
(1)4)	of translational and librational displacement of molecule
	when excited by lattice mode (jq) .
	volume of direct unit cell (1.1 <i>a</i>).
**	volume of reciprocal unit cell (1.5).
$V_b = V$	volume of crystal (1.14);
•	Totalie of crystal (1.17),

```
or potential energy of molecule (2.5);
                     or potential energy of crystal (3.6).
 V_{\nu}(\mathbf{u})
                     single-atom potential of atom \kappa (5.15).
                     velocity of longitudinal elastic wave.
 v_{\mathbf{L}}
                     velocity of transverse elastic wave.
 v_{\mathbf{T}}
                    average sound velocity (4.103).
 v_a
                    incident neutron velocity (4.28).
 vo
                    scattered neutron velocity (4.28).
 W.
                    exponent of temperature factor (i.e. T_{\kappa} \equiv \exp(-W_{\kappa}))
                    (4.9).
 \alpha_i (i = 1, 2, 3) interaxial angles of direct unit cell (1.1b).
 \alpha (= 1, 2, 3)
                    label for three Cartesian axes xyz (2.5);
                    or label for atom in a molecule (6.14).
\alpha, \beta, \gamma
                    elements of force-constant matrix (3.29).
\alpha, \beta, \gamma, \delta
                    force constants representing coefficients in expansion of
                    single-atom potential V(\mathbf{u}) in powers of \mathbf{u} (5.15).
\alpha_i(\mathbf{q})
                    angle between propagation direction and polarisation
                    direction of lattice mode (jq) (7.8).
\beta_i (i = 1, 2, 3) interaxial angles of reciprocal unit cell (1.6).
                    3×3 matrix representing anisotropic temperature factor
B(K)
                    of atom \kappa (4.64).
                    Grüneisen's constant (5.1).
\gamma_{G}
                    Kronecker delta: \delta_{ij} = 1 for i = j
\delta_{ii}
                                              = 0 for i \neq j (1.3).
δ
                    Dirac delta function.
ζ
                    wave-vector coordinate (3.31).
2\theta
                    scattering angle (= 2\theta_{\rm B} for Bragg scattering).
\theta_{\mathrm{B}}
                    Bragg angle.
\Theta_{\mathbf{D}}
                    Debye temperature.
\Theta(\kappa | j\mathbf{q})
                    amplitude of librational displacement of molecule k due
                    to mode of vibration (j\mathbf{q}) of crystal (6.5).
\theta(\kappa)
                    instantaneous librational displacement of rigid molecule
                    \kappa from its equilibrium position (6.6).
                    label for atom in isolated molecule (2.1);
K
                    or label for atom in (primitive) unit cell (3.1);
                    or label for molecule in (primitive) unit cell (6.3).
λ
                    wavelength of X-rays or neutrons (1.8b).
                    frequency (1.19).
                    density.
\Phi(\kappa\kappa')
                    3 × 3 force-constant matrix describing interaction between
                    atoms \kappa and \kappa' in isolated molecule (2.4).
                    3 \times 3 force-constant matrix describing interaction in a
\Phi\begin{pmatrix} \kappa & \kappa' \\ l & l' \end{pmatrix}
                    crystal between atom \kappa in cell l and atom \kappa' in cell l' (3.5).
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Glossary | xvi

 $\omega_{\rm D}$

6×6 force- and torque-constant matrix describing inter- $\Phi^{\text{mol}}\binom{\kappa}{l}$ action in a molecular crystal between molecule κ in cell land molecule κ' in cell l'. volume coefficient of expansion (5.1). χ circular frequency ($\equiv 2\pi\nu$). **(1)** circular frequency of lattice mode (jq) (3.15). $\omega_i(\mathbf{q})$ frequency of longitudinal mode. ω_{L} frequency of transverse mode. $\omega_{\mathbf{T}}$ Debye cut-off frequency (4.106).

Contents

Preface	e ,	page ix
Glossa	ry of important symbols	xi
Part I.	Theory of lattice dynamics	1
Chapte	er 1. Introduction and basic concepts	3
1.1.	Scope of text	3
1.2.	Nomenclature	4
1.3.	Some basic concepts of diffraction theory	4
	1.3.1. Direct lattice and reciprocal lattice	5
	1.3.2 Bragg scattering. The Ewald sphere	6
	1.3.3. The Fourier transform in diffraction theory	9
1.4.	•	11
	1.4.1. The Einstein model	11
Chapte	er 2. The internal modes of vibration of molecules	14
2.1.	Classical treatment of molecular vibrations	14
	2.1.1. Atomic displacements	14
	2.1.2. Definition of interatomic forces	15
	2.1.3. Equations of motion	18
2.2.	Energy and quantisation of normal modes	22
2.3.	Example: the CO ₂ molecule	25
2.4.	Example: monatomic linear chain	29
Chapte	er 3. The lattice modes of vibration of crystals	34
3.1.	Born-von Kármán theory: travelling waves in a finite but	8 0
	unbounded crystal	34
	3.1.1. The Brillouin zone	36
	3.1.2. Equations of motion	39
	3.1.3. Quantisation of normal modes. Phonons	43
3.2.	The force-constant matrices: symmetry restrictions	44
3.3.	The lattice vibrations of solids with simple crystal structur	res 51
	3.3.1. The cubic close-packed structure	51
	3.3.2. The rock-salt structure	

10000	
Contents	
Contents	vı
Contents	7.4

3.4. 3.5.	The lattice modes of molecular crystals page The calculation of integral vibrational properties	e 67 77
Part II.	Effect of temperature on Bragg intensities	79
Chapter 4.1.	4. The atomic temperature factor (Debye-Waller factor) General expression for intensity of X-ray scattering by a	81
	crystal in thermal motion	82
4.2.	Intensity of thermal neutron scattering	88
4.3.	The temperature factor and the probability density function	92
4.4.	The temperature factor in crystal-structure analysis	95
	4.4.1. The mean-square displacement matrix	95
	4.4.2. The ellipsoid of thermal vibration	97
	4.4.3. The anisotropic temperature-factor coefficients β_{ij}	99
	4.4.4. The equivalent isotropic mean-square displacement	101
	4.4.5. Symmetry restrictions on the anisotropic temperature	
	factor	102
	4.4.6. An example: sodium nitrate	110
	4.4.7. Thermal motion and interatomic distances	115
4.5.	Lattice-dynamical considerations	117
	4.5.1. Monatomic cubic crystals	118
	4.5.2. Diatomic cubic and close-packed hexagonal crystals	119
	4.5.3. Born S -matrix	120
4.6.	The Debye approximation	122
	4,6.1. Modified Debye model	128
4.7.	Experimental measurements of atomic temperature factors:	
	comparison with theoretical values	130
	4.7.1. Face-centred cubic metals. Variation with temperature	
	of mean-square displacement	134
	4.7.2. Potassium chloride and sodium chloride	140
Chapter	5 Anharmonic contributions to the atomic temperature factor	142
5.1.	Quasi-harmonic theory of atomic temperature factor	142
5.2.	Formal lattice-dynamical treatment of anharmonic Debye-	
	Waller factor	144
5.3.	Isolated-atom treatment of anharmonicity	146
	5.3.1. The isolated-atom potential for an atom occupying	
	a site of cubic symmetry	147
	5.3.2. Temperature factor for an atom at a site of cubic	
	symmetry	150

	(*)	Conte	ents vii
5.4.	Dawso	on structure-factor formalism	page 151
	5.4.1.	Cubic close-packed structure	153
	5.4.2.	Rock-salt structure	154
	5.4.3.	Diamond structure. 'Forbidden' reflections	154
	5.4.4.	Fluorite structure	157
5.5.	Experi	mental results	159
	5.5.1.	Aluminium	159
	5.5.2.	Potassium chloride and sodium chloride	160
	5.5.3.	Silicon	165
	5.5.4.	Crystals with fluorite structure	168
5.6.	Higher	order-cumulants treatment of anharmonicity	173
		temperature factor of molecular crystals	175
6.1.	100	gid-body model	175
6.2.		ption of mean thermal displacement of a rigid	
	mole		176
	6.2.1.	Atomic and molecular displacements. Relation	8
		between atomic B-matrix and molecular T, L,	
		matrices	180
		Symmetry restrictions on elements of T, L and S	182
6.3.	-	l-structure refinements based on the rigid-body	
	mod		186
6.4.		bution of internal modes to the temperature factor	
6.5.		ethylenetetramine (HMT)	193
6.6.		ength correction from librational motion	197
6.7.	Higher	cumulants and librational motion	198
Part III	. Therm	nal diffuse scattering	205
Chapter	7. The	rmal diffuse scattering of X-rays	207
7.1.	Introdu	uction	207
7.2.	Intensi	ty formulae	208
	7.2.1.	The long-range Q-dependence	209
	7.2.2.	The frequency term	211
	7.2.3.	The structure-factor term	212
7.3.	Therma	al diffuse scattering experiments with X-rays	213
	7.3.1.	Determination of elastic constants: scattering from	m
		long-wavelength acoustic modes	213
	7.3.2.	Determination of phonon dispersion curves for	
		monatomic crystals	215
	7.3.3.	Molecular crystals	216

~	
Contents	V111

		Experimental difficulties	page	
	7.3.5.	General comments on X-ray experiments		217
Chapter	8. Inela	astic scattering of slow neutrons		219
8.1.	Conser	vation rules for coherent inelastic scattering		219
8.2.	The me	easurement of phonon dispersion curves using the	e	1
	triple	e-axis spectrometer		220
	8.2.1.	The constant-Q method		222
8.3.	Phonor	n dispersion relations: measurement and interpre	ta-	
	tion			225
	8.3.1.	Elements		225
	8.3.2.	Ionic solids		228
	8.3.3.	Molecular solids		231
Chapter	9. The	influence of thermal diffuse scattering on the		232
measure	ment of	f Bragg intensities		
9.1.	Introdu	uction		232
9.2.	Theory	of correction factor for X-rays		233
	9.2.1.	Single crystals		233
	9.2.2.	Powders		238
9.3.	Theory	of correction factor for neutrons		239
9.4.	Experi	mental tests of theory of TDS correction factor		241
Append	lix 1. A	summary of matrix algebra		245
Append	lix 2. Pr	robability density function of an isotropically		
	vibratii	ng atom		259
B ibliogi	raphy			261
Referen	ces			263
Index				270

Part I. Theory of lattice dynamics