

Major American Universities Ph. D.
Qualifying Questions and Solutions

Problems and Solutions on Thermodynamics and Statistical Mechanics

Compiled by:
The Physics Coaching Class
University of Science and
Technology of China

Edited by:
Yung-Kuo Lim

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PROBLEMS AND SOLUTIONS ON THERMODYNAMICS AND
STATISTICAL MECHANICS**

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PREFACE

This series of physics problems and solutions which consists of seven parts – Mechanics, Electromagnetism, Optics, Atomic Nuclear and Particle Physics, Thermodynamics and Statistical Physics, Quantum Mechanics, Solid State Physics – contains a selection of 2550 problems from the graduate school entrance and qualifying examination papers of seven U.S. universities – California University Berkeley Campus, Columbia University, Chicago University, Massachusetts Institute of Technology, New York State University Buffalo Campus, Princeton University, Wisconsin University – as well as the CUSPEA and C.C. Ting's papers for selection of Chinese students for further studies in U.S.A. and their solutions which represent the effort of more than 70 Chinese physicists.

The series is remarkable for its comprehensive coverage. In each area the problems span a wide spectrum of topics while many problems overlap several areas. The problems themselves are remarkable for their versatility in applying the physical laws and principles, their up-to-date realistic situations, and their scanty demand on mathematical skills. Many of the problems involve order of magnitude calculations which one often requires in an experimental situation for estimating a quantity from a simple model. In short, the exercises blend together the objectives of enhancement of one's understanding of the physical principles and practical applicability.

The solutions as presented generally just provide a guidance to solving the problems rather than step by step manipulation and leave much to the student to work out for him/herself, of whom much is demanded of the basic knowledge in physics. Thus the series would provide an invaluable complement to the textbooks.

In editing no attempt has been made to unify the physical terms and symbols. Rather, they are left to the setters' and solvers' own preferences so as to reflect the realistic situation of the usage today.

The present volume for Thermodynamics and Statistical Physics comprises 367 problems.

Lim Yung Kuo
Editor

INTRODUCTION

Solving problems in school work is the exercise of mental faculties, and examination problems are usually picked from the problems in school work. Working out problems is a necessary and important aspect in studies of Physics.

The *Major American University Ph.D. Qualifying Questions and Solutions* is a series of books which consists of seven volumes. The subjects of each volume and their respective referees (in parentheses) are as follows:

1. Mechanics (Qiang Yan-qi, Gu En-pu, Cheng Jia-fu, Li Ze-hua, Yang De-tian)
2. Electromagnetism (Zhao Shu-ping, You Jun-han, Zhu Jun-jie)
3. Optics (Bai Gui-ru, Guo Guang-can)
4. Atomic, Nuclear and Particle Physics (Jin Huai-cheng, Yang Bao-zhong, Fan Yang-mei)
5. Thermodynamics and Statistical Physics (Zheng Jiu-ren)
6. Quantum Mechanics (Zhang You-de, Zhu Dong-pei, Fan Hong-yi)
7. Solid Physics and Comprehensive Topics (Zhang Jia-lu, Zhou You-yuan, Zhang Shi-ling)

The books cover almost all aspects of University Physics and contain 2550 problems, most of which are solved in detail.

These problems are carefully chosen from 3100 problems, some of which came from the China-U.S. Physics Examination and Application Program, others were selected from the Ph.D. Qualifying Examination on Experimental High Energy Physics, sponsored by Chao Chong Ting. The rest came from the graduate entrance examination questions of seven famous American universities during the last decade; they are: Columbia University, University of California at Berkeley, Massachusetts Institute of Technology, University of Wisconsin, University of Chicago, Princeton University and State University of New York, Buffalo.

In general, examination problems in physics in American universities do not involve too much Mathematics; however, they are to some extent characterized by the following three aspects: some problems involving various frontier subjects and overlapping domains of science are selected by professors directly from their own research work and show a "modern style", some problems involve a wider field and require a quick mind to analyse, and the methods used for solving the other problems are simple and practical which shows a full "touch of physics". From these, we think that these problems as a whole embody, to some extent, the characteristics of Ameri-

can science and culture and the features of the way of thinking of American education.

Just so, we believe it is worthwhile to collect and solve these problems and introduce them to students and teachers, even though the work is strenuous. About a hundred teachers and graduate students took part in this time-consuming job.

There are 367 problems in this volume which is divided into two parts: part I consists of 159 problems on Thermodynamics, part II consists of 208 problems on Statistical physics. Each part contains five sections.

The depth of knowledge involved in solving these problems is not beyond the contents of common textbooks on Thermodynamics and Statistical Physics used in colleges and universities in China, although the range of the knowledge and the techniques needed in solving some of these problems go beyond what we are familiar with. Furthermore, some new scientific research results are introduced into problems in school work, that will benefit not only the study of established theories and knowledge, but also the combination of teaching and research work by enlivening academic thoughts and making minds more active.

The people who contributed to solving these problems of this volume are Feng Ping, Wang Hai-da, Yao De-min and Jia Yun-fa. Associate professor Zheng Jiu-ren and Mr. Zheng Xin are referees of English of this volume.

15 October 1989

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PART I

THERMODYNAMICS

1. THERMODYNAMIC STATES AND THE FIRST LAW (1001-1030)

1001

Describe briefly the basic principle of the following instruments for making temperature measurements and state in one sentence the special usefulness of each instrument: constant-volume gas thermometer, thermocouple, thermistor.

(Wisconsin)

Solution:

Constant-volume gas thermometer: It is made according to the principle that the pressure of a gas changes with its temperature while its volume is kept constant. It can approximately be used as an ideal gas thermometer.

Thermocouple thermometer: It is made according to the principle that thermoelectric motive force changes with temperature. The relation between the thermoelectric motive force and the temperature is

$$\epsilon = a + bt + ct^2 + dt^3 ,$$

where ϵ is the electric motive force, t is the difference of temperatures of the two junctions, a, b, c and d are constants. The range of measurement of the thermocouple is very wide, from -200°C to 1600°C . It is used as a practical standard thermometer in the range from 630.74°C to 1064.43°C .

Thermistor thermometer: We measure temperature by measuring the resistance of a metal. The precision of a thermistor made of pure platinum is very good, and its range of measurement is very wide, so it is usually used as a standard thermometer in the range from 13.81K to 903.89K.

1002

Describe briefly three different instruments that can be used for the accurate measurement of temperature and state roughly the temperature range in which they are useful and one important advantage of each instrument. Include at least one instrument that is capable of measuring temperatures down to 1K.

(Wisconsin)

Solution:

1. *Magnetic thermometer:* Its principle is Curie's law $\chi = C/T$, where χ is the susceptibility of the paramagnetic substance used, T is its absolute temperature and C is a constant. Its advantage is that it can measure temperatures below 1K.

2. *Optical pyrometer:* It is based on the principle that we can find the temperature of a hot body by measuring the energy radiated from it, using the formula of radiation. While taking measurements, it does not come into direct contact with the measured body. Therefore, it is usually used to measure the temperatures of celestial bodies.

3. *Vapor pressure thermometer:* It is a kind of thermometer used to measure low temperatures. Its principle is as follows. There exists a definite relation between the saturation vapor pressure of a chemically pure material and its boiling point. If this relation is known, we can determine temperature by measuring vapor pressure. It can measure temperatures greater than 14K, and is the thermometer usually used to measure low temperatures.

1003

A bimetallic strip of total thickness x is straight at temperature T . What is the radius of curvature of the strip, R , when it is heated to temperature $T + \Delta T$? The coefficients of linear expansion of the two metals are α_1 and α_2 , respectively, with $\alpha_2 > \alpha_1$. You may assume that each metal has thickness $x/2$, and you may assume that $x \ll R$.

(Wisconsin)

Solution:

We assume that the initial length is l_0 . After heating, the lengths of the mid-lines of the two metallic strips are respectively

$$l_1 = l_0(1 + \alpha_1 \Delta T), \quad (1)$$

$$l_2 = l_0(1 + \alpha_2 \Delta T). \quad (2)$$

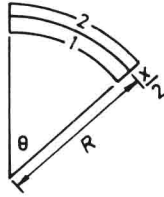


Fig. 1.1.

Assuming that the radius of curvature is R , the subtending angle of the strip is θ , and the change of thickness is negligible, we have

$$l_2 = \left(R + \frac{x}{4}\right) \theta, \quad l_1 = \left(R - \frac{x}{4}\right) \theta,$$

$$l_2 - l_1 = \frac{x}{2} \theta = \frac{x}{2} \frac{l_1 + l_2}{2R} = \frac{x l_0}{4R} [2 + (\alpha_1 + \alpha_2) \Delta T]. \quad (3)$$

From (1) and (2) we obtain

$$l_2 - l_1 = l_0 \Delta T (\alpha_2 - \alpha_1), \quad (4)$$

(3) and (4) then give

$$R = \frac{x}{4} \frac{[2 + (\alpha_1 + \alpha_2) \Delta T]}{(\alpha_2 - \alpha_1) \Delta T}.$$

1004

An ideal gas is originally confined to a volume V_1 in an insulated container of volume $V_1 + V_2$. The remainder of the container is evacuated. The partition is then removed and the gas expands to fill the entire container. If the initial temperature of the gas was T , what is the final temperature? Justify your answer.

(Wisconsin)

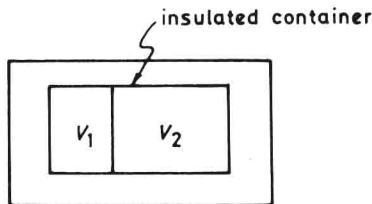


Fig. 1.2.

Solution:

This is a process of adiabatic free expansion of an ideal gas. The internal energy does not change; thus the temperature does not change, that is, the final temperature is still T .

1005

An insulated chamber is divided into two halves of volumes. The left half contains an ideal gas at temperature T_0 and the right half is evacuated. A small hole is opened between the two halves, allowing the gas to flow through, and the system comes to equilibrium. No heat is exchanged with the walls. Find the final temperature of the system.

(Columbia)

Solution:

After a hole has been opened, the gas flows continuously to the right side and reaches equilibrium finally. During the process, internal energy of the system E is unchanged. Since E depends on the temperature T only for an ideal gas, the equilibrium temperature is still T_0 .

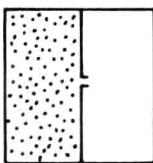


Fig. 1.3.

1006

Define heat capacity C_v and calculate from the first principle the numerical value (in calories/ $^{\circ}\text{C}$) for a copper penny in your pocket, using your best physical knowledge or estimate of the needed parameters.

(UC, Berkeley)

Solution:

$C_v = (dQ/dT)_v$. The atomic number of copper is 64 and a copper penny is about 32 g, i.e., 0.5 mol. Thus $C_v = 0.5 \times 3R = 13 \text{ J/K}$.

1007

Specific heat of granite may be: 0.02, 0.2, 20, 2000 cal/g·K.

(Columbia)

Solution:

The main component of granite is CaCO_3 ; its molecular weight is 100. The specific heat is $C = 3R/100 = 0.25 \text{ cal/g}\cdot\text{K}$. Thus the best answer is $0.2 \text{ cal/g}\cdot\text{K}$.

1008

The figure below shows an apparatus for the determination of C_p/C_v for a gas, according to the method of Clement and Desormes. A bottle G , of reasonable capacity (say a few litres), is fitted with a tap H , and a manometer M . The difference in pressure between the inside and the outside can thus be determined by observation of the difference h in heights of the two columns in the manometer. The bottle is filled with the gas to be investigated, at a very slight excess pressure over the outside atmospheric pressure. The bottle is left in peace (with the tap closed) until the temperature of the gas in the bottle is the same as the outside temperature in the room. Let the reading of the manometer be h_i . The tap H is then opened for a very short time, just sufficient for the internal pressure to become equal to the atmospheric pressure (in which case the manometer reads $h = 0$). With the tap closed the bottle is left in peace for a while, until the inside temperature has become equal to the outside temperature. Let the final reading of the manometer be h . From the values of h_i and h_f it is possible to find C_p/C_v . (a) Derive an expression for C_p/C_v in terms of h_i and h_f in the above experiment. (b) Suppose that the gas in question is oxygen. What is your theoretical prediction for C_p/C_v at 20°C , within the framework of statistical mechanics?

(UC, Berkeley)

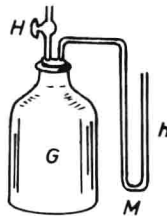


Fig. 1.4.

Solution:

(a) The equation of state of ideal gas is $pV = nkT$. Since the initial and final T, V of the gas in the bottle are the same, we have $p_f/p_i = n_f/n_i$.

Meanwhile, $n_f/n_i = V/V'$, where V' is the volume when the initial gas in the bottle expands adiabatically to pressure p_0 . Therefore

$$\frac{V}{V'} = \left(\frac{p_0}{p_i}\right)^{\frac{1}{\gamma}}, \quad \frac{p_f}{p_i} = \left(\frac{p_0}{p_i}\right)^{\frac{1}{\gamma}},$$

$$\gamma = \frac{\ln \frac{p_i}{p_0}}{\ln \frac{p_i}{p_f}} = \frac{\ln \left(1 + \frac{h_i}{h_0}\right)}{\ln \left(1 + \frac{h_i}{h_0}\right) - \ln \left(1 + \frac{h_f}{h_0}\right)}.$$

Since $h_i/h_0 \ll 1$ and $h_f/h_0 \ll 1$, we have $\gamma = h_i/(h_i - h_f)$.

(b) Oxygen consists of diatomic molecules. When $t = 20^\circ\text{C}$, only the translational and rotational motions of the molecules contribute to the specific heat. Therefore

$$C_v = \frac{5R}{2}, \quad C_p = \frac{7R}{2}, \quad \gamma = \frac{7}{5}.$$

1009

(a) Starting with the first law of thermodynamics and the definitions of c_p and c_v , show that

$$c_p - c_v = \left[p + \left(\frac{\partial U}{\partial V}\right)_T \right] \left(\frac{\partial V}{\partial T}\right)_p$$

where c_p and c_v are the specific heat capacities per mole at constant pressure and volume, respectively, and U and V are energy and volume of one mole.

(b) Use the above results plus the expression

$$p + \left(\frac{\partial U}{\partial V}\right)_T = T \left(\frac{\partial p}{\partial T}\right)_V$$

to find $c_p - c_v$ for a Van der Waals gas

$$\left(p + \frac{a}{V^2}\right)(V - b) = RT.$$

Use that result to show that as $V \rightarrow \infty$ at constant p , you obtain the ideal gas result for $c_p - c_v$.

(SUNY, Buffalo)

Solution:

(a) From $H = U + pV$, we obtain

$$\left(\frac{\partial H}{\partial T}\right)_p = \left(\frac{\partial U}{\partial T}\right)_p + p \left(\frac{\partial V}{\partial T}\right)_p .$$

Let $U = U[T, V(T, p)]$. The above expression becomes

$$\left(\frac{\partial H}{\partial T}\right)_p = \left(\frac{\partial U}{\partial T}\right)_V + \left[p + \left(\frac{\partial U}{\partial V}\right)_T\right] \left(\frac{\partial V}{\partial T}\right)_p .$$

Hence

$$c_p - c_v = \left[p + \left(\frac{\partial U}{\partial V}\right)_T\right] \left(\frac{\partial V}{\partial T}\right)_p .$$

(b) For the Van der Waals gas, we have

$$\begin{aligned} \left(\frac{\partial p}{\partial T}\right)_V &= \frac{R}{(V-b)} , \\ \left(\frac{\partial V}{\partial T}\right)_p &= R / \left[\frac{RT}{V-b} - \frac{2a(V-b)}{V^3} \right] . \end{aligned}$$

Hence,

$$c_p - c_v = \frac{R}{1 - 2a(1-b/V)^2/VRT} ,$$

When $V \rightarrow \infty$, $c_p - c_v \rightarrow R$, which is just the result for an ideal gas.

1010

One mole of gas obeys Van der Waals equation of state. If its molar internal energy is given by $u = cT - a/V$ (in which V is the molar volume, a is one of the constants in the equation of state, and c is a constant), calculate the molar heat capacities C_v and C_p .

(Wisconsin)

Solution:

$$\begin{aligned} C_v &= \left(\frac{\partial u}{\partial T}\right)_V = c , \\ C_p &= \left(\frac{\partial u}{\partial T}\right)_p + p \left(\frac{\partial V}{\partial T}\right)_p = \left(\frac{\partial u}{\partial T}\right)_V + \left[\left(\frac{\partial u}{\partial V}\right)_T + p\right] \\ &\quad \times \left(\frac{\partial V}{\partial T}\right)_p = c + \left(\frac{a}{V^2} + p\right) \left(\frac{\partial V}{\partial T}\right)_p . \end{aligned}$$

From the Van der Waals equation

$$(p + a/V^2)(V - b) = RT,$$

we obtain

$$\left(\frac{\partial V}{\partial T}\right)_p = R \left/ \left(p - \frac{a}{V^2} + \frac{2ab}{V^3}\right)\right.$$

Therefore

$$C_p = c + \frac{R \left(p + \frac{a}{V^2}\right)}{p - \frac{a}{V^2} + \frac{2ab}{V^3}} = c + \frac{R}{1 - \frac{2a(V-b)^2}{RTV^3}}.$$

1011

A solid object has a density ρ , mass M , and coefficient of linear expansion α . Show that at pressure p the heat capacities C_p and C_v are related by

$$C_p - C_v = 3\alpha M p / \rho.$$

(Wisconsin)

Solution:

From the first law of thermodynamics $dQ = dU + p dV$ and $\left(\frac{dU}{dT}\right)_p \approx \left(\frac{dU}{dT}\right)_v$ (for solid), we obtain

$$C_p - C_v = \left(\frac{dQ}{dT}\right)_p - \left(\frac{dU}{dT}\right)_v = p \frac{dV}{dT}. \quad (*)$$

From the definition of coefficient of linear expansion $\alpha = \alpha_{\text{solid}}/3 = \frac{1}{3V} \frac{dV}{dT}$, we obtain

$$\frac{dV}{dT} = 3\alpha V = 3\alpha \frac{M}{\rho}.$$

Substituting this in (*), we find

$$C_p - C_v = 3\alpha \frac{M}{\rho} p.$$