

科技资料

Visual Communications and Image Processing '90

Part 2

PROCEEDINGS

Visual Communications and Image Processing '90

Fifth in a Series

Murat Kunt
Chair/Editor

1-4 October 1990
Lausanne, Switzerland

Sponsored by
SPIE—The International Society for Optical Engineering

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Swiss Federal Institute of Technology

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IEEE Circuits and Systems Society
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Optical Society of America



Published by
SPIE—The International Society for Optical Engineering
P.O. Box 10, Bellingham, Washington 98227-0010 USA



Volume 1360
Part Two of Three Parts

SPIE (The Society of Photo-Optical Instrumentation Engineers) is a nonprofit society dedicated to advancing engineering and scientific applications of optical, electro-optical, and optoelectronic instrumentation, systems, and technology.



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Please use the following format to cite material from this book:

Author(s), "Title of Paper," *Visual Communications and Image Processing '90*, Murat Kunt, Editor, SPIE Vol. 1360, page numbers (1990).

Library of Congress Catalog Card No. 90-53318
ISBN 0-8194-0421-7

SPIE—The International Society for Optical Engineering
P.O. Box 10, Bellingham, Washington 98227-0010 USA
Telephone 206/676-3290 (Pacific Time) • Fax 206/647-1445

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Printed in the United States of America.

CONTENTS

Conference Committee	xiv
Introduction	xvii

Part One

SESSION 1A	HUMAN VISUAL SYSTEM AND NEURAL-NETWORK-BASED PROCESSING.
1360-01	Human visual quality criterion S. Comes, B. Macq, Univ. Catholique de Louvain (Belgium)..... 2
1360-02	Multiscale image coding using the Kohonen neural network M. Antonini, M. Barlaud, P. Mathieu, Univ. de Nice-Sophia Antipolis (France); J. C. Feauveau, Univ. Paris-Sud (France). 14
1360-03	Target cuing: a heterogeneous neural network approach H. M. McCauley, Naval Weapons Ctr. (USA). 27
1360-04	Bilevel quantization using dithering and Hopfield theory J. J. Hwang, M. H. Lee, Chonbuk National Univ. (Korea). 36
1360-05	New bidirectional neural network and application to binary image recognition S. Zhang, A. G. Constantinides, Imperial College of Science and Technology (UK); L. Zou, Xian Jiaotong Univ. (China). 41
1360-06	NNE-CA: the implementation of neural network emulator board K. Park, K. Cha, J. Choi, Chung-Ang Univ. (Korea)..... 49
1360-07	Multistaged neural network architecture for position invariant shape recognition J. I. Minnix, Stanford Telecommunications Inc. (USA); E. S. McVey, R. M. Iñigo, Univ. of Virginia (USA). 58
1360-08	Image compression using a neural network with learning capability of variable function of a neural unit R. Kohno, M. Arai, H. Imai, Yokohama National Univ. (Japan). 69
SESSION 1B	MASSIVELY PARALLEL COMPUTER ARCHITECTURES
1360-09	ASP: a parallel computing technology (Invited Paper) R. M. Lea, Brunel Univ. (UK). 78
1360-11	Benchmarking the ASP for computer vision A. Krikelis, Brunel Univ. and Aspex Microsystems Ltd. (UK). 92
1360-12	SCC-100 parallel processor for real-time imaging W. J. Jacobi, W. B. Kendall, L. A. Wadsworth, Space Computer Corp. (USA). 104
1360-14	MEGA Node: an implementation of a coarse-grain totally reconfigurable parallel machine M. B. Blum, C. Burrer, Telmat Informatique (France). 109
SESSION 1C	NONLINEAR IMAGE PROCESSING
1360-16	Median-based algorithms for image sequence processing (Invited Paper) B. Alp, P. Haavisto, T. Järskä, K. Oistämö, Y. Neuvo, Tampere Univ. of Technology (Finland). 122
1360-17	Nonlinear quincunx interpolation filtering A. Lehtonen, M. Renfors, Nokia Research Ctr. (Finland). 135

(continued)

VISUAL COMMUNICATIONS AND IMAGE PROCESSING '90

Volume 1360

1360-18	Minimax optimization over the class of stack filters M. Gabbouj, E. J. Coyle, Purdue Univ. (USA).	143
1360-19	Morphological filtering of noisy images L. Koskinen, J. Astola, Tampere Univ. (Finland); Y. Neuvo, Tampere Univ. of Technology (Finland).	155
1360-20	Morphological filtering and iteration H. J. Heijmans, Ctr. for Mathematics and Computer Science (Netherlands).	166
1360-21	Quantitative comparison of median-based filters T. G. Campbell, J. M. H. du Buf, Swiss Federal Institute of Technology (Switzerland).	176
1360-22	Nonlinear spatial filtering of FLIR images M. J. Pérez-Luque, C. Muñoz, N. García, Univ. Politécnica de Madrid (Spain).	188
SESSION 1D MATHEMATICAL MORPHOLOGY AND FRACTALS		
1360-23	Links: definition and properties J. Serra, Ecole des Mines de Paris (France).	202
1360-24	Minimal search for the optimal mean-square digital gray-scale morphological filter E. R. Dougherty, Rochester Institute of Technology (USA).	214
1360-25	Fractal image coding based on a theory of iterated contractive image transformations A. E. Jacquin, Georgia Institute of Technology (USA).	227
1360-26	Determining watersheds in digital pictures via flooding simulations P. Soille, L. Vincent, Ecole des Mines de Paris (France).	240
1360-27	Digital euclidean skeletons F. Meyer, Ecole des Mines de Paris (France).	251
1360-28	Mathematical morphology on the sphere J. B. Roerdink, Ctr. for Mathematics and Computer Science (Netherlands).	263
1360-29	Antiskelton: some theoretical properties and application M. Schmitt, Thomson-CSF (France).	272
1360-30	Image mosaic and interpolation by multiresolution morphological pyramids S. Pei, H. Tsai, National Taiwan Univ. (Taiwan).	284
1360-31	Subband image decomposition by mathematical morphology S. Pei, National Taiwan Univ. (Taiwan); F. Chen, National Taiwan Institute of Technology (Taiwan).	293
SESSION 1E VLSI IMPLEMENTATION AND SYSTEM ARCHITECTURES I		
1360-32	Reconfigurable architecture for real-time 3-D parameter estimation from image sequences F. M. Hugen, M. J. Korsten, Z. Houkes, Univ. of Twente (Netherlands).	304
1360-33	Real-time VLSI architecture for geometric image transformations M. Zhao, J. Gobert, Labs. d'Electronique Philips (France); O. Schirvanian, N. Demassieux, Télécom Paris Univ. (France).	316
1360-34	CCD focal-plane real-time image processor E. Eid, E. R. Fossum, Columbia Univ. (USA).	327
1360-35	Mapping technique for VLSI/WSI implementation of multidimensional systolic arrays M. B. Abdelrazik, Brunel Univ. (UK).	332
1360-36	Mixed digital/analog VLSI array architectures for image processing M. Soma, T. Alexander, Univ. of Washington (USA).	341

VISUAL COMMUNICATIONS AND IMAGE PROCESSING '90

Volume 1360

1360-37	Modular Image Processor: an efficient chip set for real-time image processing H. Waldburger, J. Dufour, G. Concordel, Thomson-TRT Défense (France).	349
1360-38	Analog parallel processor hardware for high-speed pattern recognition T. Daud, R. Tawel, H. Langenbacher, S. P. Eberhardt, A. P. Thakoor, Jet Propulsion Lab. (USA).	359
1360-39	Foveating vision systems architecture: image acquisition and display Y. Y. Zeevi, R. Ginosar, Technion—Israel Institute of Technology (Israel).	371
SESSION 1F	VLSI IMPLEMENTATION AND SYSTEM ARCHITECTURES II	
1360-40	Parallel architecture for real-time video communication L. A. de Sá, V. M. Silva, F. Perdigão, S. Faria, P. Assunção, Univ. of Coimbra (Portugal).	380
1360-41	VLSI components for a 560-Mbit/s HDTV codec K. Grüger, P. Pirsch, Univ. Hannover (FRG); J. Kraus, J. Reimers, Deutschen Bundespost TELEKOM (FRG).	388
1360-42	VLSI architectures for the hierarchical block-matching algorithm for HDTV applications L. De Vos, Siemens AG (FRG).	398
1360-43	VLSI architecture and implementation of a multifunction, forward/inverse discrete cosine transform processor M. Maruyama, H. Uwabu, I. Iwasaki, H. Fujiwara, T. Sakaguchi, Graphic Communication Technologies, Ltd. (USA); M. Sun, M. L. Liou, Bell Communications Research (USA).	410
SESSION 1G	3-D IMAGE PROCESSING	
1360-44	Determining vanishing points using Hough transform E. Lutton, H. Maître, J. Lopez-Krahe, Télécom Paris (France).	420
1360-45	Monocular correspondence detection for symmetrical objects by template matching G. Vilmar, P. W. Besslich, Univ. of Bremen (FRG).	431
1360-46	3-D reconstruction using a limited number of projections C. Klifa, Télécom Paris (France); B. Lavayssière, EDF-DER (France).	443
1360-47	Mathematical morphology for 3-D object segmentation and partial matching I. Bloch-Boulanger, Télécom Paris and Rhône-Poulenc Santé (France); H. Maître, F. J. Schmitt, Télécom Paris (France).	455
SESSION 1H	IMAGE SEQUENCE CODING I	
1360-48	Motion estimation for coding of moving video at 8 kbit/s with Gibbs-modeled vectorfield smoothing C. Stiller, Aachen Univ. of Technology (FRG).	468
1360-49	Some variants of universal pattern-matching interframe coding T. Saito, R. Abe, T. Komatsu, Kanagawa Univ. (Japan); H. Harashima, Univ. of Tokyo (Japan).	477
1360-50	Video coding using a pyramidal Gabor expansion T. Ebrahimi, T. R. Reed, Ecole Polytechnique Federale de Lausanne (Switzerland); M. Kunt, Swiss Federal Institute of Technology (Switzerland).	489
1360-51	Block testing in a variable resolution spatially interpolative moving image sequence coder P. J. Cordell, R. J. Clarke, Heriot-Watt Univ. (UK).	503
1360-52	Effective exploitation of background memory for coding of moving video using object mask generation W. Guse, RWTH Aachen (FRG); M. Gilge, International Computer Science Institute (FRG); B. Hürtgen, RWTH Aachen (FRG).	512

(continued)

VISUAL COMMUNICATIONS AND IMAGE PROCESSING '90

Volume 1360

1360-54	Very low rate coding of motion video using 3-D segmentation with two change-detection masks S. Lee, N. Kim, H. Son, Kyungpook National Univ. (Korea).....	524
1360-55	Visual pattern image sequence coding P. Silsbee, A. C. Bovik, D. Chen, Univ. of Texas/Austin (USA).....	532
SESSION 11	HIERARCHICAL VIDEO CODING	
1360-56	Subband coding of video employing efficient recursive filter banks and advanced motion compensation J. H. Husoy, H. Gronning, T. A. Ramstad, Norwegian Institute of Technology (Norway).....	546
1360-57	Interframe hierarchical address-vector quantization N. M. Nasrabadi, Worcester Polytechnic Institute (USA).....	558
1360-58	Refinement system for hierarchical video coding F. Bosveld, R. L. Lagendijk, J. Biemond, Delft Univ. of Technology (Netherlands).....	575
1360-59	Design of an HDTV subband codec considering CMOS-VLSI constraints U. Pestel, B. Schmale, Univ. Hannover (FRG).....	587
SESSION 1J	HIERARCHICAL IMAGE CODING	
1360-60	Image subband coding using an efficient recursive filter bank with complex signals H. Gronning, J. H. Husoy, T. A. Ramstad, Norwegian Institute of Technology (Norway).....	598
1360-61	Perfect reconstruction binomial QMF-wavelet transform A. N. Akansu, New Jersey Institute of Technology (USA); R. A. Haddad, Polytechnic Univ. (USA); H. Caglar, New Jersey Institute of Technology (USA).....	609
1360-62	Generalized quad-trees: a unified approach to multiresolution image analysis and coding R. Wilson, M. Todd, Univ. of Warwick (UK); A. D. Calway, Linköping Univ. (Sweden).....	619
1360-63	Three-dimensional adaptive Laplacian pyramid image coding S. Sallent, L. Torres, L. Gils, ETSI Telecomunicación (Spain).....	627
1360-184	Image representation using binary space partitioning trees H. Radha, AT&T Bell Labs. and Columbia Univ. (USA); R. Leonardi, B. Naylor, AT&T Bell Labs. (USA); M. Vetterli, Columbia Univ. (USA).....	639
 Part Two		
SESSION 2A	DIGITAL IMAGE PROCESSING IN MEDICINE I	
1360-66	Three-dimensional reconstruction and lateral views in optical microscopy T. Tommasi, B. Bianco, V. Murino, A. Oneto, A. Diaspro, Univ. of Genoa (Italy).....	652
1360-67	Improved resolution of medical 3-D x-ray computed-tomographic images C. Odet, G. Jacquemod, F. Peyrin, R. Goutte, INSA (France).....	658
1360-68	Extraction of morphometric information from dual echo magnetic resonance brain images T. Sandor, F. A. Jolesz, J. Tieman, R. Kikinis, M. LeMay, M. Albert, Harvard Medical School (USA).....	665
1360-70	Diagnostic digital image processing of human corneal endothelial cell patterns B. R. Masters, Georgia Institute of Technology (USA).....	676
1360-71	Region-oriented 3-D segmentation of NMR datasets: a statistical model-based approach T. Aach, H. Dawid, RWTH Aachen (FRG).....	690
1360-72	Enhancement and segmentation for NMR images of blood flow in arteries G. Yang, P. Burger, Imperial College of Science, Technology and Medicine (UK).....	702

VISUAL COMMUNICATIONS AND IMAGE PROCESSING '90

Volume 1360

SESSION 2B	DIGITAL IMAGE PROCESSING IN MEDICINE II	
1360-74	Analysis of x-ray hand images for bone age assessment J. Serrat, J. Vitrià, J. J. Villanueva, Univ. Autònoma de Barcelona (Spain).	716
1360-75	ITS: a practical picture archiving and communication system T. Lei, W. Sewchand, Univ. of Maryland School of Medicine (USA).	724
1360-76	Open system architecture for distributed image-reference database in radiological applications A. Bellini, G. Bucci, Univ. di Firenze (Italy).	726
SESSION 2C	HDTV	
1360-78	Motion field restoration using vector median filtering on high-definition television sequences T. Koivunen, A. Nieminen, Nokia Consumer Electronics (Finland).	736
1360-79	Motion-adaptive four-channel HDTV subband/DCT coding G. Schamel, Heinrich-Hertz-Institut für Nachrichtentechnik (FRG).	743
1360-80	Compression and channel-coding algorithms for high-definition television signals L. Alparone, G. Benelli, A. F. Fabbri, Univ. di Firenze (Italy).	754
1360-81	Source coding of HDTV with compatibility to TV M. Breeuwer, P. H. de With, Philips Research Labs. (Netherlands).	765
SESSION 2D	PARALLEL PROCESSING	
1360-82	Xputer use in image processing and digital signal processing R. W. Hartenstein, A. G. Hirschbiel, K. Lemmert, M. Riedmüller, K. Schmidt, M. Weber, Univ. Kaiserslautern (FRG).	778
1360-83	Parallel architectures for the postprocessing of SAR images L. Alparone, F. Boragine, S. Fini, F. de Stefani, Univ. di Firenze (Italy).	790
1360-85	Transputer-based embedded system for METEOSAT image data compression M. H. Versteeg, R. A. Hogendoorn, A. Monkel, National Aerospace Lab. NLR (Netherlands).	803
SESSION 2E	IMAGE CODING AND TRANSMISSION I	
1360-86	Source coding of super high definition images with discrete cosine transform M. Nomura, T. Fujii, N. Ohta, NTT Transmission Systems Labs. (Japan).	814
1360-87	New variable-rate VQ coding scheme apply in HDTV Y. Feng, Prime Computervision (USA); K. Zhang, Worcester Polytechnic Institute (USA).	826
1360-88	Clustering algorithm for entropy-constrained vector-quantizer design W. A. Finamore, D. P. de Garrido, IBM/Rio Scientific Ctr. (Brazil); W. A. Pearlman, Rensselaer Polytechnic Institute (USA).	837
1360-89	Variable block-sized vector quantization of gray-scale images with unconstrained tiling J. L. Boxerman, Massachusetts Institute of Technology (USA); H. J. Lee, Tetra Systems Inc. (USA).	847
SESSION 2F	IMAGE CODING AND TRANSMISSION II	
1360-90	New technique of linear-phase QMF filter design for subband coding J. Jeon, J. Kim, Korea Advanced Institute of Science and Technology (Korea).	860
1360-91	Gain-adaptive trained transform trellis code for images D. Kim, W. A. Pearlman, Rensselaer Polytechnic Institute (USA).	868

(continued)

VISUAL COMMUNICATIONS AND IMAGE PROCESSING '90

Volume 1360

1360-92	DS-DCT: the double-shift DCT image coding for low-bit-rate image transmission D. Cai, Institute of Electronics (China); Y. Chen, Univ. of Tokyo (Japan).	879
1360-93	Efficient error-resilient codes for sparse image coefficients N. Cheng, N. G. Kingsbury, Univ. of Cambridge (UK).	888
1360-94	Optimum quantization for subband coders L. Vandendorpe, B. Macq, Univ. Catholique de Louvain (Belgium).	898
1360-95	Statistical dependence between orientation filter outputs used in a human-vision-based image code B. Wegmann, C. Zetsche, Technische Univ. München (FRG).	909
1360-96	Coding gains of pyramid structures in progressive image transmission S. H. Park, S. U. Lee, Seoul National Univ. (Korea).	924
1360-97	Image data compression using hybrid POLA-VQ technique C. Lee, Ministry of Communications (Taiwan); R. Ju, Ministry of Communications and National Central Univ. (Taiwan); T. Liu, Ministry of Communications (Taiwan); B. Jeng, Ministry of Communications and National Central Univ. (Taiwan); J. Huang, K. Kan, Ministry of Communications (Taiwan).	936
SESSION 2G	EDGE/BOUNDARY DETECTION	
1360-99	Object contours and boundaries in color images R. Ronfard, Ecole des Mines de Paris and Télécom Paris (France).	946
1360-100	Attributed tree data structure for representing the descriptions of object contours in images Z. Ren, W. Ameling, Technische Hochschule Aachen (FRG); P. J. Jensch, Univ. Oldenburg (FRG).	956
1360-101	Two design techniques for 2-D FIR LoG filters P. Siohan, D. Pelé, V. Ouvrard, CCETT (France).	970
1360-102	Adaptable edge quality metric R. N. Strickland, D. K. Chang, Univ. of Arizona (USA).	982
SESSION 2H	NEUROMORPHOLOGY OF BIOLOGICAL VISION I	
1360-103	Neuromorphology of biological vision: a basis for machine vision M. M. Gupta, Univ. of Saskatchewan (Canada).	998
1360-104	Color-subspace-based color-coordinate system J. P. Parkkinen, Univ. of Kuopio (Finland) and Univ. of Iowa (USA); J. Hallikainen, T. Jaaskelainen, Univ. of Kuopio (Finland).	1010
1360-105	GRUPO: a 3-D structure recognition system S. T. Acton, A. C. Bovik, Univ. of Texas/Austin (USA).	1018
SESSION 2I	NEUROMORPHOLOGY OF BIOLOGICAL VISION II	
1360-106	Global stability in nonlinear lateral inhibition G. F. McLean, Univ. of Victoria (Canada); M. E. Jernigan, Univ. of Waterloo (Canada).	1032
1360-107	Dynamic neural network for visual memory M. M. Gupta, G. K. Knopf, Univ. of Saskatchewan (Canada).	1044
1360-108	Binocular fusion inferences in a log-polar decision space N. C. Griswold, N. Kehtarnavaz, Texas A&M Univ. (USA).	1056
1360-109	Dense color stereo J. R. Jordan III, A. C. Bovik, Univ. of Texas/Austin (USA).	1069

1360-111	Image invariance with changes in distance: the effect of a nonuniform visual system E. Peli, Eye Research Institute and Harvard Medical School (USA); J. Yang, Northeastern Univ. (USA); R. B. Goldstein, Eye Research Institute and Harvard Medical School (USA).	1079
SESSION 2J	IMAGE SEQUENCE CODING II	
1360-112	Coding of moving video at 1 mbit/s: movies on CD B. Hürtgen, RWTH Aachen (FRG); M. Gilge, International Computer Science Institute (FRG); W. Guse, RWTH Aachen (FRG).	1092
1360-113	Vector quantization with 3-D gradient motion compensation C. Lee, M. Nadler, Virginia Polytechnic Institute and State Univ. (USA).	1104
1360-114	Two-layers constant-quality video coding for ATM environments F. M. Pereira, Instituto Superior Técnico (Italy); L. Masera, Centro Studi e Lab. Telecomunicazioni (Italy).	1114
1360-115	Image sequence representation using polar-separable filters T. G. Campbell, T. R. Reed, M. Kunt, Swiss Federal Institute of Technology (Switzerland).	1126
1360-117	Visual model weighted DCT vector quantization for variable bit-rate video coding F. Lavagetto, S. Zappatore, Univ. di Genova (Italy).	1134
1360-118	Encoding of sign language image sequences at very low rate C. Huang, C. H. Wu, National Tsing Hua Univ. (Taiwan).	1140
1360-119	Real-time facial action image synthesis system driven by speech and text S. Morishima, Seikei Univ. (Japan); K. Aizawa, H. Harashima, Univ. of Tokyo (Japan).	1151
1360-120	Image modeling for digital TV codecs J. Leduc, Univ. Catholique de Louvain (Belgium).	1160
1360-121	Video signal processing using vector median K. Oistämö, Y. Neuvo, Tampere Univ. of Technology (Finland).	1171
1360-123	Image analysis for face modeling and facial image reconstruction H. Agawa, G. Xu, Y. Nagashima, F. Kishino, ATR Communication Systems Research Labs. (Japan).	1184
1360-124	General motion estimation and segmentation S. Wu, J. Kittler, Univ. of Surrey (UK).	1198
1360-126	Adaptive algorithms for pel-recursive displacement estimation L. Böröczky, Technical Univ. of Budapest (Hungary); J. N. Driessen, J. Biemond, Delft Univ. of Technology (Netherlands).	1210
1360-127	Distributed detection methods for displacement estimation S. N. Efstratiadis, A. K. Katsaggelos, Northwestern Univ. (USA).	1222
1360-180	Control analysis of video packet loss in ATM networks D. Lee, Columbia Univ. (USA); K. Tzou, Bell Communications Research (USA); S. Li, Univ. of Texas/Austin (USA).	1232

Part Three

SESSION 3A	SEGMENTATION/CLASSIFICATION	
1360-128	Surface defect detection using adaptive image modeling P. Salembier, Swiss Federal Institute of Technology (Switzerland).	1246
1360-129	Thresholding three-dimensional image Y. J. Zhang, J. J. Gerbrands, E. Backer, Delft Univ. of Technology (Netherlands).	1258

(continued)

VISUAL COMMUNICATIONS AND IMAGE PROCESSING '90

Volume 1360

1360-130	Edge point/region cooperative segmentation specific to 3-D scene reconstruction application P. Bonnin, Etablissement Technique Central de l'Armement (France); B. Zavidovique, ETCA and Univ. de Paris XI (France).	1270
1360-131	Sequential classification into m multivariate populations using the information based on small samples N. A. Nechval, Civil Aviation Engineers Institute (USSR).	1282
SESSION 3B	TEXTURE	
1360-132	Adaptive classification of textured images using moments and autoregressive models L. Sukissian, A. Tirakis, S. Kollias, National Technical Univ. of Athens (Greece).	1296
1360-133	Classification of textures in aerial images G. Brunet, J. Devars, ENSEA (France).	1307
1360-134	Frequency and orientation selective texture measures using linear symmetry and Laplacian pyramid J. Bigün, Swiss Federal Institute of Technology (Switzerland).	1319
1360-135	Texture classification using transform vector quantization G. F. McLean, Univ. of Victoria (Canada).	1332
SESSION 3C	IMAGE RESTORATION	
1360-136	Tutorial review of recent developments in digital image restoration (Invited Paper) M. I. Sezan, Eastman Kodak Co. (USA); A. M. Tekalp, Univ. of Rochester (USA).	1346
1360-137	Maximum-likelihood blur identification R. L. Lagendijk, J. Biemond, Delft Univ. of Technology (Netherlands).	1360
1360-138	Optimal constraint parameter estimation for constrained image restoration S. J. Reeves, Auburn Univ. (USA); R. M. Mersereau, Georgia Institute of Technology (USA).	1372
1360-139	Multiple input adaptive image restoration algorithms A. K. Katsaggelos, Northwestern Univ. (USA).	1381
1360-140	Robust estimation of local orientations in images using a multiresolution approach R. Wilson, Univ. of Warwick (UK); S. C. Clippingdale, NHK Science and Technical Research Labs. (Japan); A. H. Bhalerao, Univ. of Warwick (UK).	1393
1360-141	Image restoration using biorthogonal wavelet transform J. M. Bruneau, M. Barlaud, P. Mathieu, Univ. de Nice-Sophia Antipolis (France).	1404
1360-142	Stochastic model-based approach for simultaneous restoration of multiple misregistered images C. Srinivas, General Electric Corporate Research and Development (USA); M. D. Srinath, Southern Methodist Univ. (USA).	1416
1360-143	Optical methods for iterative image restoration A. K. Katsaggelos, T. E. DeRoux, M. E. Marhic, Northwestern Univ. (USA).	1428
SESSION 3D	DIGITAL IMAGE PROCESSING	
1360-144	Stability analysis of multichannel linear-predictive systems Y. Öztürk, Ege Univ. Izmir (Turkey) and San Diego State Univ. (USA); H. Abut, San Diego State Univ. (USA).	1442
1360-145	Group delay equalization of multidimensional recursive filters F. T. Tehrani, R. E. Ford, California State Univ./Fullerton (USA).	1454

VISUAL COMMUNICATIONS AND IMAGE PROCESSING '90

Volume 1360

1360-146	Perceptually relevant model for aliasing in the triplet-stripe filter CCD image sensor R. A. Beuker, Philips Research Labs. (Netherlands); F. W. Hoeksema, Univ. of Twente (Netherlands).	1463
1360-147	Gram-Gabor approach to optimal image representation M. Porat, Y. Y. Zeevi, Technion—Israel Institute of Technology (USA).	1474
1360-148	Fast method of geometric picture transformation using logarithmic number systems and its applications for computer graphics T. Kurokawa, Aichi Institute of Technology (Japan); T. Mizukoshi, Oki Technosystems Lab. (Japan).	1479
1360-149	Space-variant filtering of images through a hybrid implementation of the Wigner distribution function C. Gonzalo, Instituto de Optica (Spain).	1491
1360-150	Curved shadow generation by ray tracing in the plane A. Tokuta, Univ. of South Florida (USA).	1499
1360-151	Computation network for visible surface reconstruction H. Hsieh, W. Chang, National Chiao Tung Univ. (Taiwan).	1504
SESSION 3E	JPEG/MPEG ALGORITHMS AND IMPLEMENTATION	
1360-152	Color image transmission system for prepress with ADCT compression algorithm H. Hasegawa, M. Sugiura, H. Ono, N. Kurakami, T. Omachi, NEC Corp. (Japan).	1518
1360-153	Design of a multifunction video decoder based on a motion-compensated predictive-interpolative coder K. Yang, S. Singhal, Bell Communications Research (USA); D. J. Le Gall, C-Cube Microsystems (USA).	1530
1360-155	Video coding for digital storage media using hierarchical intraframe scheme K. Kamikura, H. Watanabe, NTT Human Interface Labs. (Japan).	1540
1360-156	Principal devices and hardware volume estimation for moving picture decoder for digital storage media M. Konoshima, O. Kawai, K. Matsuda, Fujitsu Labs. Ltd. (Japan).	1551
1360-183	Comparing motion-interpolation structures for video coding A. Puri, R. Aravind, AT&T Bell Labs. (USA).	1560
1360-185	Encoding of motion video sequences for the MPEG environment using arithmetic coding E. Viscito, C. A. Gonzales, IBM/Thomas J. Watson Research Ctr. (USA).	1572
SESSION 3F	VISION SCIENCE AND TECHNOLOGY FOR SPACE	
1360-157	Digital image gathering and minimum mean-square error restoration S. K. Park, College of William and Mary (USA); S. E. Reichenbach, Univ. of Nebraska (USA).	1578
1360-158	Information theoretical assessment of image gathering and coding for digital restoration (Invited Paper) F. O. Huck, NASA/Langley Research Ctr. (USA); S. John, Science and Technology Corp. (USA); S. E. Reichenbach, Univ. of Nebraska (USA).	1590
1360-159	Image coding by edge primitives R. Alter-Gartenberg, Old Dominion Univ. (USA); R. Narayanswamy, Science and Technology Corp. (USA).	1608
1360-160	Photon detection with parallel-asynchronous processing D. D. Coon, A. G. Perera, Microtronics Associates, Inc. (USA).	1620

(continued)

VISUAL COMMUNICATIONS AND IMAGE PROCESSING '90

Volume 1360

1360-161	Vision-based planetary rover navigation B. H. Wilcox, Jet Propulsion Lab. (USA)	1628
1360-162	Expert imaging system G. B. Westrom, Odetics Inc. (USA)	1634
1360-163	Hybrid image processing R. D. Juday, NASA/Johnson Space Ctr. (USA)	1647
SESSION 3G	PATTERN RECOGNITION	
1360-164	Retrieval of script information appearing on bank checks for automatic reading purposes J. C. Salomé, M. Leroux, H. Oiry, French Post Technical Research Dept. (France); A. Saad, Univ. of Rennes (France)	1652
1360-165	Use of discrete-state Markov process for Chinese-character recognition B. Jeng, National Central Univ. and Ministry of Communications (Taiwan); C. Shih, S. Sun, T. Wu, B. Chien, Ministry of Communications (Taiwan); M. Chang, National Central Univ. (Taiwan)	1663
1360-166	Underground radar system utilizing pattern-recognition technique in the frequency domain Y. Nagashima, J. Masuda, R. Arioka, S. Kouno, NTT Corp. (Japan); Y. Sudoh, Airec Engineering Corp. (Japan)	1671
1360-167	Rotation- and translation-invariant pattern recognition based on distance transformations J. Möschen, B. J. Hosticka, Fraunhofer Institute of Microelectronic Circuits and Systems (FRG)	1682
1360-168	Description and matching of density variation for personal identification through facial images K. Takahashi, T. Sakaguchi, T. Minami, O. Nakamura, Kogakuin Univ. (Japan)	1694
1360-169	Footprint image processing expert system with friendly user interface Y. Hattori, T. Minami, O. Nakamura, Kogakuin Univ. (Japan)	1705
1360-170	Extraction of arbitrary shapes from a noisy binary image using pseudo view field tracer T. Nagao, T. Agui, M. Nakajima, Tokyo Institute of Technology (Japan)	1719
1360-171	2-D invariant color pattern recognition using complex log-mapping transform S. Pei, M. Hsu, National Taiwan Univ. (Taiwan)	1729
1360-181	Zak-Gabor representation of images I. Gertner, City Univ. of New York (USA); Y. Y. Zeevi, Technion—Israel Institute of Technology and Rutgers Univ. (USA)	1738
SESSION 3H	IMAGE SEQUENCE CODING III	
1360-173	Fast codebook search algorithm in vector quantization J. Huguet, L. Torres, Univ. Politecnica de Cataluna (Spain)	1750
1360-174	Digital HDTV compression at 44 Mbps using parallel motion-compensated transform coders H. Hang, R. Leonardi, B. G. Haskell, R. L. Schmidt, H. Bheda, J. H. Othmer, AT&T Bell Labs. (USA)	1756
1360-175	Hybrid DCT encoding of TV and HDTV: a comparative study N. García, J. I. Ronda, F. Jaureguizar, Univ. Politécnica de Madrid (Spain)	1773
1360-176	Development of video teleconference system Y. Kobayashi, Tokyo Electric Power Co. (Japan); K. Toudo, NEC Corp. (Japan)	1783

VISUAL COMMUNICATIONS AND IMAGE PROCESSING '90

Volume 1360

1360-177	Motion-compensated adaptive interframe/intraframe prediction K. Xie, L. Van Eycken, A. J. Oosterlinck, Univ. of Leuven (Belgium).	1798
1360-179	Analysis of a 2-D DCT image coding scheme with motion compensation and vector quantization R. Picco, F. L. Bellifemine, A. Chimienti, National Research Council (Italy).	1810
1360-116	Pyramidal encoding for packet video transmission (Additional Paper) L. Salgado, A. Sanz, Univ. Politécnica de Madrid (Spain).	1822
1360-172	Asynchronous conditional-replenishment video codec for 64-kbits/s channels (Additional Paper) J. M. Menéndez, C. Muñoz, Univ. Politécnica de Madrid (Spain).	1834
1360-73	Enhancement and delineation of lung tumors in local x-ray chest images (Additional Paper) F. Zhou, Univ. Hospital Zurich (Switzerland); H. B. Zhou, Univ. Zurich (Switzerland); X. Q. Wu, Northern Jiaotong Univ. (China); L. J. He, Beijing Institute of Tumor Studies (China)	1846
	Author Index	1854

Three-dimensional reconstruction and lateral views in optical microscopy

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ABSTRACT

The three-dimensional properties of an optical microscope are analyzed and a defocusing technique is proposed to recover the spatial distribution of the specimens under investigation. Limitations and real capabilities of the 3D reconstruction are pointed out. As a result, lateral views are obtained by means of operations in the spatial frequency domain.

In such a way, it is possible to represent side views of an object within the angular aperture range of the microscope. A theory concerning image formation is discussed and simulations of side view reconstructions are reported.

1. INTRODUCTION

Three-dimensional (3D) reconstruction of observed objects is one of major task of today's imaging science. In particular, in the microscope imaging of biological specimens, the goal of obtaining a spatial distribution of an intensive physical property of the objects should be achieved avoiding alterations to the structures.

By using an optical sectioning instead of a physical one it is possible to recover internal planes, not destroying or altering the specimens. The 3D properties of an optical microscope have been discussed in many articles (see for example [1,2]) and restored images have been obtained in particular situations [3]. However, the problem is not so simple as it may appear in the scientific literature, where the various methodologies often are based on a priori information which permits a satisfactory reconstruction.

We present a general technique based on defocusing, with no a priori information. In such a way, limits and possibilities of the reconstruction process are stressed, and we prove that a total 3D reconstruction is not possible, since in the defocusing operation a part of information is irreparably lost. This fact can be easily seen by a spectral analysis and a calculation of the Optical Transfer Function (OTF) of the system leads to define the limits of restoration. On the other hand, we analyze the possibilities of the defocusing method and to this end we used a simplified but effective model of the microscope. The model is obtained under the hypotheses of geometrical optics. This means to disregard diffraction, i.e., to state that the system is rigorously stigmatic. The assumption is not correct, but the error is a minor one: one must bear in mind that if a 3D reconstruction has some limits in a diffraction-free microscope, it can hardly be presumed that a better behaviour will be exhibited by any instrument with diffraction limitations.

In the next section the image formation theory is presented. In section 3, we analyze the restoration task. Simulation and results are discussed in the last section.

2. IMAGE FORMATION

The properties of out-of-focus images have been extensively studied in the last decades (see for example [4]). Assuming the microscope as a linear system, if the plane of acquisition (e.g. a TV camera target) is not properly focused, then the relationship between the actual image intensity $i(x,y)$ and the ideal (in focus) one $i_0(x,y)$ is simply a convolution by some function $s(x,y)$.

$$i(x,y) = i_0(x,y) * s(x,y) \quad (1)$$

Equation (1) holds if the image plane is perpendicular to the optical z-axis.

Then, x and y are coordinates perpendicular to the z -axis; $s(x,y)$ is a function depending on the imaging system used and on the degree of defocusing. Once $s(x,y)$ and $i(x,y)$ are known, an inverse convolution yields the required in-focus image $i_0(x,y)$.

The 3D image formation of microscopic objects follows a similar approach. In this case $s(x,y,z)$ is the 3D Point Spread Function (described later on), $i_0(x,y,z)$ represents the 3D object, i.e., the 3D transmission coefficient while $i(x,y,z)$ is the image obtained; more precisely, it is a set of 2D images at fixed z , formed by the in-focus planes and by the out-of-focus planes under and over the in-focus ones. The result is a set of blurred images and the goal consists in reconstructing the in-focus planes, i.e., the object, starting from $i(x,y,z)$. In a computerized analysis, z is a discrete quantity, and every plane $i(x,y,z)$ is a digitized image.

We consider incoherent illumination, as one normally has in microscopy; introducing the 3D Fourier Transform I, I_0 and S of i, i_0 and s respectively, the well-known convolution property yields

$$I(f_x, f_y, f_z) = I_0(f_x, f_y, f_z) S(f_x, f_y, f_z) \quad (2)$$

If we consider geometric optics, an expression for s is:

$$s(x,y,z) = \begin{cases} \frac{A}{z^2} & \text{if } \sqrt{x^2+y^2} < |z| \tan \theta \\ 0 & \text{otherwise} \end{cases} \quad (3)$$

where A is a constant and θ is the angular aperture of the microscope. The calculation of S leads to a very important result which has a considerable effect on the 3D reconstruction. The Fourier Transform of equation (3) gives the following condition:

$$S(f_x, f_y, f_z) = 0 \quad \text{if } |f_z| > \tan \theta \sqrt{f_x^2 + f_y^2} \quad (4)$$

i.e., S is zero inside the double cone in the Fourier space, whose angle of aperture is formed by the f_z axis and $\pi/2 - \theta$. In the Appendix eq.(4) is proven and the analytic expression of S in the region where it does not vanish is calculated.

The fact that $S(f_x, f_y, f_z)$ is zero in a wide region can be explained by some considerations in the frequency domain [5].

The wave coming from the object under investigation can be decomposed as a superposition of plane waves with different directions [6]. Only the plane waves entering the microscope at angles not higher than the input angular aperture θ of the microscope contribute to the image formation. In fact, high frequencies are lost.

A more accurate expression for $s(x,y,z)$ can be used [7] which does not neglect diffraction, but the conclusion we obtain is always valid, i.e., S has regions where it vanishes.

3. 3-D RECONSTRUCTION

We aim at obtaining internal planes of the specimens and a restoration has to be performed to eliminate the contribution of the out-of-focus planes. Starting from equation (1), an inverse convolution yields i_0 , i.e., the set of in-focus planes. The procedure could be the following one: calculate i from the images, through a Fourier Transform; measure or calculate s and deduce its transform. Then (2) yields I_0 , from which an inverse transform derives the desired i_0 . This simple operation cannot be performed since S is equal to zero in entire regions of the frequency domain. This implies that a perfect 3-D reconstruction is not possible.

On the contrary, a correct information can be obtained about the projections of the object in all directions comprised within an angle relative to the optical axis. By using the projection theorem, if we imagine to superimpose the two frames: f_x, f_y, f_z , and x, y, z , then the projection of the object parallel to a direction z' , is the inverse Fourier transform of the function obtained cutting the Fourier space by a

plane through the origin perpendicular to z' .

Hence if zz' angle is $< \theta$ such projection is obtained. If not, the above said plane is in the region where S vanishes, and no information can be recovered. Analytically, let us consider a new reference frame x', y', z' , obtained by adding two orthogonal axes x' and y' to z' , leaving the origin unchanged, and let f_x, f_y, f_z be the spatial frequencies conjugate to x', y', z' . Then, if A is the orthonormal matrix that links the $x y z$ and $x' y' z'$ frames, one has :

$$dx dy dz = \det(A) dx' dy' dz' = dx' dy' dz'$$

and

$$x f_x + y f_y + z f_z = x' f'_x + y' f'_y + z' f'_z \quad (5)$$

then

$$I_0(f'_x, f'_y, f'_z) = \iiint i_0(x', y', z') \exp(-2\pi j(x' f'_x + y' f'_y + z' f'_z)) dx' dy' dz' \quad (6)$$

If one considers the plane $f'_z = 0$, then (6) can be written in the following way :

$$\begin{aligned} I_0(f'_x, f'_y, 0) &= \iiint i_0(x', y', z') \exp(-2\pi j(x' f'_x + y' f'_y)) = \\ &= \iint \exp(-2\pi j(x' f'_x + y' f'_y)) \int i_0(x', y', z') dz' dx' dy' = \\ &= F \left(\int i_0(x', y', z') dz' \right) \end{aligned} \quad (7)$$

Equation (7), which expresses the projection theorem, gives the Fourier transform of the projection of i_0 along the z' direction, perpendicular to the plane $f'_z = 0$.

Once one has calculated $I_0(f'_x, f'_y, f'_z)$, it is sufficient to calculate the Fourier transform in the plane $f'_z = 0$, by means of interpolations of the values of I_0 . An inverse transform of equation (7) gives the projection along the z' direction. We have pointed out that I_0 cannot be recovered; nevertheless, in the region where S does not vanish, we can apply a Wiener filter on I_0 , so obtaining a Fourier transform to be used in the projection theorem.

4. RESULTS AND DISCUSSION

Simulations and actual experiments have been performed.

A three dimensional image has been simulated, consisting in a pyramid with eight planes. The point spread function is a double cone. A 3-D defocused image has been obtained by means of a convolution between the original image and the cone. Starting from this blurred image, a restoration has been performed and side views have been obtained.

The calculations have been performed on $32 \times 32 \times 16$ matrices, in which some zero points have been added to avoid aliasing effects.

Fig. 1 shows a defocused set of planes. Fig. 2 shows the reconstructed images of the pyramid. From these images the different point of view are evident.

In the case of actual experiments difficulties arise when getting the set of images, since it is difficult to adjust with accuracy the defocusing. The software implemented needs a set of images where the in-focus planes are taken at a constant sz distance, and this task is not easy to get if an automatic controlled adjustment of the focus is not available.

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