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A Two-Step Perturbation Method in Nonlinear Analysis of Beams, Plates and Shells

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二次摄动法 (英文版)

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About the Author

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Preface

This book, written in memory of Professor WZ Chien (1912–2010) with great respect, discusses a two-step perturbation method and its applications in the nonlinear analysis of elastic structures. The capability to predict the nonlinear response of beams, plates and shells when subjected to thermal and mechanical loads is of prime interest to structural analysis. In fact, many structures are subjected to high load levels that may result in nonlinear load–deflection relationships due to large deformations. One of the important problems deserving special attention is the study of their nonlinear response to large deflection, postbuckling and nonlinear vibration.

The major difference between the linear analysis and the nonlinear analysis of structures lies in that the principle of superposition is not valid in the latter. Approximate analytical methods, for example, the Ritz method and the Galerkin method, have been used mainly to study nonlinear bending, postbuckling and nonlinear vibration of beams, plates and shells. It was proved that, for nonsymmetric cross-ply laminated plates and functionally graded material (FGM) plates with four edges simply supported subjected to uniaxial or biaxial compression, or uniform temperature rise, bifurcation buckling did not exist due to the stretching/bending coupling effect. Unfortunately, for nonsymmetric cross-ply laminated plates and FGM plates, the Ritz method or finite element method usually obtain physically incorrect solutions that are inconsistent with the prebuckled state. Further, in the traditional perturbation method, the perturbation parameter is no longer a small perturbation parameter in the large deflection, postbuckling and large amplitude vibration region when the plate/shell deflection is sufficiently large. Hence, the accuracy and effectiveness of traditional perturbation solutions for stronger nonlinear problems are doubted by many researchers.

A two-step perturbation method was first proposed by Shen and Zhang (1988) for the postbuckling analysis of isotropic plates. This approach gives explicit analytical expressions for all the variables in the postbuckling range. This approach provides a good physical insight into the problem considered, and the influence of all the parameters on the solution can be assessed easily. The advantage of this method is that it is unnecessary to guess the forms of solutions which can be obtained step by step, and such solutions satisfy both governing equations and boundary conditions accurately in the asymptotic sense. This approach is then successfully used in solving many nonlinear bending, postbuckling and nonlinear vibration problems of beams, plates and shells made of advanced composite materials. This approach may find more extensive applications in the nonlinear analysis of nanoscale structures.

This book comprises seven chapters involving the latest research materials. The present chapter and section titles are a significant indication of the total content. Each chapter

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contains adequate introductory material so that an engineering graduate who is familiar with a basic understanding of beams, plates and shells will be able to follow it. The advantages and disadvantages of the traditional perturbation method are introduced in Chapter 1. A two-step perturbation method and its application in the nonlinear analysis of beams, plates and shells are presented in detail in each chapter. Some difficult tasks in the nonlinear analysis of elastic structures are included, for example: the nonlinear analysis of Euler–Bernoulli beams based on an exact expression of the curvature is presented in Chapter 2; the nonlinear vibration analysis of functionally graded fiber-reinforced composite laminated plates in hygro-thermal environments is presented in Chapter 3; the geometrically nonlinear bending analysis of shear deformable plates with four free edges resting on elastic foundations is presented in Chapter 4; the contact postbuckling analysis of composite laminated plates resting on tensionless elastic foundations subjected to thermal and mechanical loads is presented in Chapter 5; the nonlinear vibration of functionally graded fiber-reinforced composite laminated cylindrical shells without or with piezoelectric fiber-reinforced composite actuators is presented in Chapter 6; the contact postbuckling analysis of anisotropic cylindrical shells surrounded by an elastic medium subjected to mechanical loads in thermal environments is presented in Chapter 7. Most of the solutions presented in these chapters are the results of investigations made by the author and his collaborators since 1997. The results presented herein may be benchmarks for checking the validity and accuracy of other numerical solutions.

At the time of writing this book, despite a number of existing texts in the theory and analysis of plates and/or shells, there is not a single book which is devoted entirely to solve geometrically nonlinear problems of beams, plates and shells by means of a two-step perturbation method. It is hoped that this book will fill the gap to some extent and that it might be used as a valuable reference source for postgraduate students, engineers, scientists and applied mathematicians in this field.

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List of Symbols

A, B, D, E, F, H	extensional, bending-extension coupling, bending and higher order stiffness matrices
$A_{ij}^*, B_{ij}^*, D_{ij}^*, E_{ij}^*, F_{ij}^*, H_{ij}^*$	reduced stiffness matrices
a, b	length and width of a plate
d_{31}, d_{32}	piezoelectric strain constants of the k th ply
E_{ij}, F_{ij}, H_{ij}	higher order stiffness matrices
E_{11}, E_{22}	elastic moduli of a single ply
E^f, E^m	Young's moduli of the fiber and the matrix
E_Z	transverse electric field component
\bar{F}, F	stress function and its dimensionless form
G_{12}, G_{13}, G_{23}	shear moduli of a single ply
h	thickness of a plate or shell
\bar{K}_1, K_1, k_1	Winkler foundation stiffness and its two dimensionless forms
\bar{K}_2, K_2, k_2	shearing layer stiffness and its two dimensionless forms
L	length of a shell or beam
P	axial load
p	radial pressure
q	transverse distributed pressure
R	mean radius of a shell
\bar{t}, t	time and its dimensionless form
\bar{U}, \bar{V}	displacement components in the X and Y directions
V_f, V_m	fiber and matrix volume fractions
V_k	applied voltage across the k th ply
\bar{W}, W	deflection of a plate or shell and its dimensionless form
\bar{W}^*, W^*	initial geometric imperfection
X, Y, Z	a coordinate system
x, y, z	dimensionless form of a coordinate system
\bar{Z}, \bar{Z}_B	geometric parameter of a composite or isotropic shell
α_{11}, α_{22}	thermal expansion coefficients in the longitudinal and transverse directions for the k th ply

α^f, α^m	thermal expansion coefficients of the fiber and the matrix
β_{11}, β_{22}	longitudinal and transverse coefficients of hygroscopic expansion for the k th ply
β^f, β^m	swelling coefficients of fiber and matrix
β	aspect ratio of a plate ($= a/b$) or a shell ($= L/\pi R$)
Δ_x, δ_x	end-shortening and its dimensionless form
ε	a small perturbation parameter
λ^*	imperfection sensitivity parameter
λ_p, λ_p^*	dimensionless forms of axial compressive load
λ_q, λ_q^*	dimensionless forms of external pressure
λ_T, λ_T^*	dimensionless forms of thermal stress
μ	imperfection parameter
ν_{12}, ν_{21}	Poisson's ratios of a single ply
ν^f, ν^m	Poisson's ratios of fiber and matrix
ρ, ρ^f, ρ^m	mass density of a plate or shell, fiber and matrix
$\bar{\Psi}_x, \bar{\Psi}_y$	rotations of the normals about the X and Y axes
Ω_L, ω_L	linear frequency and its dimensionless form
Ω_{NL}, ω_{NL}	nonlinear frequency and its dimensionless form

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1

Traditional Perturbation Method

1.1 Introduction

The perturbation method is one of the most appropriate methods which can be used to solve various boundary-value problems in elastic structures. It provides a useful approximate analytical tool for solving a large class of nonlinear equations. The traditional perturbation method is also called the small perturbation method. Using the perturbation method, a complex nonlinear equation may be decomposed into an infinite number of relatively easy ones. In this method, the solution of the original equation is considered as the sum of the solution of each order of perturbation equations and a sequence of terms with increasing power of a small perturbation parameter as their coefficients, so that the first few terms reveal the important feature of the solution. Hence, the solution procedure is convenient compared to solving the original nonlinear equation directly.

The advantage of this method is that it provides solutions to satisfy both governing equations and boundary conditions accurately in the asymptotic sense. Unlike numerical methods, the perturbation approach provides a good physical insight into the problem considered, and the influence of all the parameters on the solution can be assessed easily. The big difference between the perturbation method and other approximate methods, like the Galerkin method and the Ritz method, is that it is not necessary to guess the forms of solutions. In contrast, the accuracy of applying the Ritz and Galerkin methods depends strongly on the choice of the admissible function which does not satisfy all the geometrical and natural boundary conditions, and usually does not satisfy equilibrium equations or motion equations.

The perturbation method is interesting because it can be used for structural nonlinear analysis in various fields such as nonlinear bending, postbuckling and large amplitude vibration of beam, plate and shell structures. However, the successful application of the perturbation method depends largely on the choice of the small perturbation parameter. This perturbation parameter may obviously appear in the original problem or may be introduced by researchers. Usually, the nondimensional load or the nondimensional deflection or both of these are selected as the perturbation parameter in the traditional perturbation method.

1.2 Load-type Perturbation Method

The load-type perturbation method is mainly used in large deflection analysis and postbuckling analysis of plates. Vincent (1931) first studied the large deflection of an isotropic circular plate subjected to uniform pressure by using a load-type perturbation method. In his study, the nondimensional load $[qr^4(1 - \nu^2)/(Eh^4)]$ is taken as a small perturbation parameter, where q is the transverse uniform pressure, h is the plate thickness, r is the radius of the circular plate and E and $\nu (= 0.25)$ are the Young's modulus and Poisson's ratio, respectively, of the plate. The boundary condition is assumed to be simply supported with or without in-plane displacements, referred to as "movable" and "immovable", respectively. The load-deflection relationship obtained by Vincent (1931) may be written as

$$\frac{W_m}{h} = 0.738 \left(\frac{qr^4}{Eh^4} \right) - 0.122 \left(\frac{qr^4}{Eh^4} \right)^3 + 0.0662 \left(\frac{qr^4}{Eh^4} \right)^5 \quad (\text{movable}) \quad (1.1)$$

and

$$\frac{W_m}{h} = 0.738 \left(\frac{qr^4}{Eh^4} \right) - 0.766 \left(\frac{qr^4}{Eh^4} \right)^3 + 2.36 \left(\frac{qr^4}{Eh^4} \right)^5 \quad (\text{immovable}) \quad (1.2)$$

where W_m is the maximum deflection of the plate.

The solutions of Equations 1.1 and 1.2 are little better than the solutions obtained by Chien (1954), in which the nondimensional central deflection (W_c/h) is used as the perturbation parameter. This is because there exists a great discrepancy between the experimental results and the theoretical predictions of Vincent (1931) when the plate deflection is sufficiently large, as reported by Chen and Guang (1981).

In contrast, Stein (1959) studied the postbuckling behavior of an isotropic rectangular plate subjected to uniaxial compression by using a load-type perturbation method. In his study, the nondimensional load $[(P - P_{cr})/P_{cr}]^{1/2}$ is taken as a small perturbation parameter, where P_{cr} is the critical buckling load for the same plate under uniaxial compression. The von Kármán equation was expressed in terms of three displacements. The boundary condition is assumed to be simply supported. The postbuckling load-shortening relationship obtained by Stein (1959) may be written as

$$\begin{aligned} \frac{12(1 - \nu^2)b^2}{4\pi^2h^2} \frac{\Delta_x}{a} &= \frac{Pb}{4\pi^2D} + \frac{1}{2} \frac{m^2}{\beta^2} \left(\frac{(m^2 + n^2\beta^2)^2}{m^4 + n^4\beta^4} \right) \left(\frac{P - P_{cr}}{P_{cr}} \right) + \frac{3}{2} \frac{m^2}{\beta^2(m^4 + n^4\beta^4)} \\ &\times \left[\frac{m^8}{(m^2 + 9n^2\beta^2)^2 - (m^2 + n^2\beta^2)^2} - \frac{n^8\beta^8}{(9m^2 + n^2\beta^2)^2 - 9(m^2 + n^2\beta^2)^2} \right] \\ &\times \left(\frac{(m^2 + n^2\beta^2)^2}{m^4 + n^4\beta^4} \right)^2 \left(\frac{P - P_{cr}}{P_{cr}} \right)^2 \end{aligned} \quad (1.3)$$

in which a and b are the length and width of the plate, $\beta = a/b$ is the plate aspect ratio, Δ_x is the plate end-shortening displacement in the X direction and $D = Eh^3/[12(1 - \nu^2)]$ is the flexural rigidity of the plate.

This load-type perturbation method was then extended to the case of postbuckling analysis of an orthotropic rectangular plate by Chandra and Raju (1973). The postbuckling load-shortening relationship was obtained for a perfect plate under uniaxial compression. Although the resultant expression for an isotropic plate is coincident with that included in the work of Stein (1959), the higher order term in the solution of Chandra and Raju (1973) is incorrect, as reported by Blazquez and Picon (2010).

From the load-deflection curve of the circular plate, the condition of $[qr^4(1 - \nu^2)/Eh^4] = 1$ is equivalent to $(W_c/h) = 0.1-0.2$, and this condition can easily be exceeded in the large deflection region. In contrast, the condition of $P < 2P_{cr}$ is easily satisfied for most plates in the postbuckling region, and therefore, the load-type perturbation method is better for use in the postbuckling analysis than in the large deflection analysis of a plate. As has been shown (Zhang and Fan, 1984), in many cases when the load-type perturbation method is used, the postbuckling load-deflection curve does not converge to the exact solution when the plate deflection is sufficiently large. Hence, it is not a good option for nonlinear analysis of plates by using the load-type perturbation method.

1.3 Deflection-type Perturbation Method

Chien (1947) is the pioneer in studying the large deflection of circular plates by using the deflection-type perturbation method. For an isotropic circular plate ($\nu = 0.3$) with a movable in-plane boundary condition, the load-deflection relationship obtained by Chien (1954) may be written as

$$\frac{qr^4}{Eh^4} = 0.87 \left(\frac{W_c}{h} \right) + 0.09152 \left(\frac{W_c}{h} \right)^3 \quad (1.4)$$

where W_c is the central deflection of the plate.

This method is easy to follow and has been applied successfully to solve many large deflection problems of plates. For example, Yeh (1953) presented the large deflection analysis of annular plates. Chien and Yeh (1954) presented the large deflection analysis of circular plates with various boundary conditions under uniformly distributed or concentrated load. Hu (1954) presented the large deflection analysis of circular plates under the combined action of uniformly distributed and concentrated loads. Chien et al. (1992) presented the large deflection analysis of elliptical plates with clamped boundary conditions subjected to uniform pressure. All these important contributions are of interest to the research community.

The large deflection analysis of rectangular plates is more complicated than that of circular plates. Chien and Yeh (1957) presented the large deflection analysis of an isotropic rectangular plate with clamped boundary conditions subjected to uniform pressure by using the deflection-type perturbation method, in which the nondimensional central deflection (W_c/h) is taken as a small perturbation parameter. By solving the von Kármán equation expressed in

terms of three displacements, the load–deflection relationship for an isotropic square plate ($\nu = 1/3$) can be written as

$$\frac{qa^4}{Dh} = 50.3815 \left(\frac{W_c}{h} \right) + 24.37716 \left(\frac{W_c}{h} \right)^3 \quad (1.5)$$

Similarly, Kan and Huang (1967) presented the large deflection analysis of a sandwich plate with clamped boundary conditions subjected to uniform pressure. By solving the nonlinear equation expressed in terms of three displacements, the load–deflection relationship for a sandwich square plate ($\nu = 0.3$) can be written as

$$\frac{qa^3}{Eh_f h_c^2} = 8.357 \left(\frac{W_c}{h_c} \right) + 3.8532 \left(\frac{W_c}{h_c} \right)^3 \quad (1.6)$$

where h_f and h_c are the thicknesses of the face sheet and core layer.

Chia (1980) wrote a good book for the nonlinear analysis of composite thin plates. This book provides a lot of examples for the large deflection analysis of orthotropic rectangular plates (Chia, 1972a), orthotropic circular plates (Nowinski, 1960), orthotropic elliptical plates (Prabhakara and Chia, 1975) and anisotropic rectangular plates (Chia, 1972b).

Moreover, Dym and Hoff (1968) studied the postbuckling of an isotropic cylindrical shell under axial compression by using the deflection-type perturbation method, in which the non-dimensional maximum deflection (W_m/h) is taken as a small perturbation parameter. For a mixed boundary-value problem of elastic cylindrical shells, the Kármán-type equation expressed in terms of a transverse displacement \bar{W} and a stress function \bar{F} is more convenient than that expressed in terms of three displacements \bar{U} , \bar{V} and \bar{W} . By solving the Kármán-type equations, the asymptotic solutions up to fourth order for the postbuckling load–shortening relationship were obtained.

Actually, in Koiter's initial postbuckling theory (Koiter, 1945, 1963), the large deflection solution of an isotropic cylindrical shell was first determined by using the deflection-type perturbation method and then performed the imperfection-sensitive analysis of the same cylindrical shell under mechanical loads, as reported by Budiansky and Amazigo (1968). Like in the case of Dym and Hoff (1968), these solutions can not predict the full postbuckling equilibrium path of the cylindrical shell. The applications of a similar solution methodology could be found in the free and forced vibration analyses of elastic structures (Rehfield, 1973, 1974).

1.4 Multi-parameter Perturbation Method

Besides the single-parameter perturbation method as described in Sections 1.2 and 1.3, a multi-parameter perturbation method is also sometimes used in the nonlinear analysis of elastic structures. Among those, Hu (1954) presented the large deflection analysis of circular plates under combined action of uniformly distributed and concentrated loads. In his study, both nondimensional uniform pressure (qr^4/Eh^4) and nondimensional concentrated load ($Pr^2/\pi Eh^4$) were taken as two small perturbation parameters. In such a case, the solution

procedure is more complicated. He found that the solution is poor when it converges slowly or can actually be divergent when these two perturbation parameters are not very small. Chien (2002) presented the large deflection analysis of a cantilever beam subjected to a uniform pressure. Unlike in the case of Hu (1954), nondimensional uniform pressure ($qL^3/12EI$) and nondimensional end displacement ($6\Delta/L$) were taken as two small perturbation parameters, where EI is the flexural rigidity of the beam, L is the undeformed length of the beam and Δ is the vertical displacement at the free end. Andrianov et al. (2005) presented the nonlinear natural in-plane vibrations of an isotropic rectangular plate with clamped boundary conditions by using a three-parameter perturbation method. Other applications of multi-parameter perturbation method could be found in Nowinski and Ismail (1965). In most cases, it is unnecessary to use multi-parameter perturbation method when the relationship of these perturbation parameters could be established.

1.5 Limitations of the Traditional Perturbation Method

In the traditional perturbation method, the nondimensional generalized displacement, for example, the mean square root of deflection or the mean square root of the slope, is also taken as a small perturbation parameter instead of the nondimensional load or nondimensional deflection (Hu, 1954; Chen and Guang, 1981). The comparison studies for large deflection of clamped circular plates (Schmidt and DaDeppo, 1974; Chen and Guang, 1981; Zheng, 1990) show that the perturbation solution derived by using the mean square root of the slope as a perturbation parameter is better than that derived by using the nondimensional load as a perturbation parameter, whereas the perturbation solution derived by using the central deflection as a perturbation parameter is the best one among others. However, Hu (1954) pointed out that the nondimensional central deflection is not a better choice for a circular plate subjected to the combined action of uniformly distributed and concentrated loads. This is due to the fact that, in such a case, the central deflection may be zero valued. Further, Vol'mir (1967) reported that there exists a depression phenomenon in the central region of the deflection curve of Chien (1954) when the plate deflection is sufficiently large. In fact, these two weaknesses can easily be improved by using the maximum deflection instead of the central deflection and replacing the linear solution properly or considering more terms in the perturbation expansion series.

Generally, it is necessary to have $\varepsilon < 1$ in the traditional perturbation method. It is worth noting that ε is no longer a small perturbation parameter in the large deflection region when the plate deflection is sufficiently large, that is, $W_m/h > 1$, or in the deep postbuckling region when the applied load is larger than two times the buckling load, that is, $(P - P_{cr})/P_{cr} > 1$, and in such a case the solution may be invalid. Blazquez and Picon (2010) reported that the two solutions based on the revised method of Chandra and Raju and the method of Shen and Zhang agree well when $P < 2P_{cr}$, whereas a discrepancy could be observed when $P > 2P_{cr}$. This is due to the fact that the revised method of Chandra and Raju is a load-type perturbation method where $[(P - P_{cr})/P_{cr}]^{1/2}$ is taken to be a small perturbation parameter, and the solution may also be invalid when $P > 2P_{cr}$. Although the theoretical limitation is that $\varepsilon < 1$, the perturbation solution of Chien is adequate for the large deflection region, even if $\varepsilon = W_c/h$ reaches 4, when compared with experimental results. It seems reasonable to conclude that the perturbation method can be used for solving stronger nonlinear problems virtually.

In order to satisfy the condition $\varepsilon < 1$, the small perturbation parameter was assumed to be $\varepsilon = h/a$ in the large amplitude vibration analysis of the plate (Bhimaraddi, 1989, 1992, 1993), or was assumed to be $\varepsilon = \bar{W}/R$ in the large amplitude vibration analysis of the shell (Chen and Babcock, 1975), where R is the mean radius of the shell.

In order to overcome the weakness of the traditional perturbation method in the nonlinear analysis of elastic structures, Shen and Zhang (1988) proposed a two-step perturbation method. This approach gives explicit analytical expressions of all the variables in the post-buckling range. In contrast to the traditional perturbation scheme, this method avoids the paradox by a two-step perturbation scheme. In the first step ε may have no physical meaning, but is definitely a small perturbation parameter. In the second step $(A_{11}^{(1)}\varepsilon)$ is taken as the second perturbation parameter relating to the nondimensional maximum deflection that may be large in the large deflection region or in the deep postbuckling region, where $A_{11}^{(1)}$ is the amplitude of the first term in the perturbation expansion of the plate deflection. This approach is successfully used in solving many nonlinear bending, postbuckling and nonlinear vibration problems of beams, plates and shells. This approach is now called the “Method of Shen and Zhang” by Blazquez and Picon (2010). This approach may find more extensive applications in the nonlinear analysis of nanoscale structures (Shen, 2010a, b, 2011; Shen et al., 2010, 2011; Shen and Zhang, 2006, 2007, 2010a, b).

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